

Effects of Λ NN three-body force in B_{Λ} values of hypernuclei

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Collaborators: Y. Yamamoto, Th.A. Rijken

Grand challenges of hypernuclear physics

Interaction: To understand baryon-baryon interaction

- 2 body interaction between baryons (nucleon, hyperon)

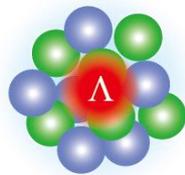
- hyperon-nucleon (YN)
 - hyperon-hyperon (YY)
- } A major issue in hypernuclear physics

Structure: To understand many-body system of nucleons and hyperon

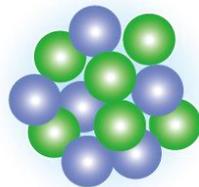
- Addition of hyperon(s) shows us new features of nuclear structure

Ex.) Structure change by hyperon(s)

- No Pauli exclusion between N and Y
 - YN interaction is different from NN
- } “Hyperon as an impurity in nuclei”



Λ hypernucleus



Normal nucleus

+



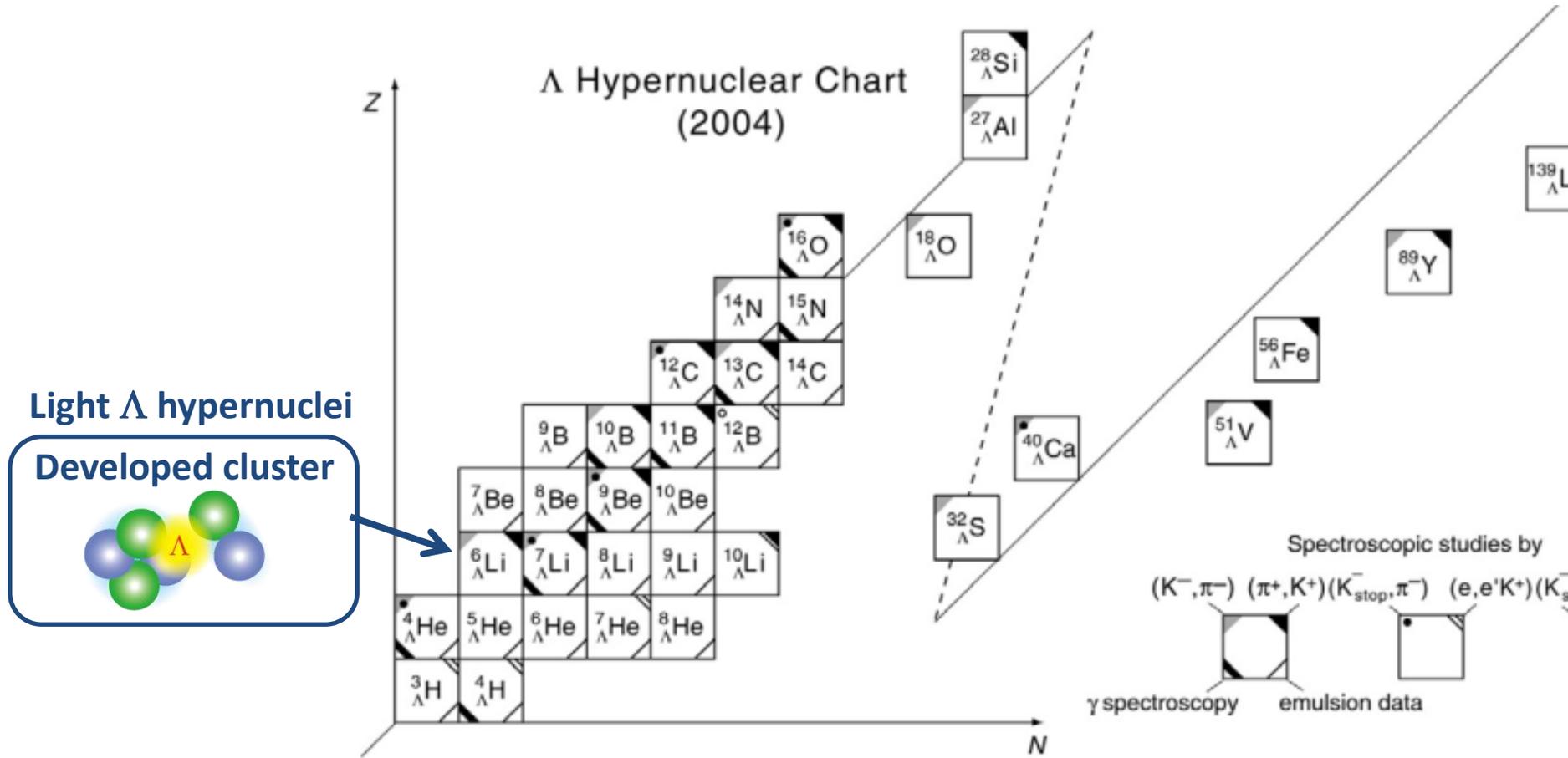
As an impurity

Today's talk: “structure of Λ hypernuclei” and “ Λ binding energy”

Structure of Λ hypernuclei

Λ hypernuclei observed so far

- Concentrated in light Λ hypernuclei
- Most of them have well pronounced cluster structure

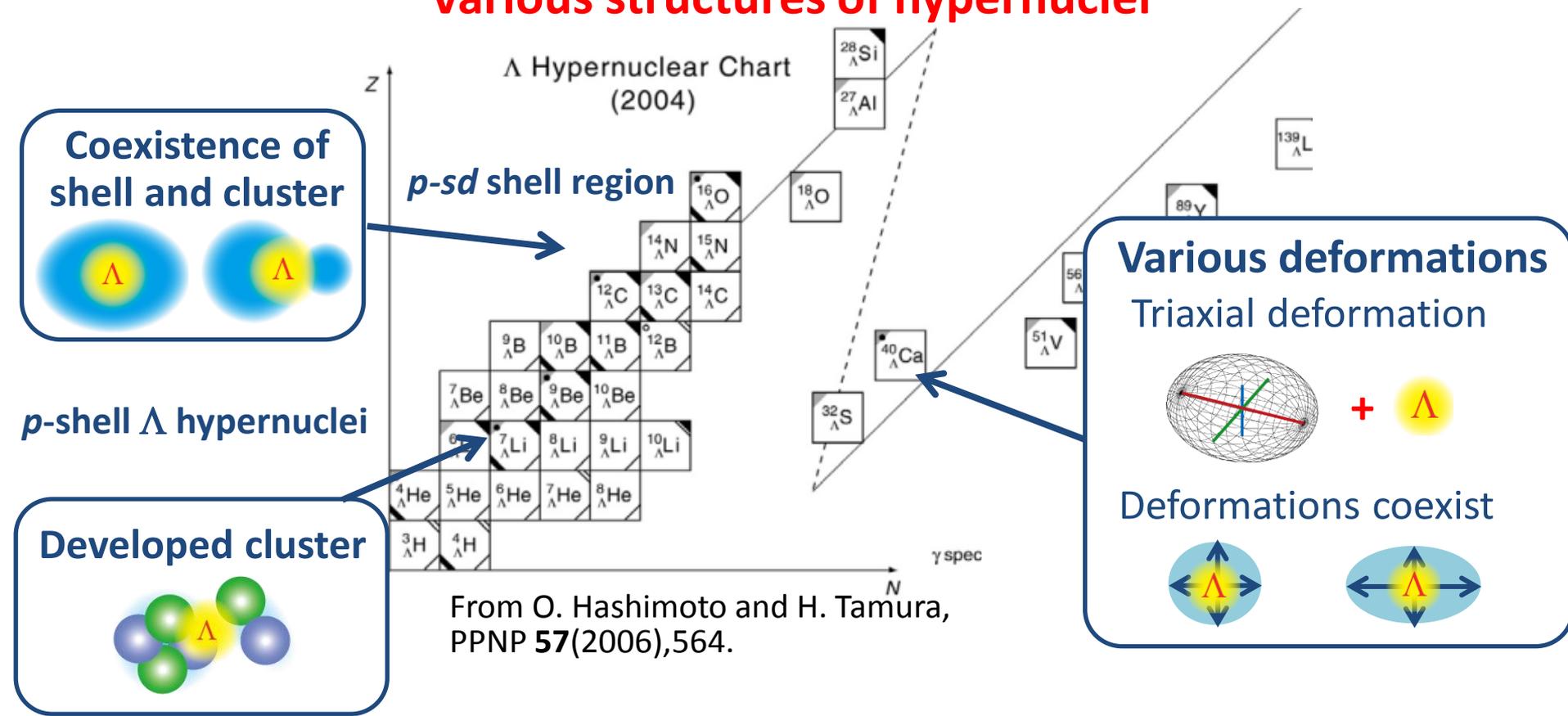


Toward heavier and exotic Λ hypernuclei

Experiments at J-PARC, JLab and Mainz etc.

- Hypernuclear chart will be extended to heavier regions

“Various structures of hypernuclei”

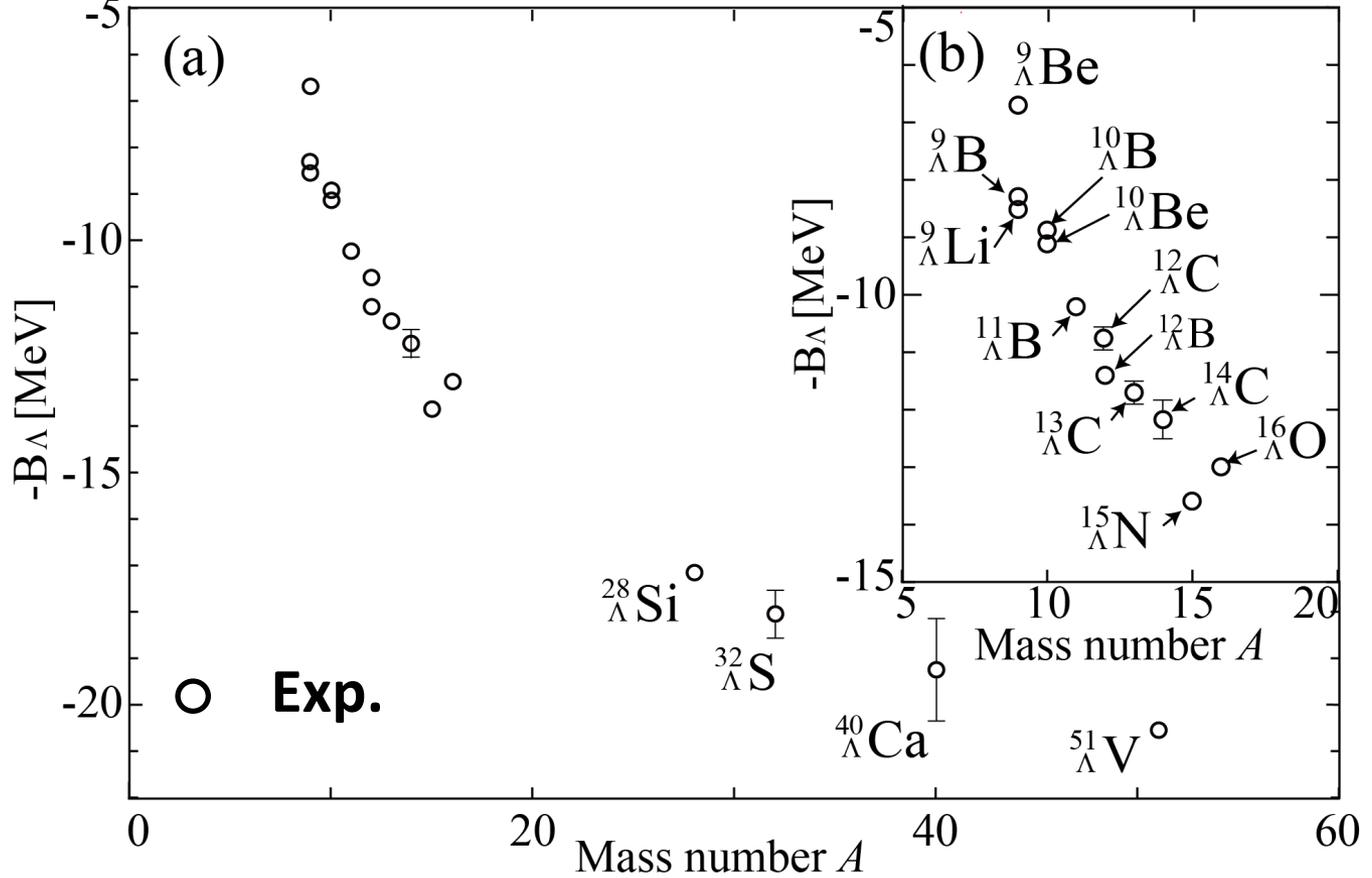


How do core nuclei affect the mass dependence of B_{Λ} ?

“clustering/deformations”, “Density dependence of interactions”

B_Λ as a function of mass number A

Observed data of Λ binding energy B_Λ ($9 \leq A \leq 51$)



Do core nuclei affect the mass dependence of B_Λ ?

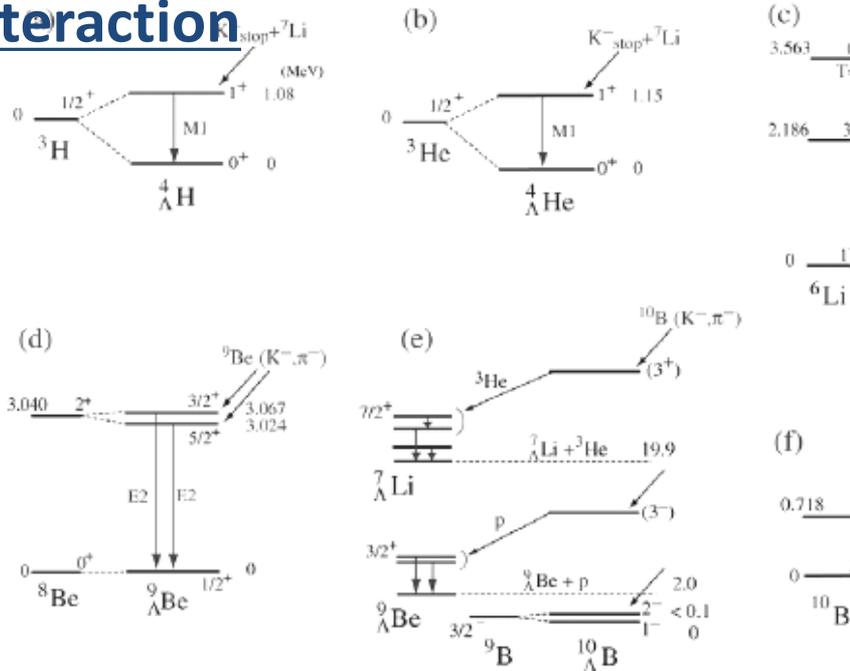
“clustering/deformations”, “Density dependence of interactions”

Bertini *et al.*, NPA83,306(1979), Davis, Juric, *et al.*, NPB52(1973), Davis, NPA547,369(1992);NPA754,3c(2005), Ajimura *et al.*, NPA639(1998)93c, Pile *et al.*, PRL66,2585(1991), Hotchi *et al.*, PRC64, 044302(2001), Hashimoto and Tamura, PPNP57,564(2006), Tang, *et al.*, PRC90,034320(2014).

Background of our research

Knowledge of ΛN two-body effective interaction

- Study of light (*s*, *p*-shell) Λ hypernuclei
 - Accurate solution of few-body problems [1]
 - ΛN G-matrix effective interactions [2]
 - Increases of experimental information [3]



Development of theoretical models

- Through the study of unstable nuclei
 - Ex.: Antisymmetrized Molecular Dynamics (AMD)[4]
 - AMD can describe dynamical changes of various structure
 - No assumption on clustering and deformation

AMD calculation by using YNG-ESC

[1] E. Hiyama, NPA **805** (2008), 190c, [2] Y. Yamamoto, *et al.*, PTP Suppl. **117** (1994), 361., [3] O. Hashimoto and H. Tamura, PPNP **57** (2006), 564., [4] Y. Kanada-En'yo *et al.*, PTP **93** (1995), 115.

We extended the AMD to hypernuclei

HyperAMD (Antisymmetrized Molecular Dynamics for hypernuclei)

◆ Hamiltonian

$$\hat{H} = \hat{T}_N + \hat{V}_{NN} + \hat{T}_\Lambda + \hat{V}_{\Lambda N}$$

NN: Gogny D1S

Λ N: YNG interaction

◆ Wave function

● Nucleon part: Slater determinant

Spatial part of single particle w.f. is described as Gaussian packet

$$\varphi_N(\vec{r}) = \frac{1}{\sqrt{A!}} \det[\varphi_i(\vec{r}_j)]$$

$$\varphi_i(r) \propto \exp\left[-\sum_{\sigma=x,y,z} v_\sigma (r - Z_i)_\sigma^2\right] \chi_i n_i \quad \chi_i = \alpha_i \chi_\uparrow + \beta_i \chi_\downarrow$$

● Single particle w.f. of Λ hyperon:

Superposition of Gaussian packets

$$\varphi_\Lambda(r) = \sum c_m \varphi_m(r)$$

$$\varphi_m(r) \propto \exp\left[-\sum_{\sigma=x,y,z} \mu v_\sigma (r - z_m)_\sigma^2\right] \chi_m \quad \chi_m = a_m \chi_\uparrow + b_m \chi_\downarrow$$

● Total w.f.:

$$\psi(\vec{r}) = \sum_m c_m \varphi_m(r_\Lambda) \otimes \frac{1}{\sqrt{A!}} \det[\varphi_i(\vec{r}_j)]$$

Λ N interaction used

YNG interaction derived from ...

ESC08c + **MPP + TBA**
 2-body including Λ N- Σ N coupling **Many-body force**

MPP: repulsion which is essential with high dens.

TBA: phenomenological 3 body attraction

$$V_{\Lambda N}(r; k_F) = \sum_{i=1}^3 (a_i + b_i k_F + c_i k_F^2) \exp(-r^2/\beta_i^2)$$

NNN part of MPP + TBA: scattering data of $^{16}\text{O} + ^{16}\text{O}$, and saturation property

TBA with Λ : to reproduce observed spectra of $^{89}_{\Lambda}\text{Y}$

→ Enough stiff EoS to give $2M_{\odot}$ maximum mass of neutron star

Yamamoto, Furumoto, Yasutake and Rijken,
 PRC**88**,022801(2013); PRC**90**,045805(2014).

| $^{89}_{\Lambda}\text{Y}$ | <i>s</i> | <i>p</i> | <i>d</i> | <i>f</i> |
|---------------------------|----------|----------|----------|----------|
| ESC | -26.5 | -19.0 | -11.4 | -4.3 |
| MPa | -24.0 | -17.3 | -10.5 | -3.9 |
| exp | -23.7 | -17.6 | -10.9 | -3.7 |

exp: O. Hashimoto and H. Tamura, PPNP **57**(2006),564.

◆ k_F determined by density

● Averaged density approximation(ADA):

$$\langle \rho \rangle = \int dr^3 \rho_N(\mathbf{r}) \rho_{\Lambda}(\mathbf{r}) \quad k_F = \left(\frac{3\pi^2 \langle \rho \rangle}{2} \right)^{1/3}$$

No free parameter except for TBA

Theoretical framework: HyperAMD

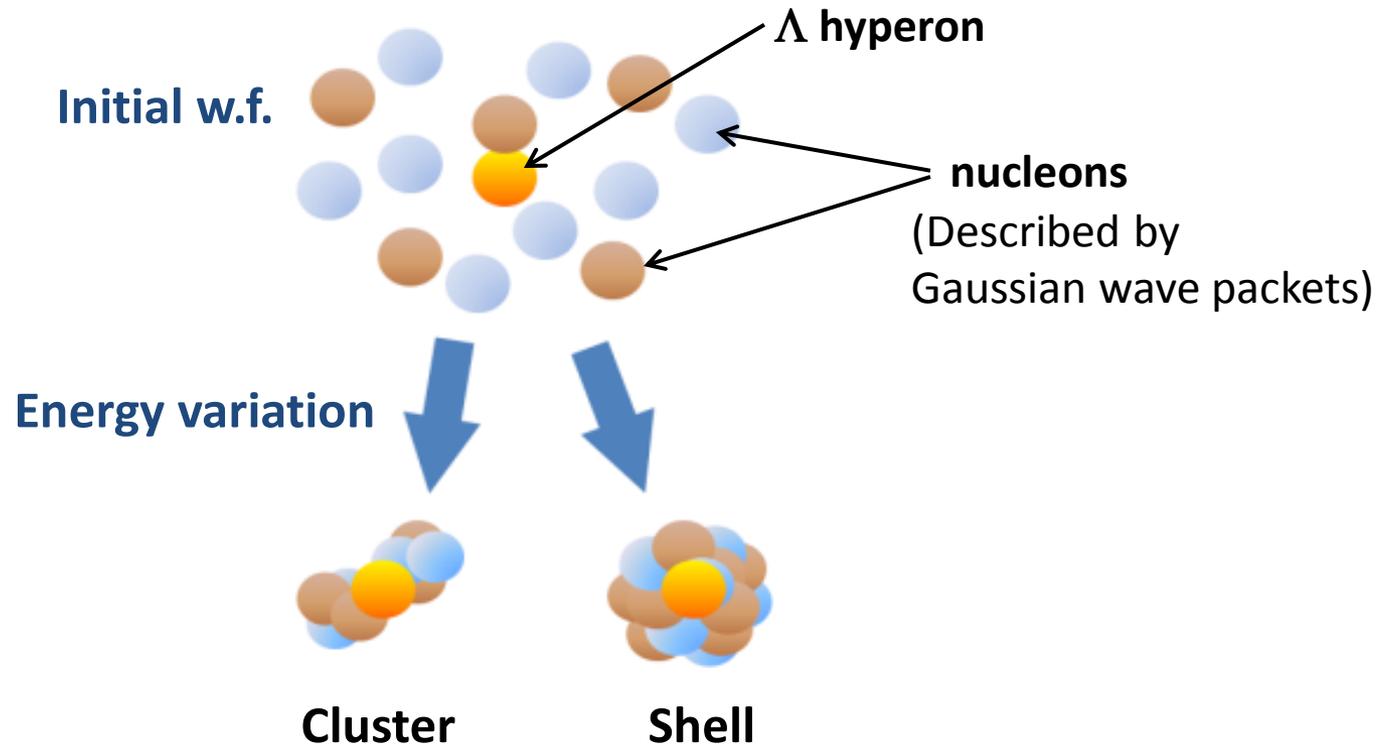
◆ Procedure of the calculation

Variational Calculation

- Imaginary time development method

$$\frac{dX_i}{dt} = \frac{\kappa}{\hbar} \frac{\partial H^\pm}{\partial X_i^*} \quad \kappa < 0$$

- Variational parameters: $X_i = Z_i, z_i, \alpha_i, \beta_i, a_i, b_i, v_i, c_i$



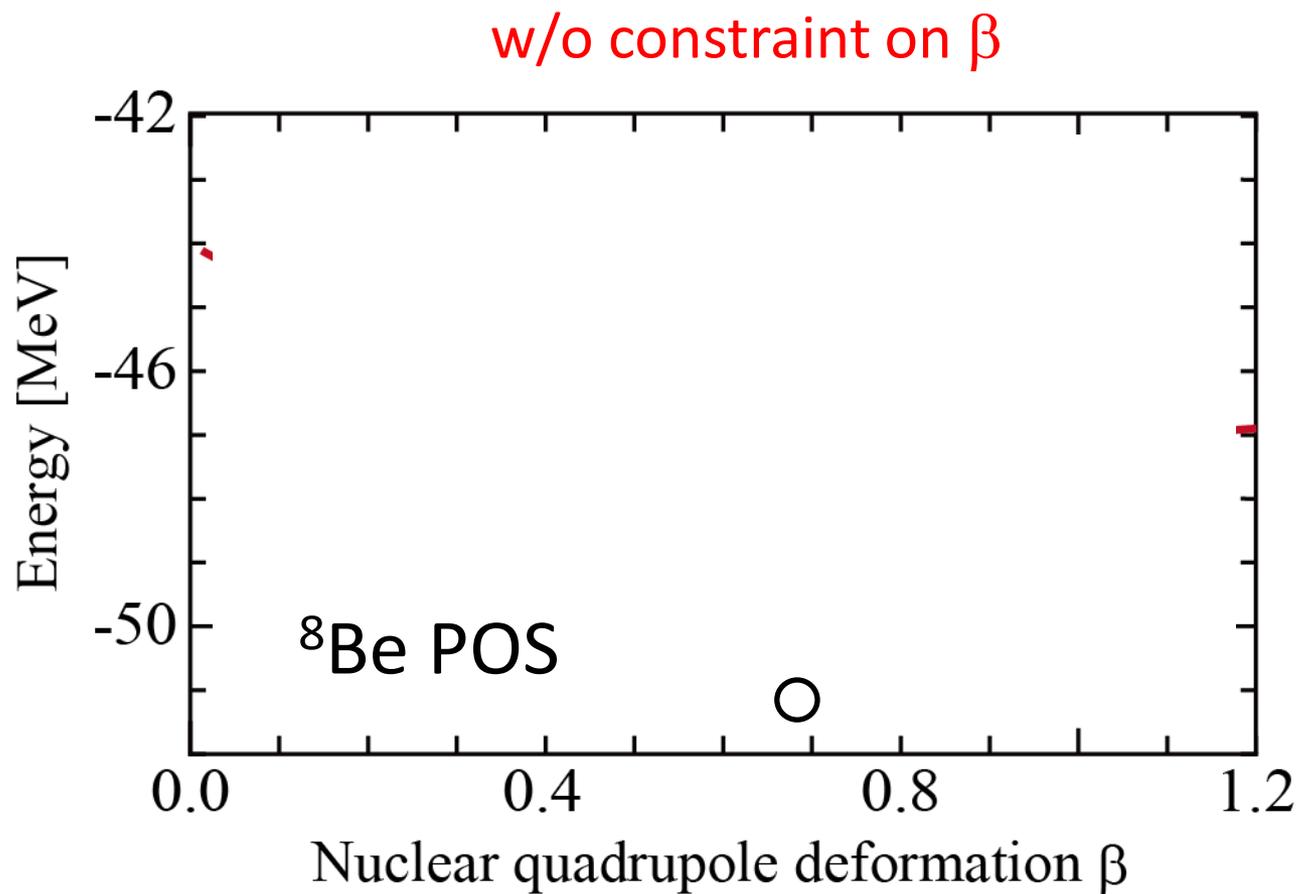
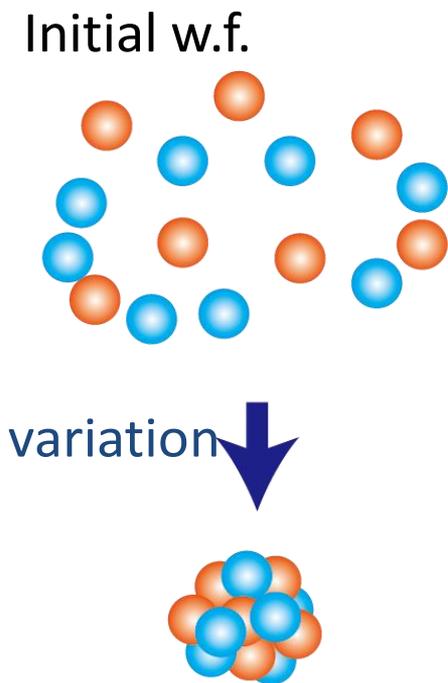
Actual calculation of HyperAMD

M. Isaka, *et al.*, PRC83(2011) 044323

M. Isaka, *et al.*, PRC83(2011) 054304

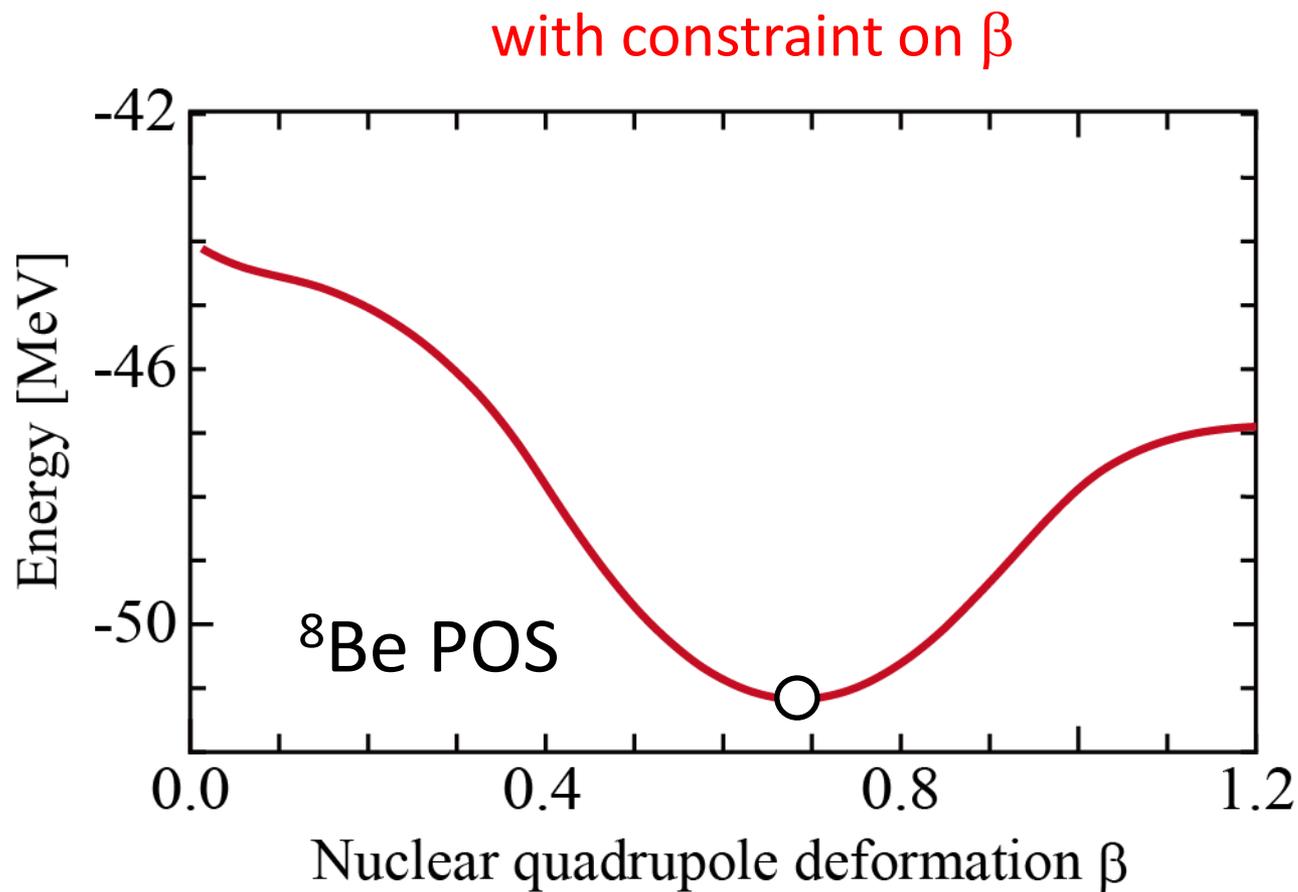
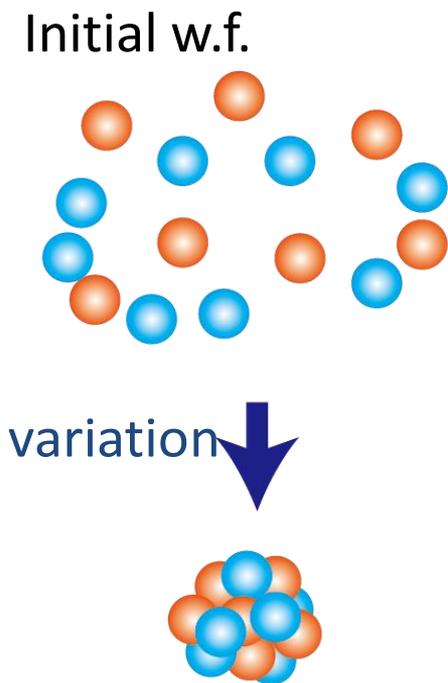
Energy variation with constraint on nuclear quadrupole deformation

Ex.) ^8Be



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Actual calculation of HyperAMD

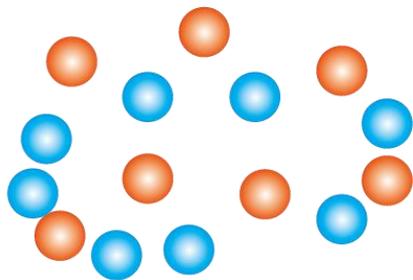
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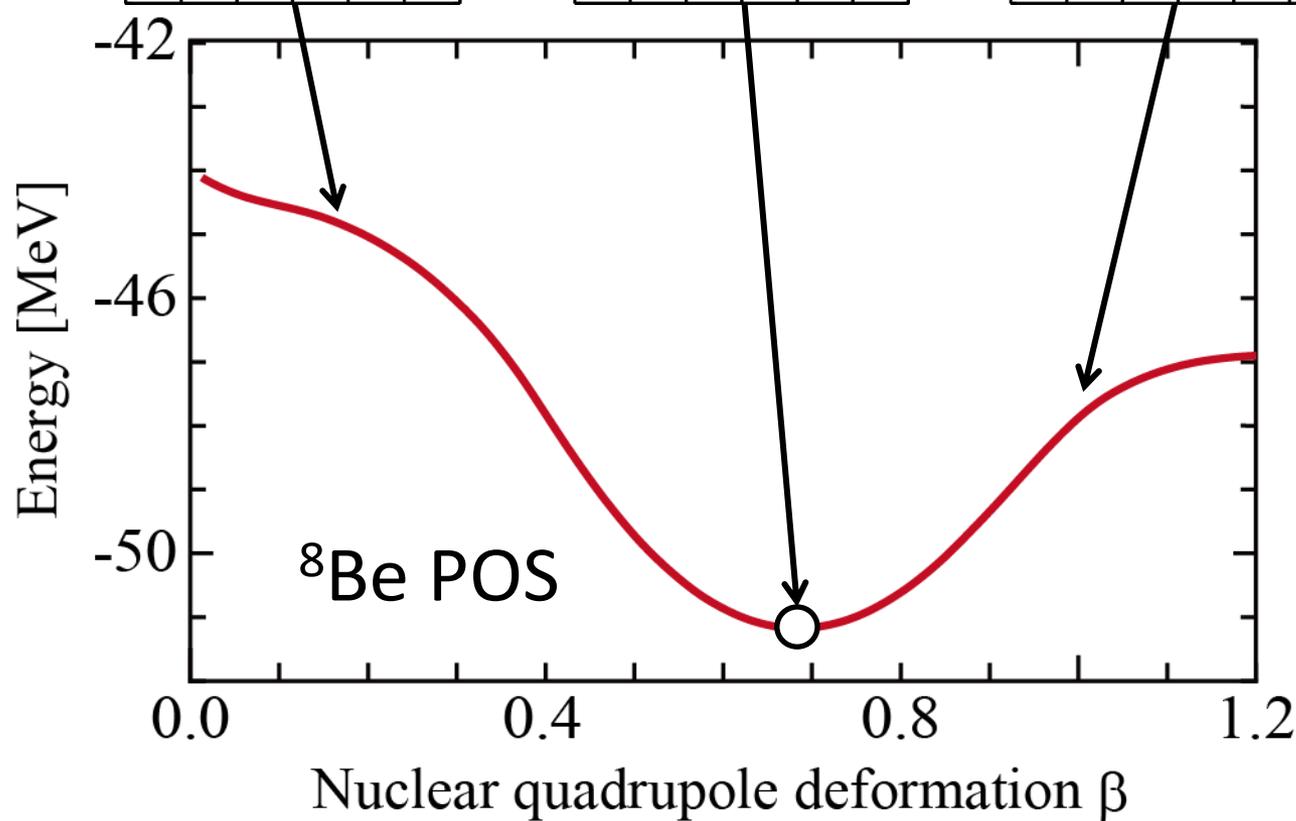
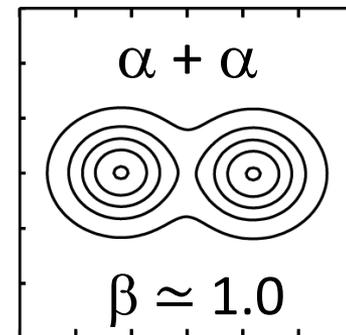
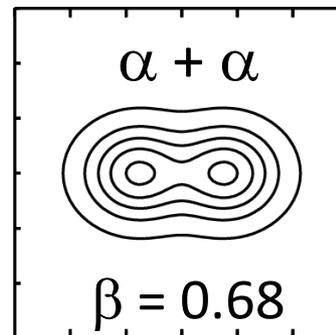
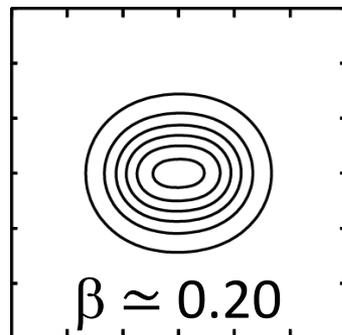
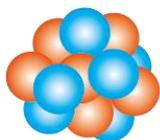
Energy variation with constraint on nuclear quadrupole deformation

Ex.) ^8Be

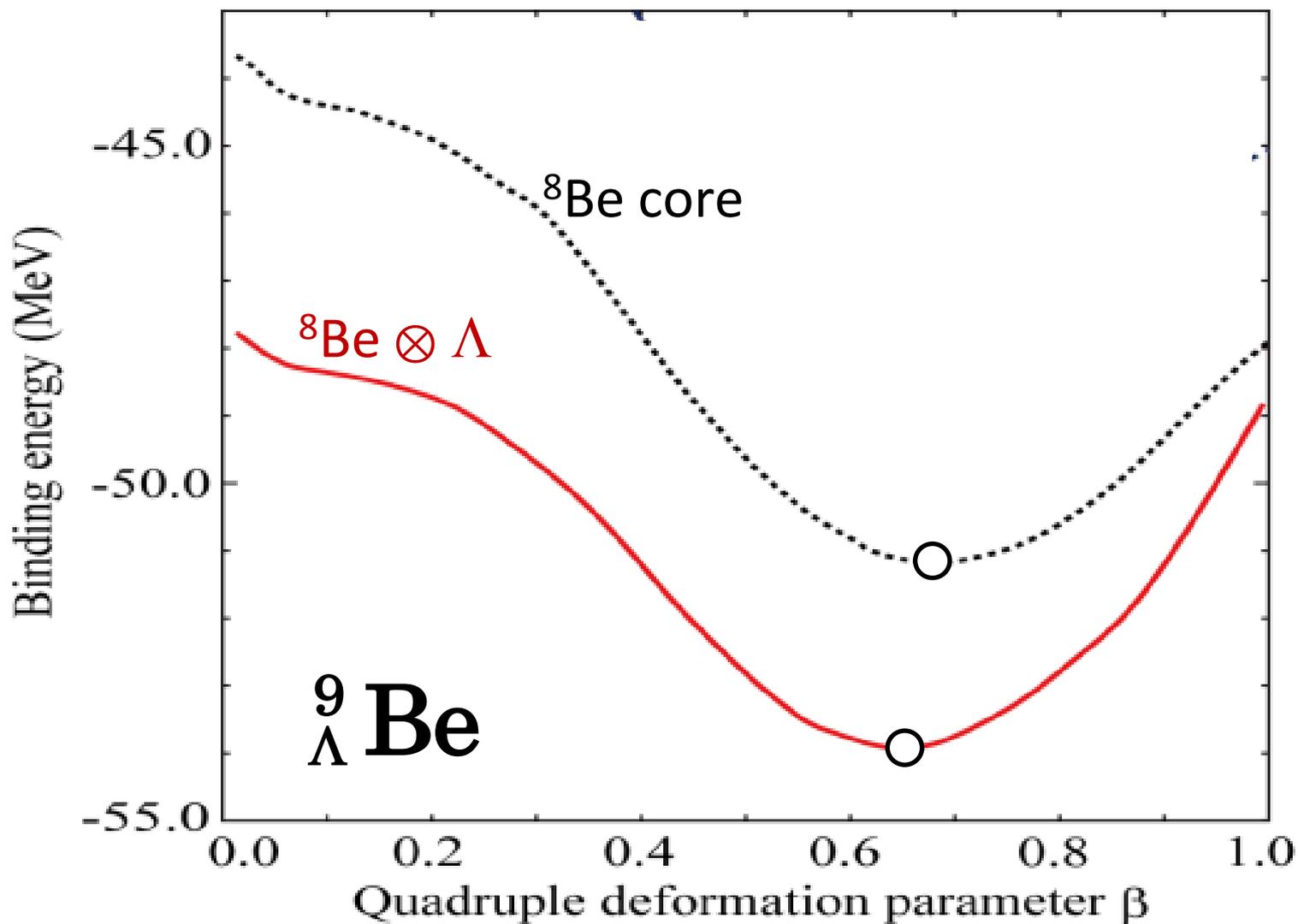
Initial w.f.



variation



◆ For hypernuclei



Theoretical framework: HyperAMD

◆ Procedure of the calculation

Variational Calculation

- Imaginary time development method $\frac{dX_i}{dt} = \frac{\kappa}{\hbar} \frac{\partial H^\pm}{\partial X_i^*}$ $\kappa < 0$
- Variational parameters: $X_i = Z_i, z_i, \alpha_i, \beta_i, a_i, b_i, v_i, c_i$

Angular Momentum Projection

$$\left| \Phi_K^s; JM \right\rangle = \int d\Omega D_{MK}^{J*}(\Omega) R(\Omega) \left| \Phi^{s+} \right\rangle$$

Generator Coordinate Method(GCM)

- Superposition of the w.f. with different configuration
- Diagonalization of $H_{sK,s'K'}^{J^\pm}$ and $N_{sK,s'K'}^{J^\pm}$

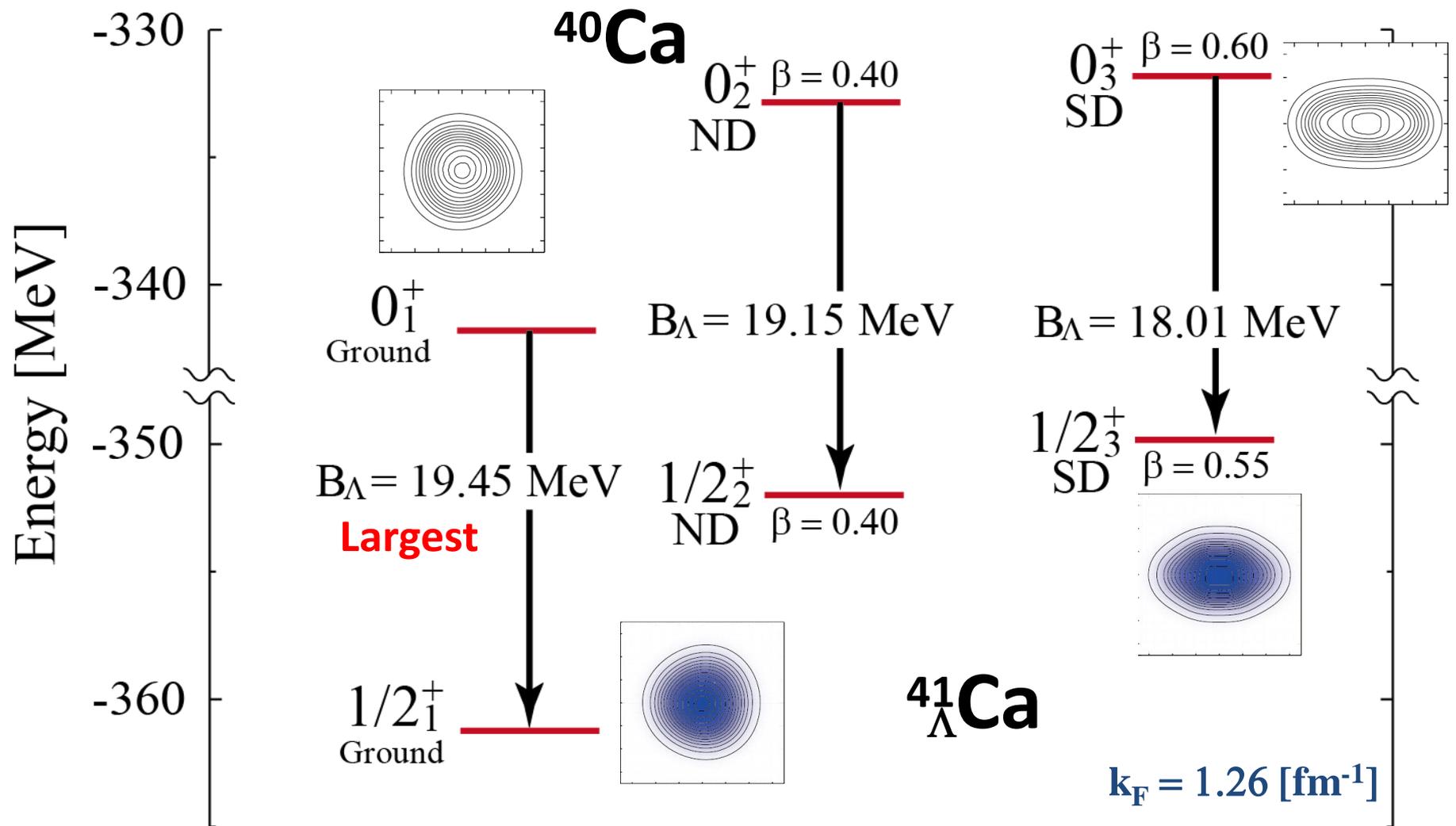
$$H_{sK,s'K'}^{J^\pm} = \left\langle \Phi_K^s; J^\pm M \left| \hat{H} \right| \Phi_{K'}^{s'}; J^\pm M \right\rangle$$
$$N_{sK,s'K'}^{J^\pm} = \left\langle \Phi_K^s; J^\pm M \left| \Phi_{K'}^{s'}; J^\pm M \right\rangle$$
$$\left| \Psi^{J^\pm M} \right\rangle = \sum_{sK} g_{sK} \left| \Phi_K^s; J^\pm M \right\rangle$$

Applications of HyperAMD ($^{41}_{\Lambda}\text{Ca}$)

Example: $^{41}_{\Lambda}\text{Ca}$

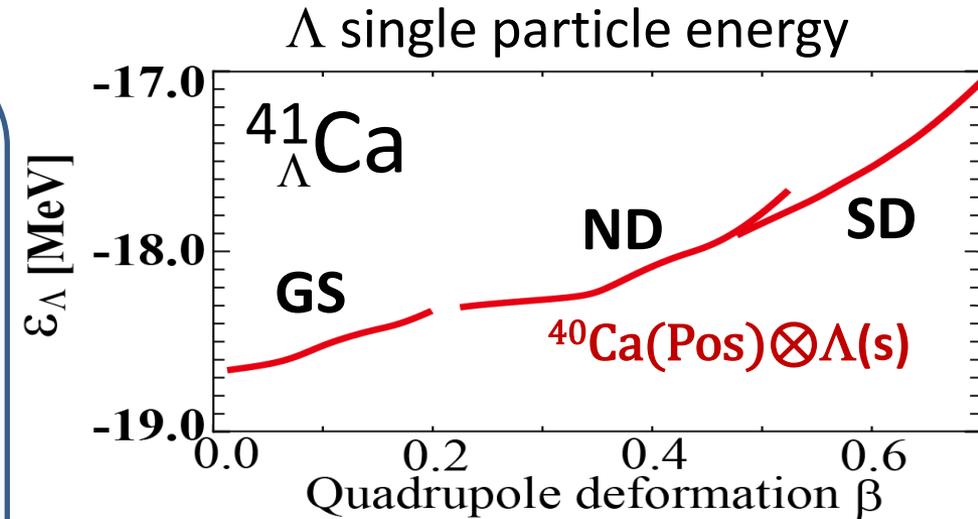
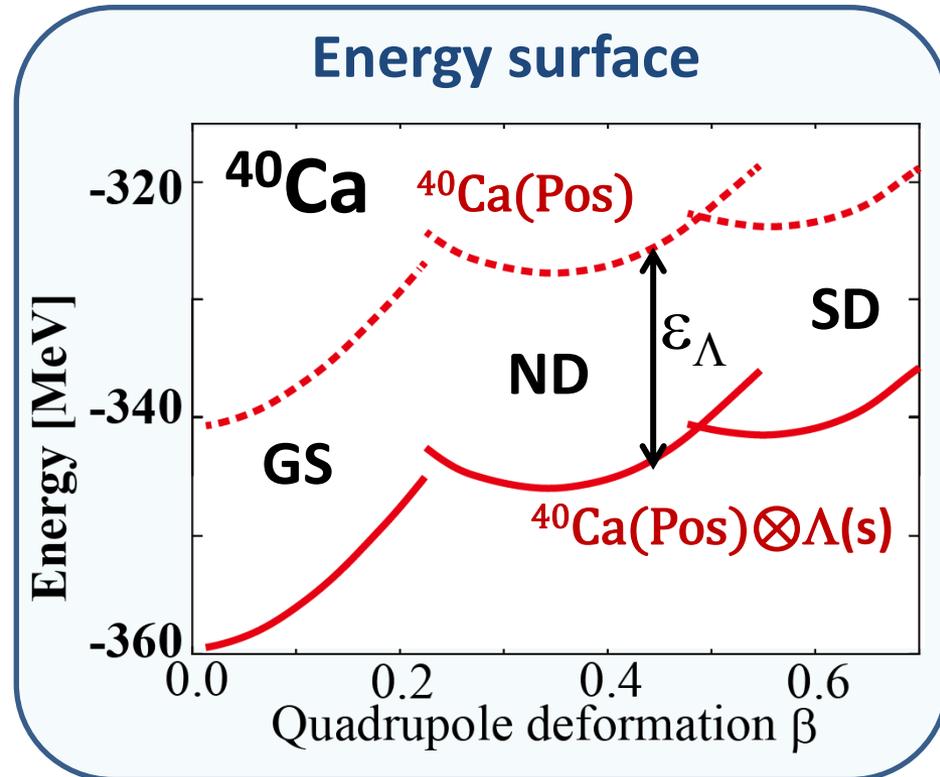
B_{Λ} is dependent on nuclear deformation

M. Isaka, *et al.*, PRC89, 024310(2014)



Applications of HyperAMD ($^{41}_{\Lambda}\text{Ca}$)

- **Definition:** $\epsilon_{\Lambda}(\beta) = E(^{41}_{\Lambda}\text{Ca})(\beta) - E(^{40}\text{Ca})(\beta)$
- ϵ_{Λ} changes within 1 - 2 MeV as β increases



| | $-B_{\Lambda}$ | T_{Λ} | $V_{\Lambda N}$ |
|----|----------------|---------------|-----------------|
| GS | -18.7 | 5.5 | -24.3 |
| ND | -18.2 | 5.5 | -23.9 |
| SD | -17.7 | 5.6 | -23.4 |

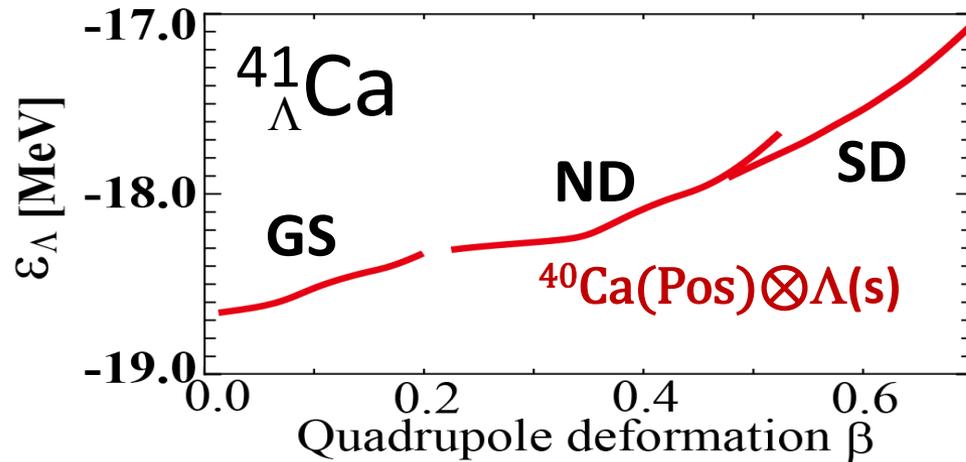
Difference of ϵ_{Λ} is mainly coming from ΛN potential energy

Why? — “overlap between Λ and core”

Applications of HyperAMD ($^{41}_{\Lambda}\text{Ca}$)

- ε_{Λ} varies due to changes of overlap between Λ and N

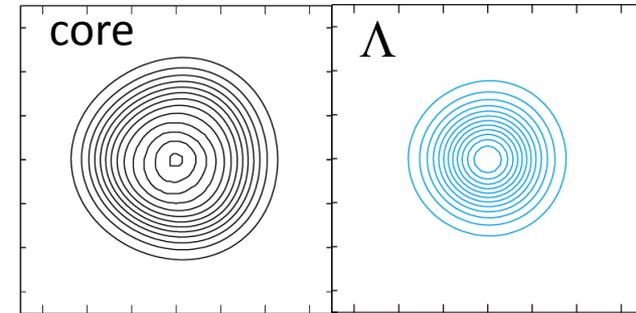
– Deformation of Λ distribution is small, while nuclear part is deformed



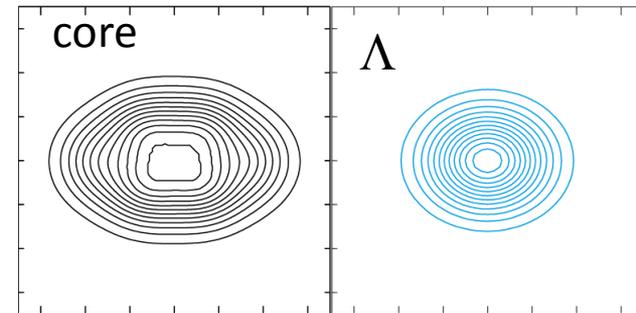
Overlap between the Λ and nucleons

$$I = \int d^3r \rho_N(\mathbf{r}) \rho_{\Lambda}(\mathbf{r})$$

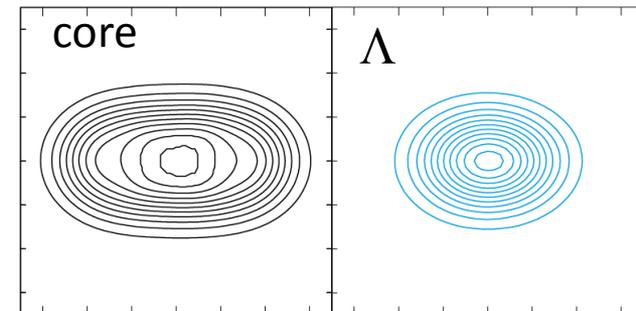
GS
 $\beta = 0.01$
 $I = 0.1364$



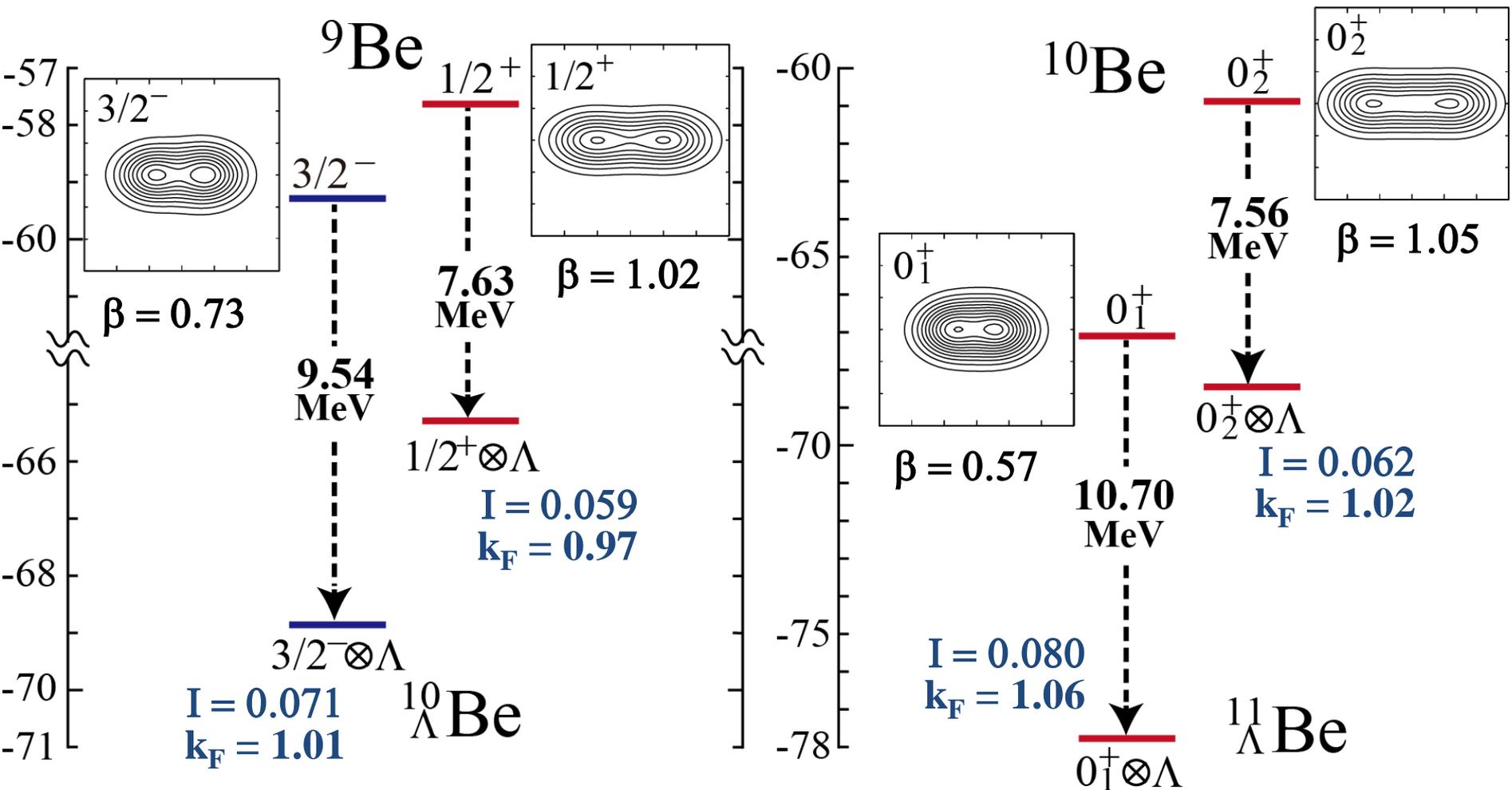
ND
 $\beta = 0.35$
 $I = 0.1356$



SD
 $\beta = 0.55$
 $I = 0.1336$



Decrease of overlap makes $V_{\Lambda N}$ shallower I [fm^{-3}]

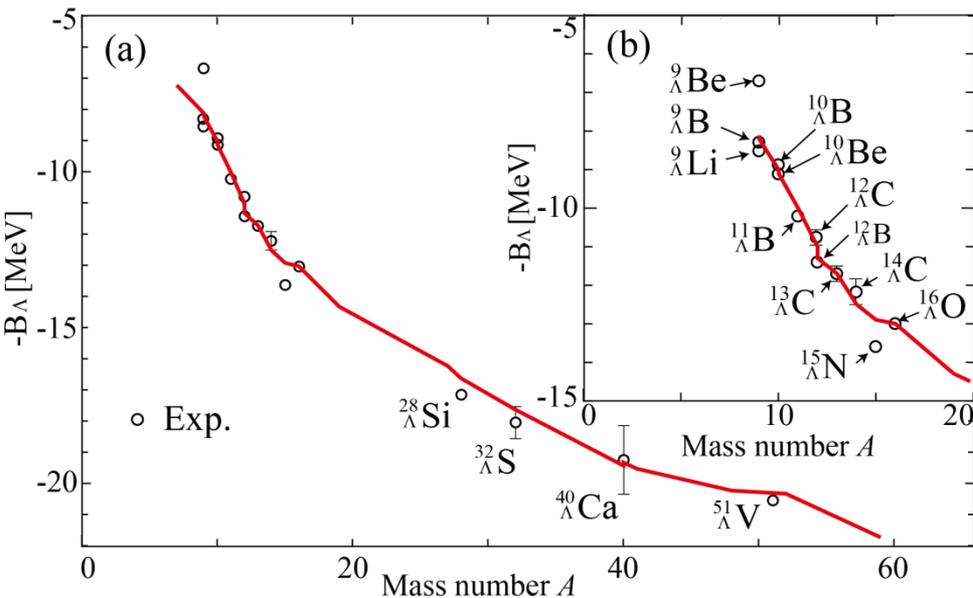
Applications of HyperAMD (${}_{\Lambda}$ Be)Example: ${}^{10}_{\Lambda}\text{Be}$, ${}^{11}_{\Lambda}\text{Be}$ Dependence of B_{Λ} on cluster structures I [fm^{-3}], k_F [fm^{-1}]Same trend as in ${}^{41}_{\Lambda}\text{Ca}$, whereas smaller k_F is used in excited states

Results and Discussions

“mass dependence of B_Λ ”

B_Λ as a function of mass number A

ESC08c + MPP + TBA
repulsive attraction

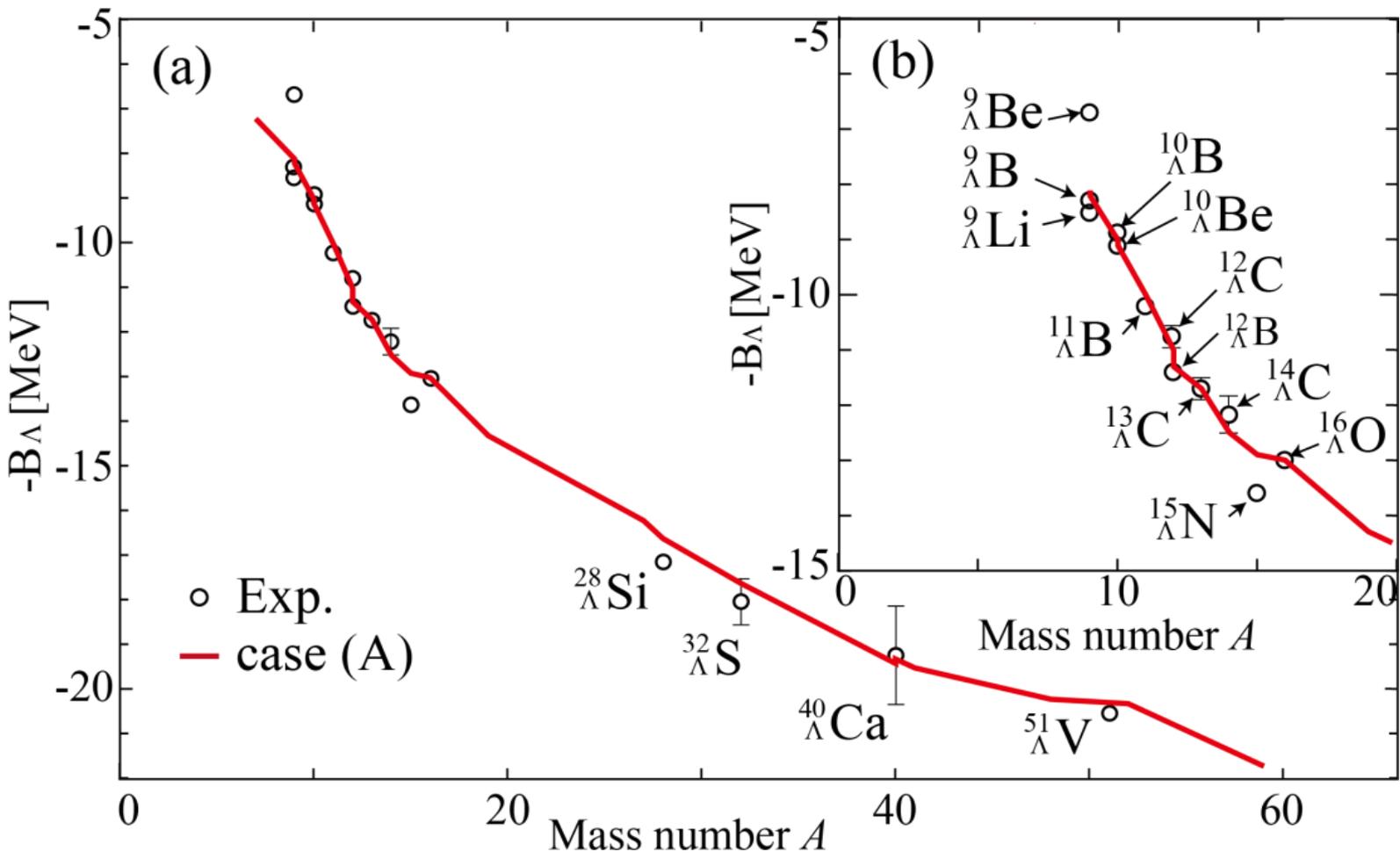


| | β | γ | $\langle\rho\rangle$ | k_F | $-B_\Lambda^{\text{calc}}$ | $-B_\Lambda^{\text{exp}}$ |
|--------------------------|---------|------------|----------------------|-------|----------------------------|---------------------------|
| $^9_\Lambda\text{Li}$ | 0.50 | 2° | 0.072 | 1.02 | -8.1 | -8.50 ± 0.12 [2] |
| $^9_\Lambda\text{Be}$ | 0.87 | 1° | 0.060 | 0.96 | -8.1 | -6.71 ± 0.04 [2] |
| $^9_\Lambda\text{B}$ | 0.45 | 2° | 0.072 | 1.02 | -8.2 | -8.29 ± 0.18 [2] |
| $^{10}_\Lambda\text{Be}$ | 0.57 | 1° | 0.077 | 1.04 | -9.0 | -9.11 ± 0.22 [2] |
| | | | | | | -8.55 ± 0.18 [3] |
| $^{10}_\Lambda\text{B}$ | 0.58 | 1° | 0.075 | 1.04 | -9.1 | -8.89 ± 0.12 [2] |
| $^{11}_\Lambda\text{B}$ | 0.50 | 29° | 0.081 | 1.06 | -10.0 | -10.24 ± 0.05 [1] |
| $^{12}_\Lambda\text{B}$ | 0.39 | 48° | 0.083 | 1.07 | -11.3 | -11.37 ± 0.06 [2] |
| | | | | | | -11.38 ± 0.02 [1] |
| $^{12}_\Lambda\text{C}$ | 0.41 | 34° | 0.086 | 1.08 | -11.0 | -10.76 ± 0.19 [1] |
| $^{13}_\Lambda\text{C}$ | 0.45 | 60° | 0.090 | 1.10 | -11.7 | -11.69 ± 0.19 [1] |
| $^{14}_\Lambda\text{C}$ | 0.45 | 31° | 0.093 | 1.11 | -12.5 | -12.17 ± 0.33 [1] |
| $^{15}_\Lambda\text{N}$ | 0.28 | 60° | 0.098 | 1.13 | -12.9 | -13.59 ± 0.15 [1] |
| $^{16}_\Lambda\text{O}$ | 0.02 | - | 0.105 | 1.16 | -13.0 | -12.96 ± 0.05 [2] |
| $^{19}_\Lambda\text{O}$ | 0.30 | 3° | 0.110 | 1.18 | -14.3 | - |
| $^{27}_\Lambda\text{Mg}$ | 0.36 | 36° | 0.125 | 1.23 | -16.2 | - |
| $^{28}_\Lambda\text{Si}$ | 0.32 | 53° | 0.125 | 1.23 | -16.6 | -17.1 ± 0.2 [9] |
| $^{32}_\Lambda\text{S}$ | 0.28 | 0° | 0.130 | 1.24 | -17.6 | -18.0 ± 0.5 [23] |
| $^{40}_\Lambda\text{K}$ | 0.01 | - | 0.136 | 1.26 | -19.4 | - |
| $^{40}_\Lambda\text{Ca}$ | 0.03 | - | 0.136 | 1.26 | -19.3 | -19.24 ± 1.1 [2] |
| $^{41}_\Lambda\text{Ca}$ | 0.13 | 12° | 0.136 | 1.26 | -19.5 | - |
| $^{48}_\Lambda\text{K}$ | 0.01 | - | 0.141 | 1.28 | -20.2 | - |
| $^{51}_\Lambda\text{V}$ | 0.18 | 2° | 0.151 | 1.31 | -20.3 | -20.51 ± 0.13 [2] |
| $^{59}_\Lambda\text{Fe}$ | 0.26 | 23° | 0.142 | 1.28 | -21.7 | - |

B_Λ as a function of mass number A

ESC08c + MPP + TBA
 repulsive attraction

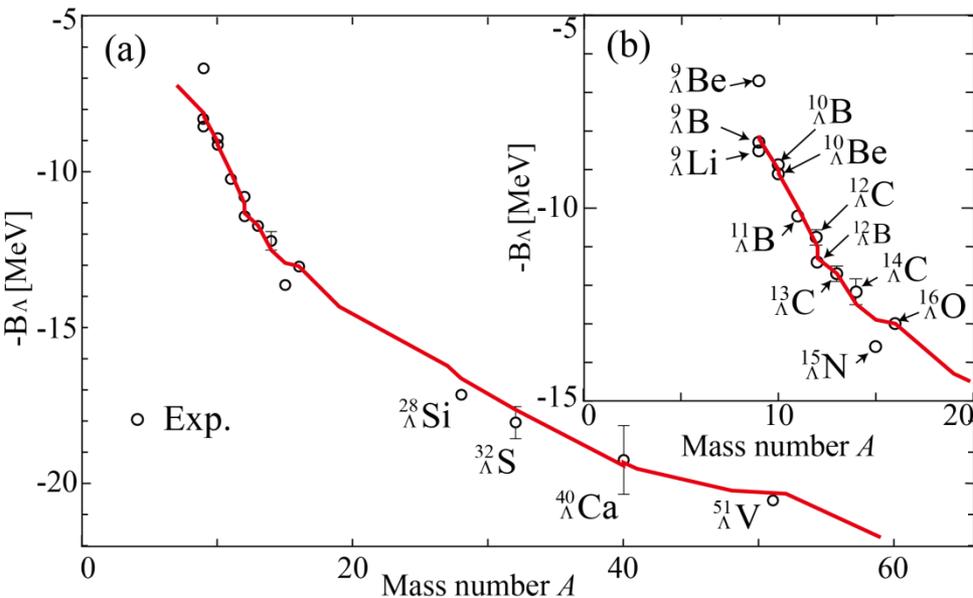
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| | | | | | | -11.37 ± 0.06 [2] |
| | | | | | | -11.38 ± 0.02 [1] |
| | | | | | | -10.76 ± 0.19 [1] |
| | | | | | | -11.69 ± 0.19 [1] |
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HyperAMD w/ ESC08c + MPP + TBA successfully reproduces B_Λ

B_Λ as a function of mass number A

ESC08c + MPP + TBA
 repulsive attraction



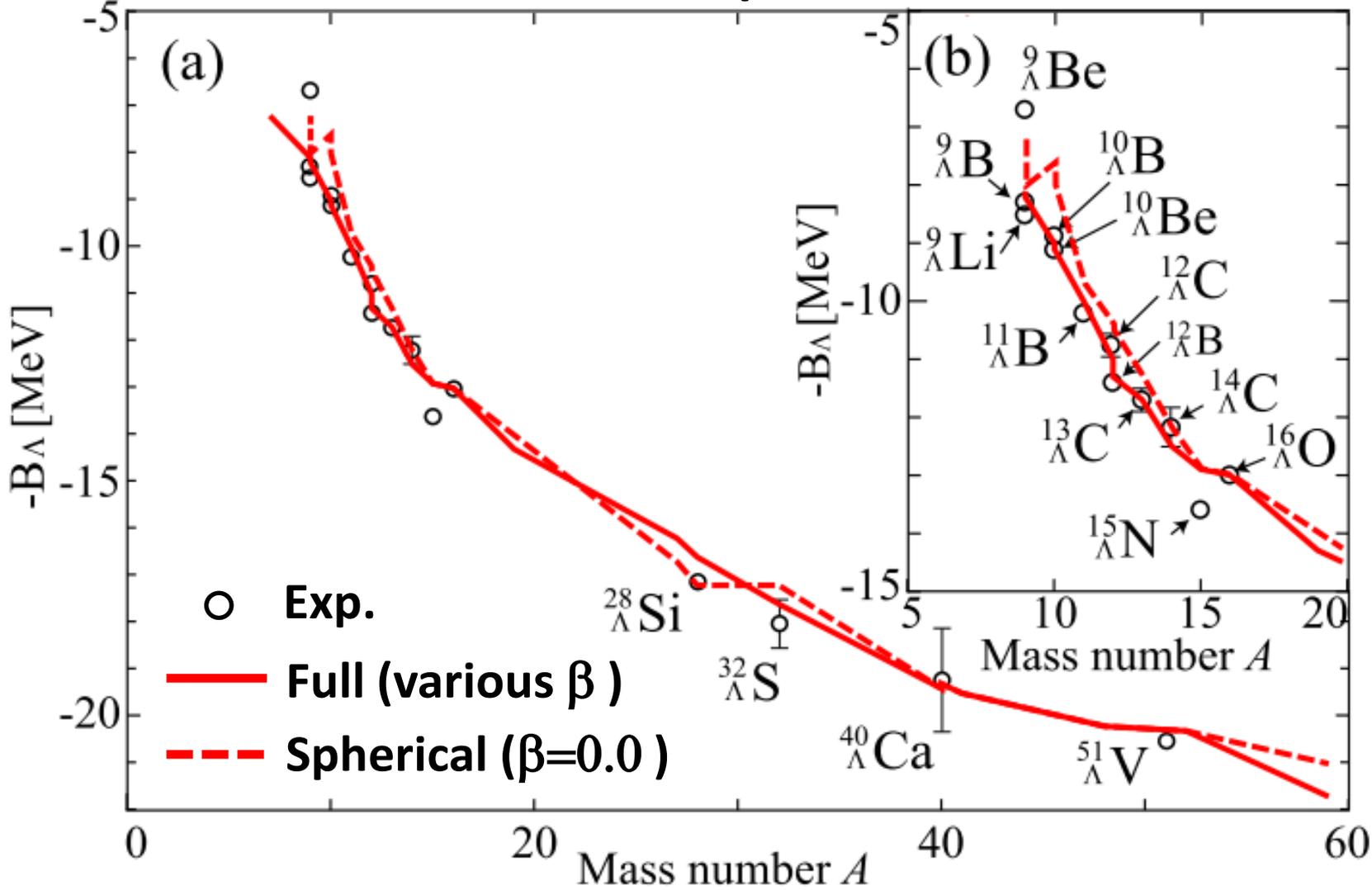
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| $^{59}_\Lambda\text{Fe}$ | 0.26 | 23° | 0.142 | 1.28 | -21.7 | - |

HyperAMD w/ ESC08c + MPP + TBA successfully reproduces B_Λ

What is essential to reproduce B_{Λ}

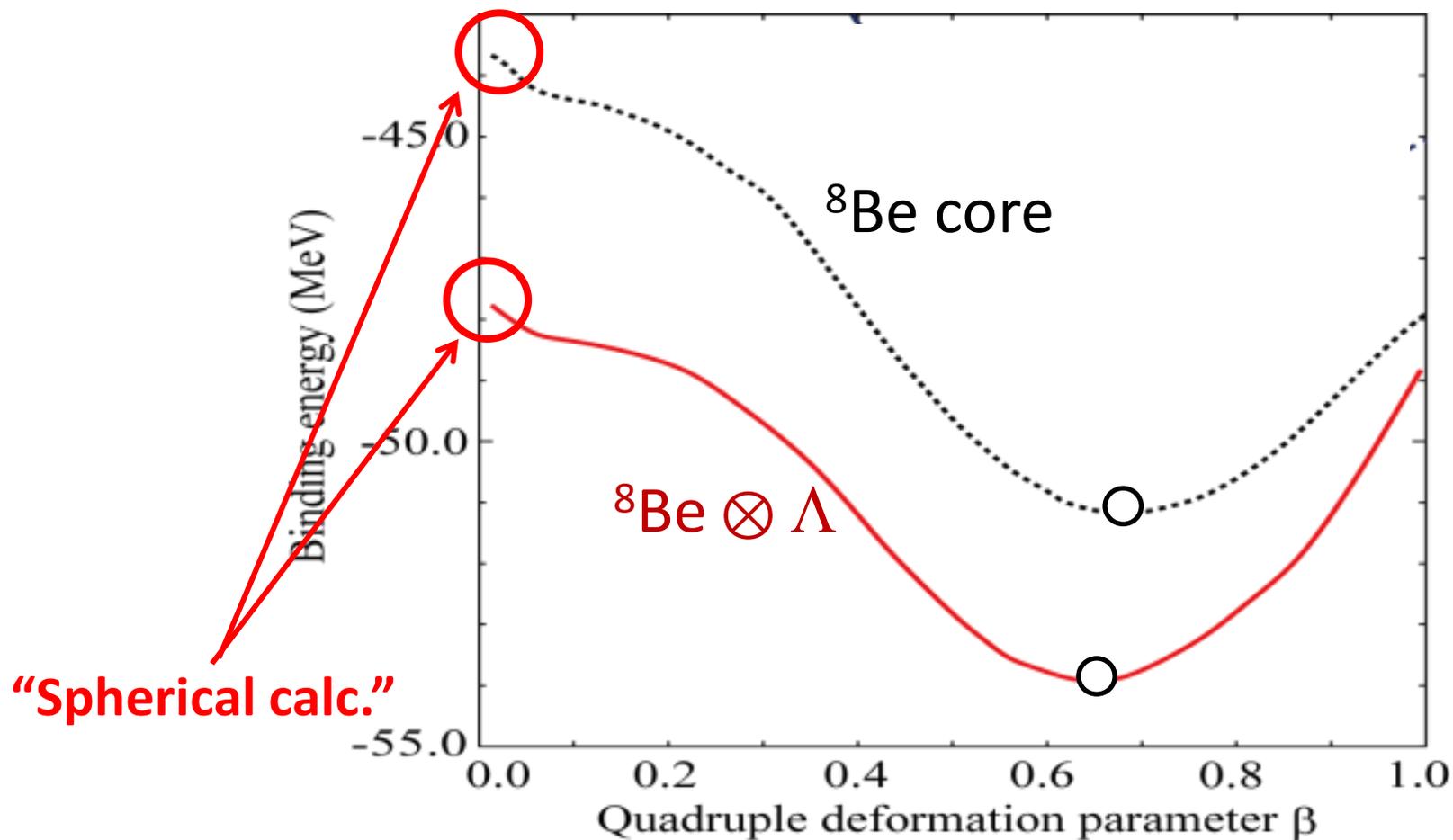
“Description of the core structure”

“Full calc.” vs. “Spherical calc.”



What is essential to reproduce B_Λ

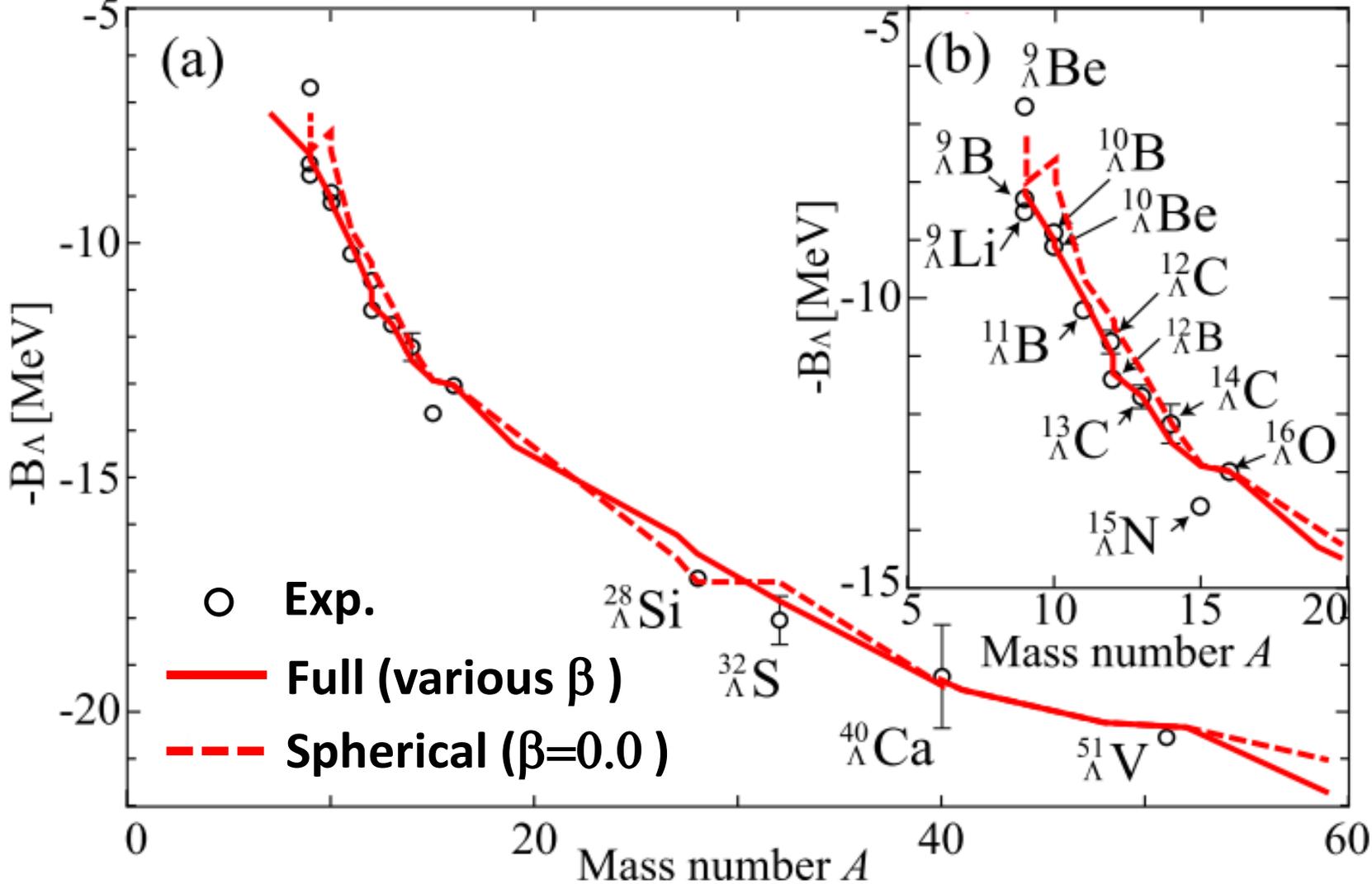
Ex. ${}^9_\Lambda\text{Be}$ “Full calc.” vs. “Spherical calc.”



“Full calc.”: all of w.f. on energy curve in GCM calc.

What is essential to reproduce B_Λ

“Description of the core structure”



Deformation of the ground states is essential to reproduce B_Λ

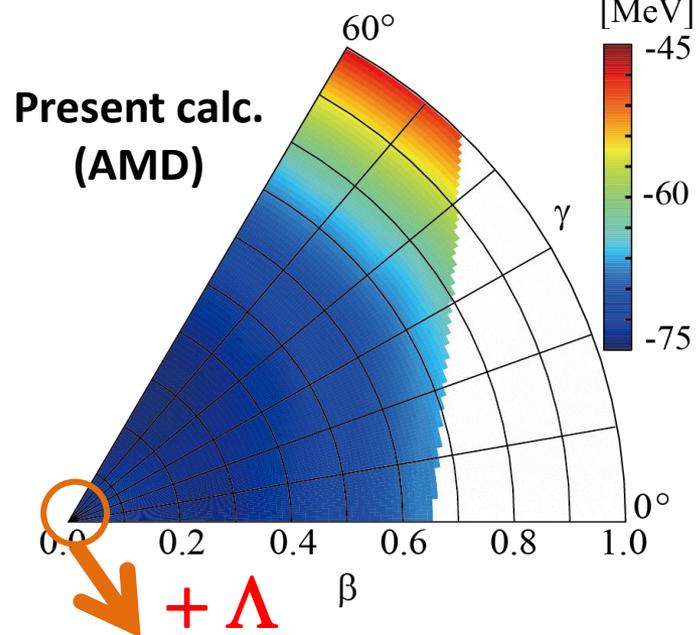
What is essential to reproduce B_Λ

“Description of the core structure”

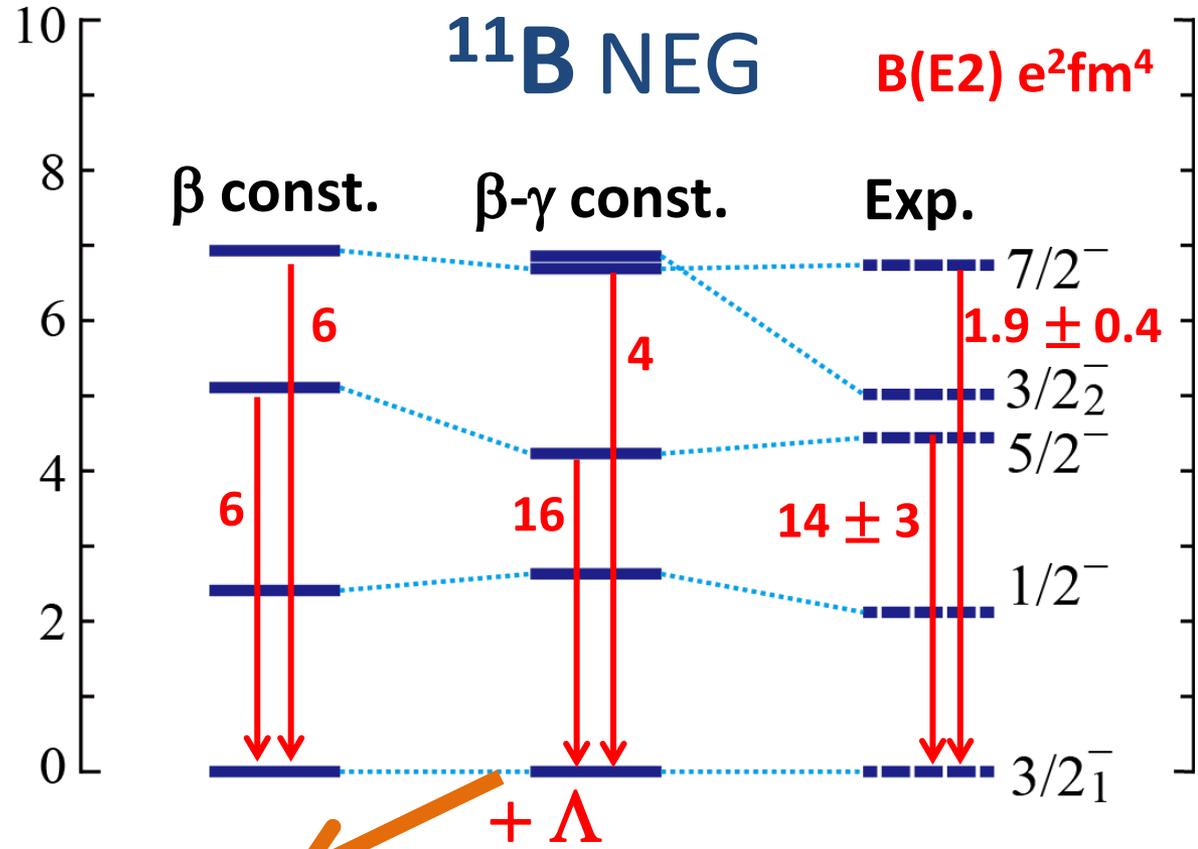
Ex.: ^{11}B

More sophisticated treatment: GCM calc. on (β, γ) plane

T. Suhara and Y. Kanada-En'yo,
PTP123,303(2010)



$^{12}_\Lambda\text{B}$ (Spherical)
 $B_\Lambda = 9.5 \text{ MeV}$
 $(k_F = 1.16 \text{ fm}^{-1})$

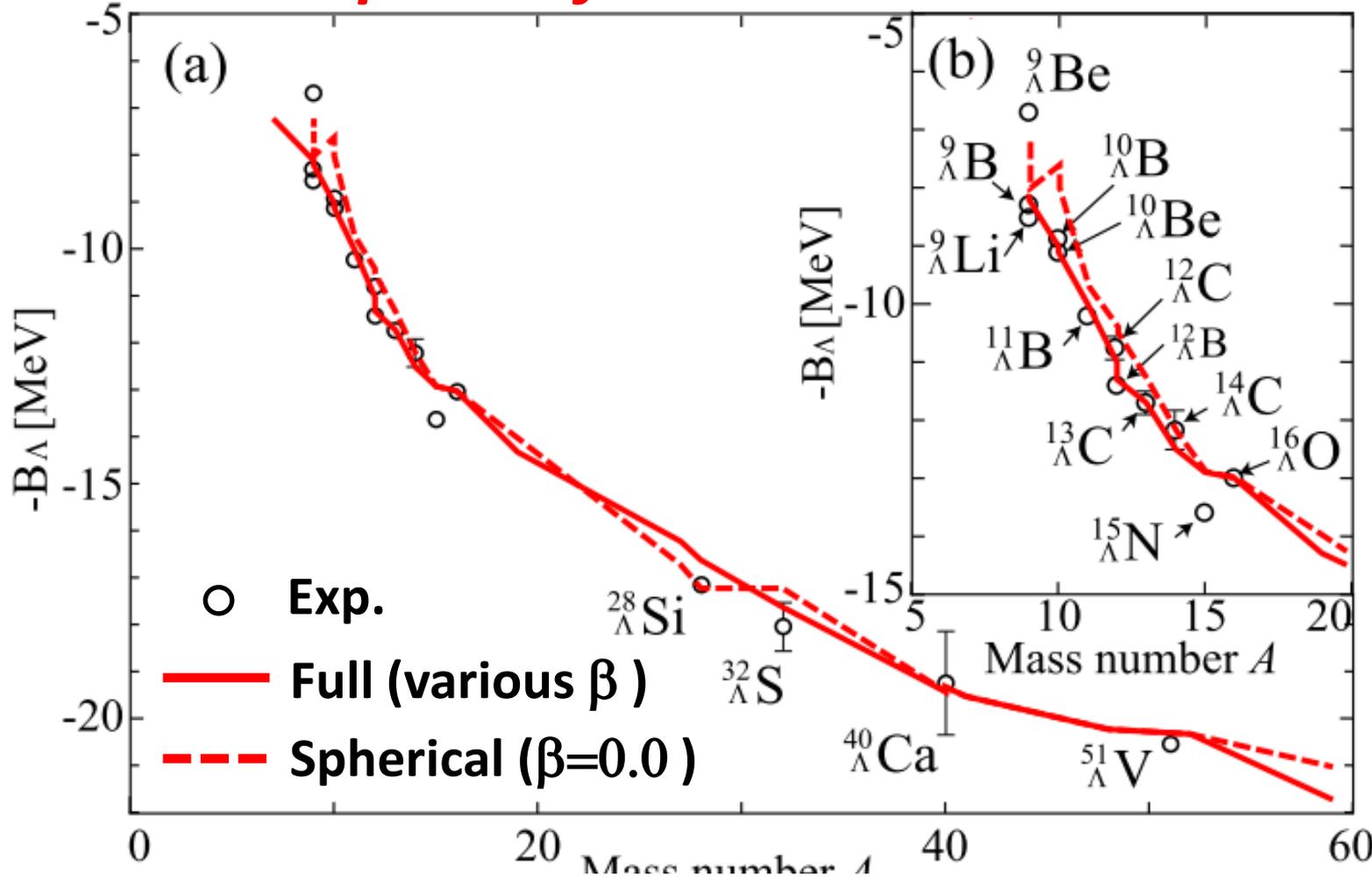


$^{12}_\Lambda\text{B}$ (β - γ)
 $B_\Lambda = 11.3 \text{ MeV}$
 $(k_F = 1.07 \text{ fm}^{-1})$

$^{12}_\Lambda\text{B}$ (EXP)
 $B_\Lambda = 11.4 \pm 0.02 \text{ MeV}$

What is essential to reproduce B_Λ

“Description of the core structure”



Deformations affect A dep. of B_Λ through k_f dep. of interaction and overlap between Λ and nucleons

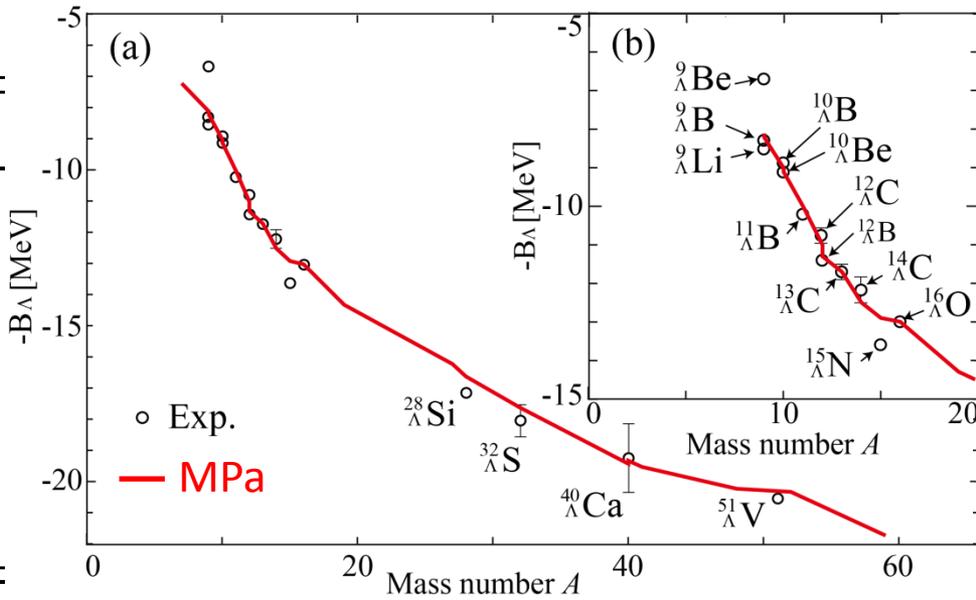
Comparison with the results with ESC08c only

◆ Effects of many-body force

ESC: ESC08c only
 MPa: ESC08c + MPP + TBA

Over-binding with ESC08c only

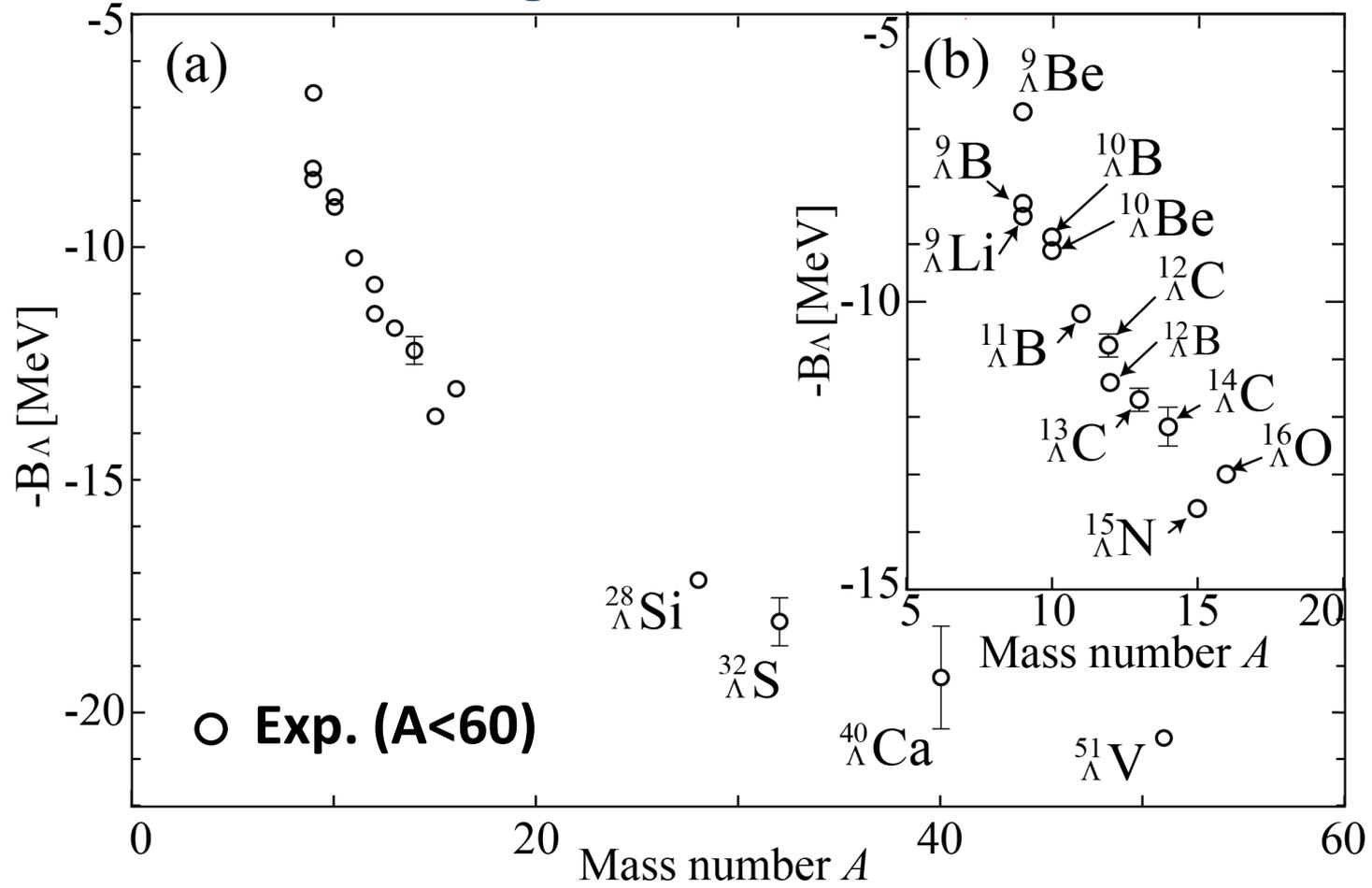
| [MeV] | ESC | MPa | Exp. |
|----------------------------|-------|-------|-------------------|
| $^{13}_{\Lambda}\text{C}$ | -11.5 | -11.7 | -11.69 ± 0.19 |
| $^{16}_{\Lambda}\text{O}$ | -13.3 | -13.0 | -12.96 ± 0.05 |
| $^{28}_{\Lambda}\text{Si}$ | -17.7 | -16.6 | -17.1 ± 0.2 |
| $^{40}_{\Lambda}\text{K}$ | -21.5 | -19.4 | — |
| $^{48}_{\Lambda}\text{K}$ | -22.6 | -20.2 | — |
| $^{51}_{\Lambda}\text{V}$ | -23.5 | -20.3 | -20.51 ± 0.13 |



- Observed B_{Λ} is well reproduced with MPP + phenomenological TBA
- Systematic data of B_{Λ} will provide a new insight to many-body force

Current status of observed B_{Λ}

Observations are not enough with $A > 16$



Systematic and accurate data of observed B_{Λ} are desired

Bertini *et al.*, NPA83,306(1979), Davis, Juric, *et al.*, NPB52(1973), Davis, NPA547,369(1992);NPA754,3c(2005), Ajimura *et al.*, NPA639(1998)93c, Pile *et al.*, PRL66,2585(1991), Hotchi *et al.*, PRC64, 044302(2001), Hashimoto and Tamura, PPNP57,564(2006), Tang, *et al.*, PRC90,034320(2014).

Summary

◆ Summary

- HyperAMD + GCM was applied with ESC08c + MPP + TBA interaction

Observed B_{Λ} are successfully reproduced in wide mass regions

● Structure of the core nuclei

- Spherical shape: deviate from observed B_{Λ}
- Description of core deformation is essential

**Density dependence
of ΛN interaction**

● Many-body (MPP + TBA) force effects

- Input: experimental data of ${}^{89}_{\Lambda}\text{Y}$
- MPP + TBA force could be determined by systematic data of B_{Λ}

◆ Future plan

- To reveal reasons for deviation of B_{Λ} with $A < 9$ (e.g. ${}^9_{\Lambda}\text{Be}$)
- Many-body force effects: from systematics of B_{Λ} , Λ in excited orbit?