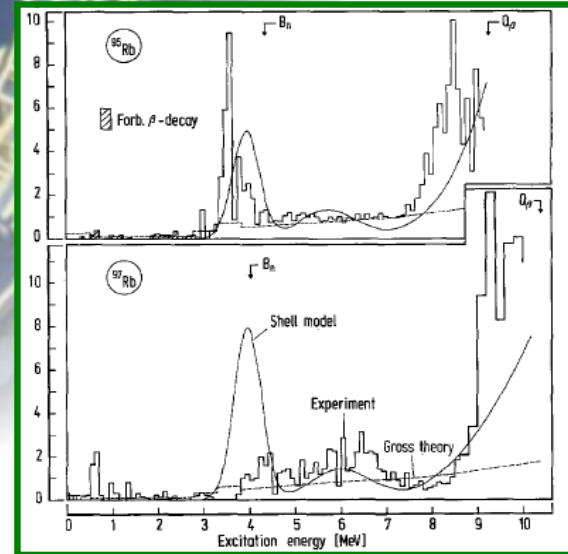
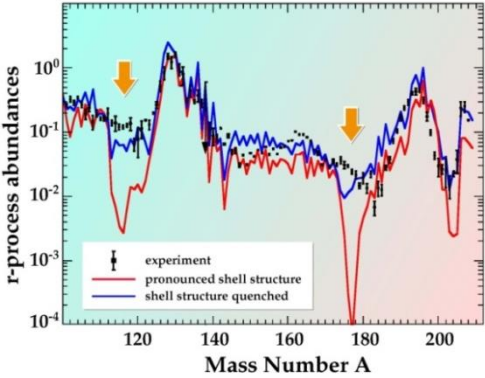


Thirty years of collaboration with Peter Möller: From RPA (1984) to FRDM & QRPA (today)



ASRC Workshop (JAEA)
Tokai, Dec. 2014

Karl-Ludwig Kratz

The starting-point: my PhD-thesis

Because of the important contribution of **halogen** isotopes to the **β -delayed neutron** groups occurring in fission, rapid separation procedures for **Br** and **I** were already developed in the 1950's by several groups.



In those days, chemical **ELEMENT** separation was (still) superior to physical **MASS** separation.

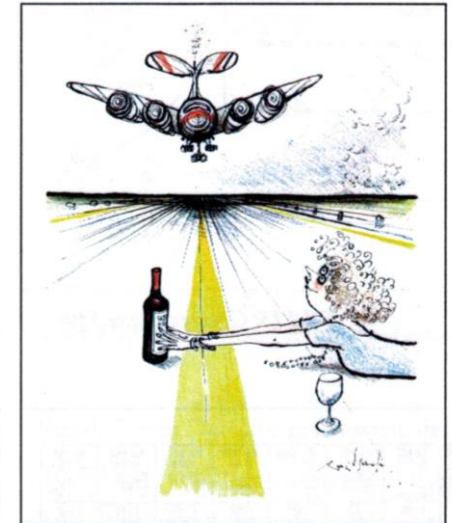
“SRAFAP”
(Students Running
As Fast As Possible)

In my **PhD thesis**, I worked out an **ultra-fast separation procedure** based upon formation of methyl-halides in a hot-atom reaction between fission halogens recoiling from a solid U-235 target into gaseous methane:



... shortest-lived isotopes detected :

0.65 s ^{91}Br , 0.25 s ^{92}Br and 0.45 s ^{141}I , 0.2 s ^{142}I



“HITSEP”
(High-Tech **SE**paration
Procedure)

Development of β dn-spectroscopy, as Mainz postdoc

Neutron detector:

He-3 ionisation chambers

Energy resolution:

10 keV for $E_n = \text{thermal peak}$

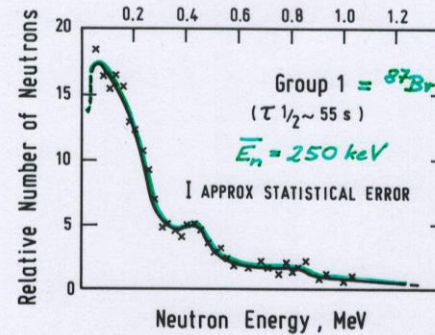
20 keV for $E_n = 2 \text{ MeV}$

Efficiency:

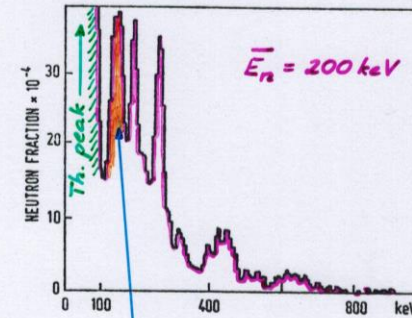
$10^{-4} - 10^{-5}$

...new spectroscopic method ?

Importance of low-energy β dn's for astrophysics

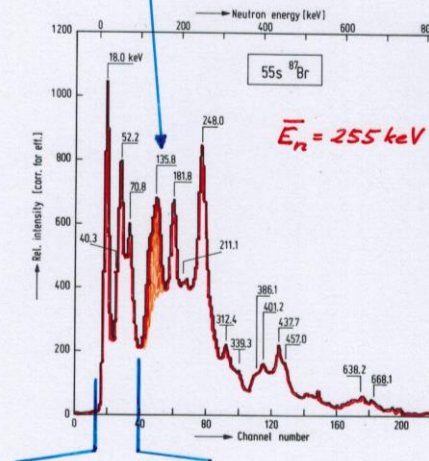


Batchelor et al.
1956



Shalev - Redstam
1974

^{87}Br



Mainz - Berkeley
Kollaboration
1976

$E_n \leq 100 \text{ keV}$ ca. 30% der DN von Grp. 1 !

β dn-emission as inverse process to n-capture

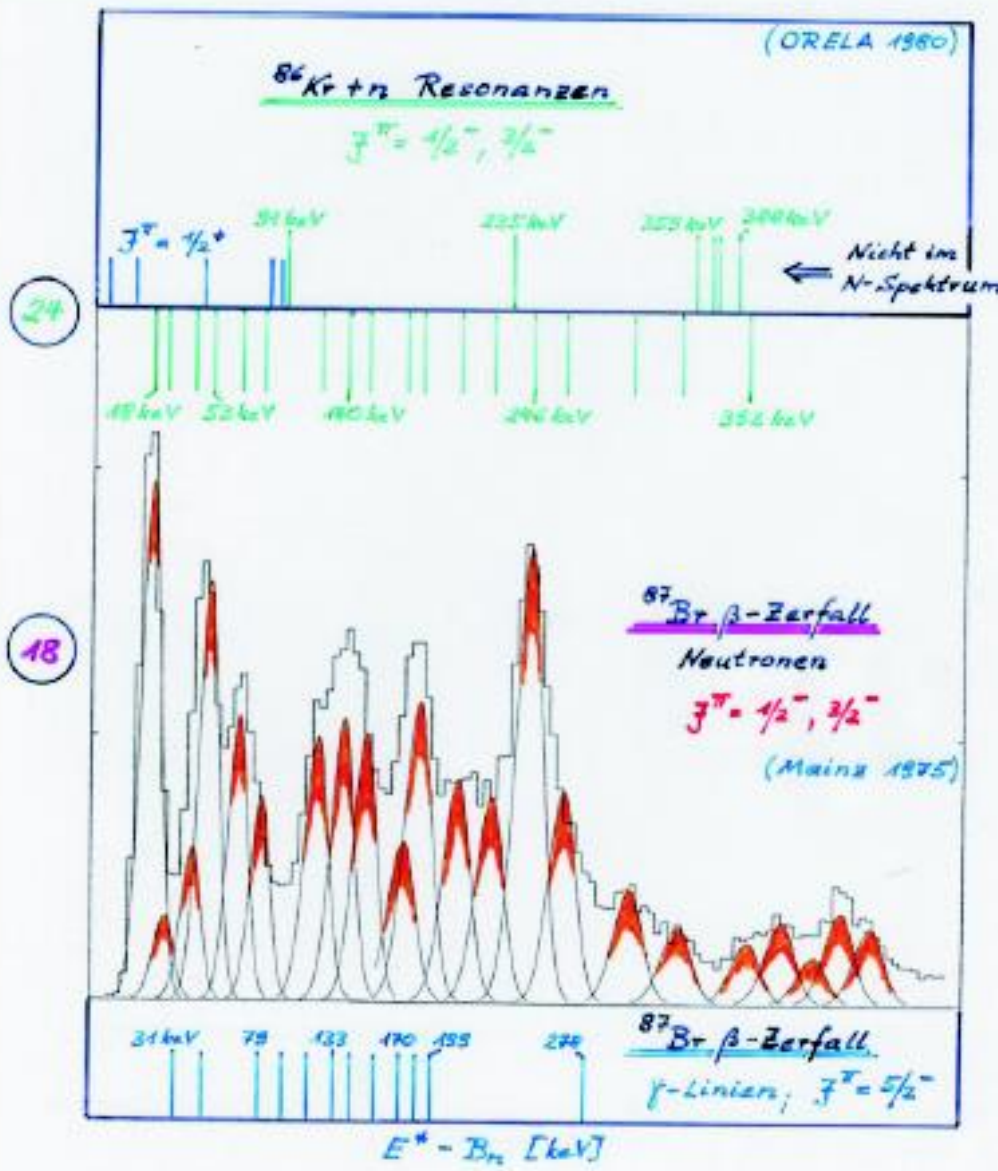
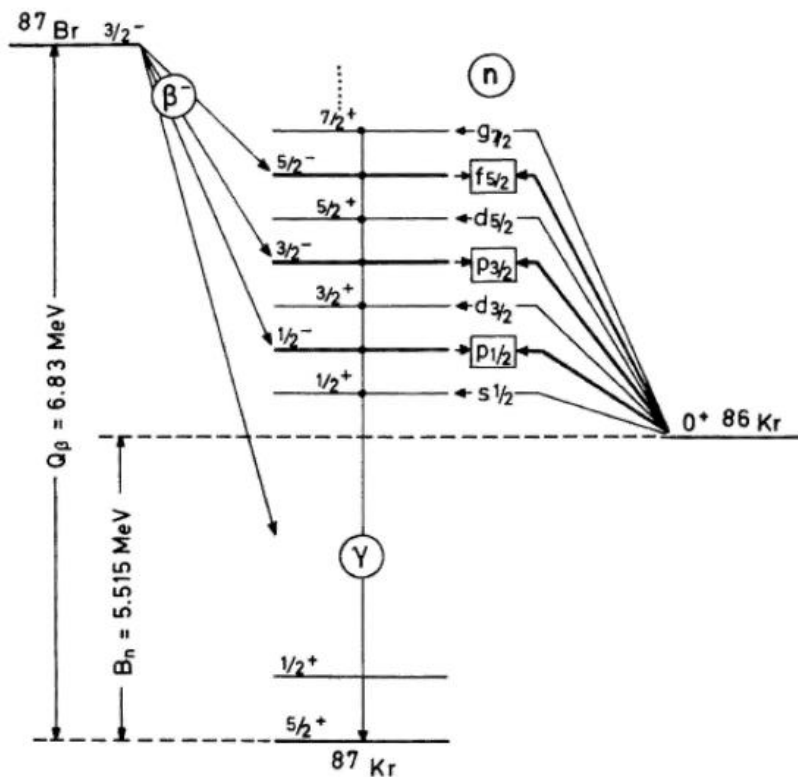
Two experiments:



at ORELA (Oak Ridge)



at TRIGA Mainz



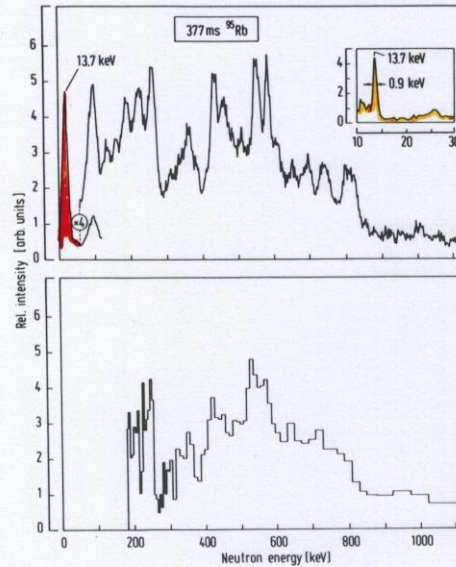
The early 1980s: the "OSTIS Period" at ILL

Detailed nuclear structure studies of neutron-rich Rb isotopes

$(Q_\beta - B_n)$ in ^{95}Rb large enough to populate many excited states in the final nucleus ^{94}Sr

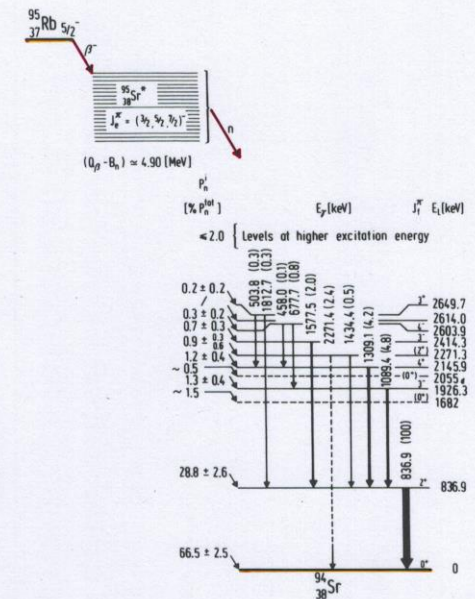
13.7 keV peak in βdn ground-state spectrum
 \Rightarrow **p-wave neutron**

$\beta\gamma$ - spectroscopy
 $\gamma\gamma$ - spectroscopy
 including angular correl.
 conv. electrons
 $\tau_{\text{level}}, Q_\beta$
 del. n. - spectroscopy
 including $n\gamma$ - coinc.
 E_n - singles
 $E_n\gamma$ - coinc.

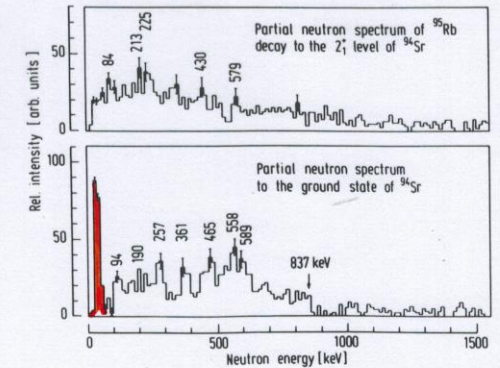


E_n - singles spectra with ^3He ion. - chambers
 TOF-techniques $\left\{ \begin{array}{l} \text{Li-glass} \\ \text{NE 213} \end{array} \right.$

here: 377ms $^{95}\text{Rb}_{58}$ d.n. - decay



$\beta(^{95}\text{Rb}) - n(^{95}\text{Sr}) - \gamma(^{94}\text{Sr})$ coinc.
 with a high-eff. ^3He long-counter



$E_n\gamma$ - coinc. spectra (^3He ion.- ch.)

β dn-Data \Rightarrow high-energy part of $S_\beta(E)$

Z. Phys. A312, 43 (1983)

The Beta-Decay of ^{95}Rb and ^{97}Rb

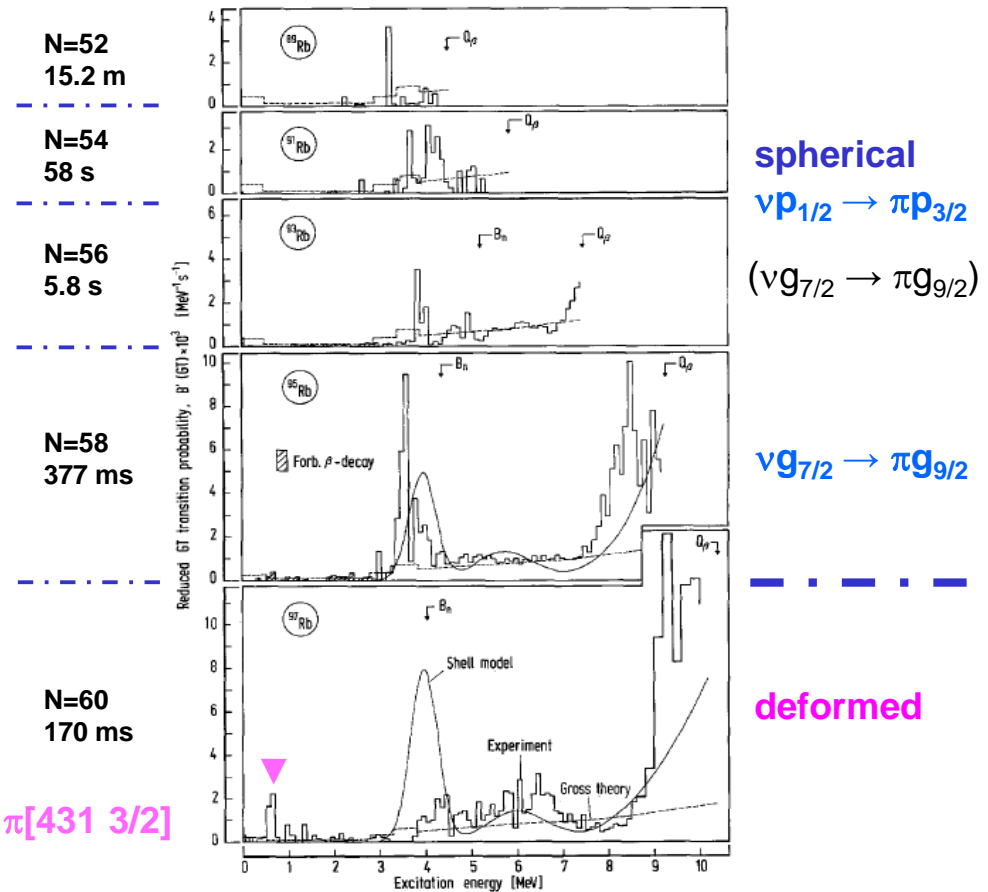
K.-L. Kratz, H. Ohm, A. Schröder et al.

Combination of γ - and β dn-data \Rightarrow exp. $S_\beta(E)$



The high-energy part of $S_\beta(E > B_n)$ is given by the β dn-energy spectra and $P_n(i)$ -values

$v[420 1/2] \Rightarrow \pi[431 3/2]$



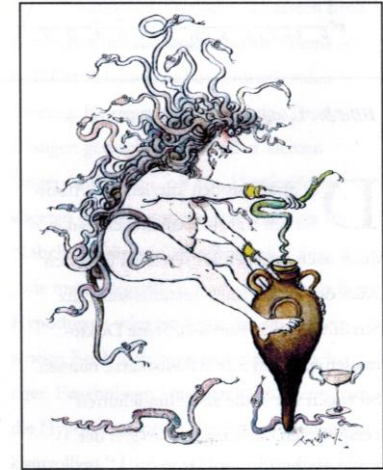
$S_{\beta}(E)$ of Rb isotopes: the starting point of my collaboration with Peter

Nucl. Phys. A417 (1984)

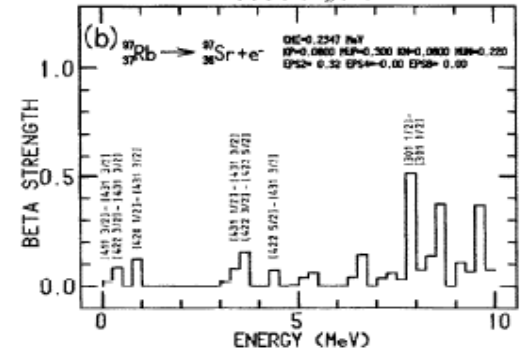
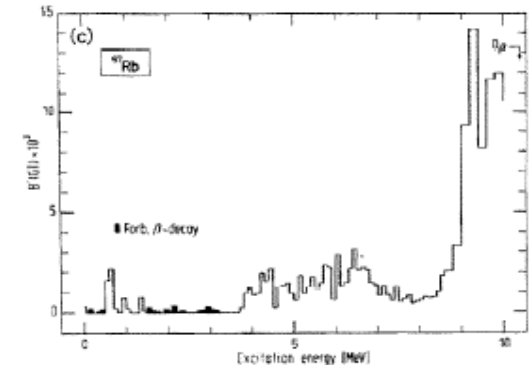
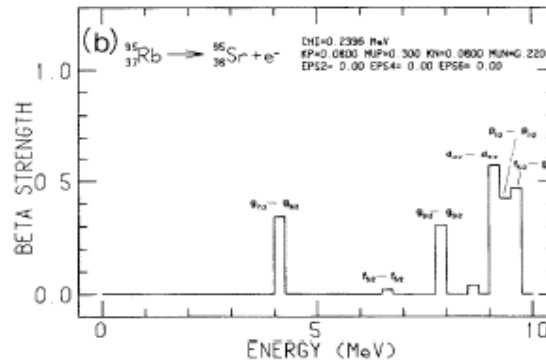
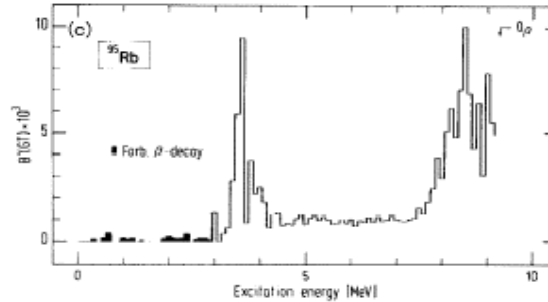
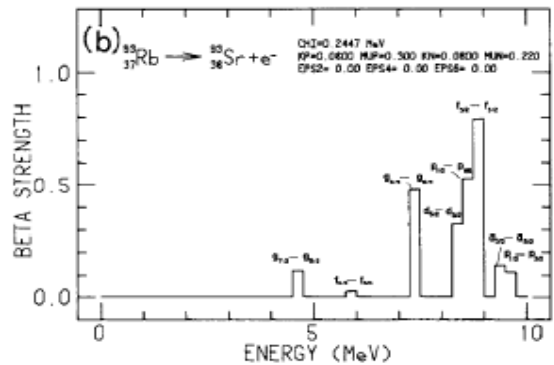
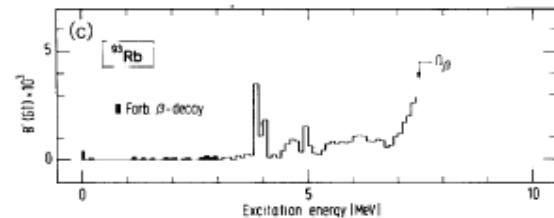
CALCULATION OF GAMOW-TELLER β -STRENGTH FUNCTIONS IN THE RUBIDIUM REGION IN THE RPA APPROXIMATION WITH NILSSON-MODEL WAVE FUNCTIONS

JOACHIM KRUMLINDE and PETER MÖLLER

Departments of Physics and Mathematical Physics, Lund University, Box 725, S22007 Lund, Sweden



The “antique” $S_{\beta}(E)$ paper

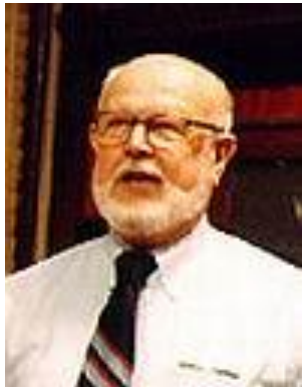


“We find good agreement between calculated and experimental spectra, provided an appropriate choice of single-particle parameters and deformation is made.”

How we became interested in “astro” and “cosmo”

In 1981, I got an invitation to a USA – EU workshop on “**Nuclear Astrophysics**”, to give a review talk about “**Beta-decay detection methods and limits**”

At this workshop, Willy Fowler gave a talk on isotopic **FUN** anomalies of Ca and Ti found in the EK-1-4-1 inclusion of the Allende meteorite
...in particular $^{48}\text{Ca}/^{46}\text{Ca} = 250 !$



“... Agreement for the ^{46}Ca and ^{49}Ti anomalies was obtained (within the assumed “**n β ”** nucleosynthesis process) by increasing the theoretical Hauser-Feshbach cross sections for $^{46}\text{K}(n,\gamma)$ and $^{49}\text{Ca}(n,\gamma)$ by a factor 10 on the basis of probable thermal (30 keV, s-wave) resonances [...] in the compound nuclei ^{47}K and ^{50}Ca , respectively. ...”

Already during this talk I wondered,
if we had already measured this resonance at **CERN / ISOLDE** via high-resolution **β -delayed neutron spectroscopy** of
 $^{50}\text{K}(\beta^-)^{50*}\text{Ca}(n)^{49}\text{Ca}$
as “inverse process” to neutron capture of
8.7 min $^{49}\text{Ca}(n,\gamma)^{50}\text{Ca}$

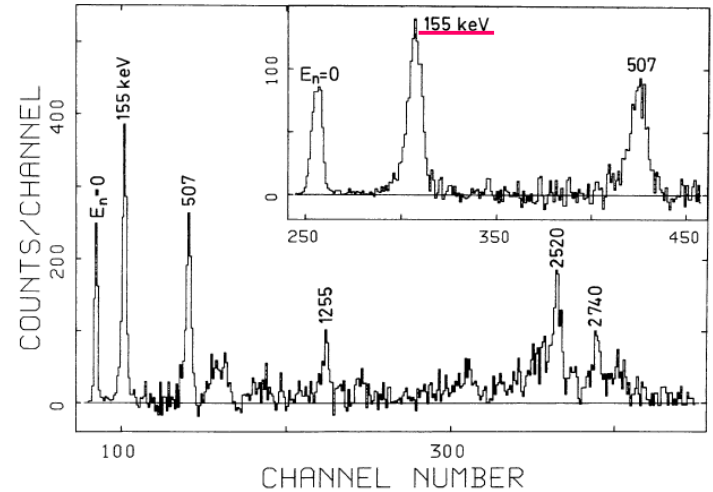
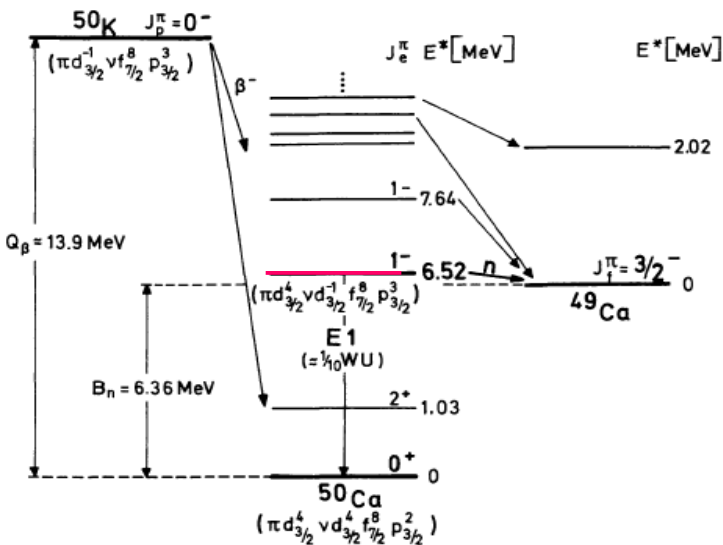
This led to the EK-1-4-1 story (I)

Astron. Astrophys. 125 (1983)

Determination of stellar neutron-capture rates for radioactive nuclei with the aid of β -delayed neutron emission

K.-L. Kratz, W. Ziegert, W. Hillebrandt & F.-K. Thielemann

we discuss the reaction $^{49}\text{Ca}(n, \gamma) ^{50}\text{Ca}$ [$T_{1/2}(^{49}\text{Ca})=8.7$ min] with its inverse process [$^{50}\text{K}(\beta^-) ^{50}\text{Ca}^*(n)$ ^{49}Ca [$T_{1/2}(^{50}\text{K})=740$ ms]. The observed non-statistical behaviour of this reaction supports recent attempts to explain the isotopic anomalies in Ti observed in meteoritic inclusions.



Low-lying s-wave resonance in ^{50}Ca does exist; however, not at 30 keV, but at 155 keV.
 $n\beta$ -process \Rightarrow r-process

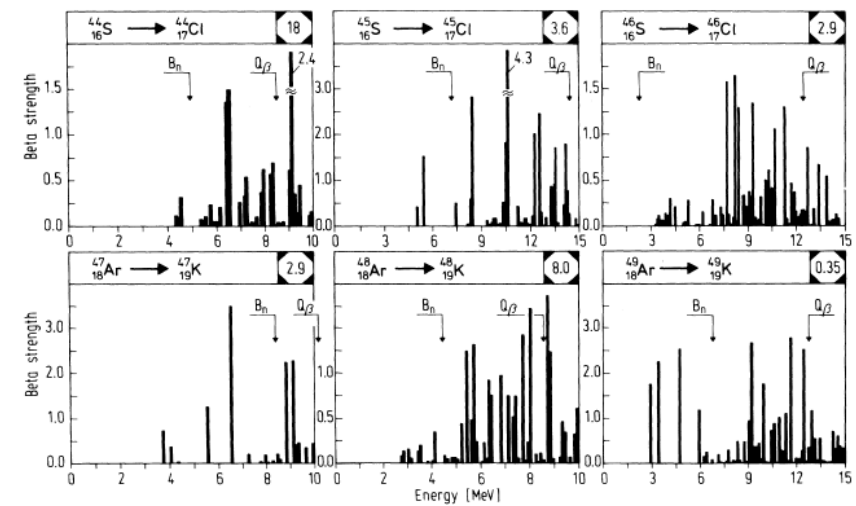
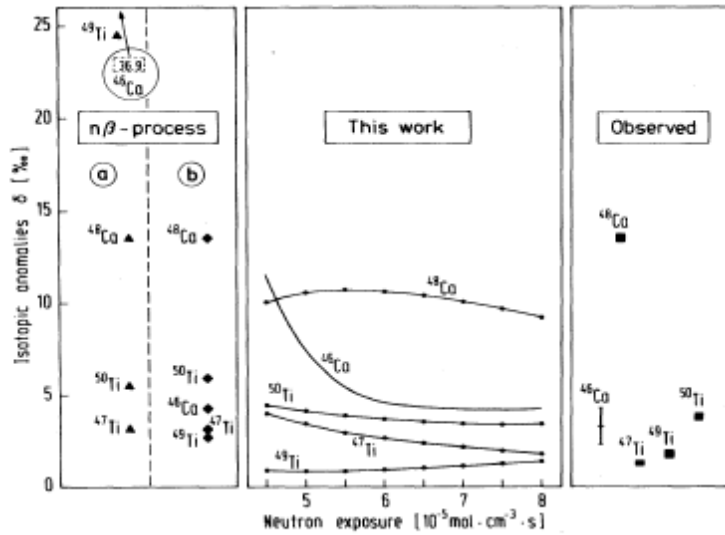
Fig. 4. Principal characteristics of β -delayed neutron decay of 740 ms ^{50}K . The neutron-emitter state at 6.52 MeV ($J^\pi=1^-$) in ^{50}Ca may have the properties of the s-wave neutron-capture state required by Sandler et al. (1982) (see text)

The EK-1-4-1 story (II)

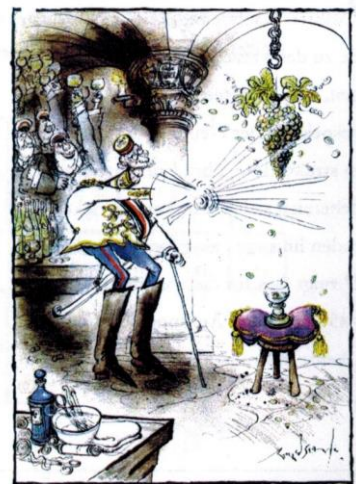
Phys. Rev. Lett. 55 (1985)

Interpretation of the Solar $^{48}\text{Ca}/^{46}\text{Ca}$ Abundance Ratio and the Correlated Ca-Ti Isotopic Anomalies in the EK-1-4-1 Inclusion of the Allende Meteorite

W. Ziegert, M. Wiescher, K.-L. Kratz
P. Möller, J. Krumlinde
 F.-K. Thielemann, W. Hillebrandt



GT strength functions from RPA



Cleverness required

Left: Ca-Ti anomalies predicted by $n\beta$ -process;
 (a) HF rates for all nuclei
 (b) HF x 10 for ^{46}K and $^{49}\text{Ca}(n,\gamma)$.
 Middle: **Our astro-calculations.**
 Right: **EK-1-4-1 observations.**

The EK-1-4-1 story (III) – 15 years later



MEMORIE DELLA
SOCIETÀ ASTRONOMICA ITALIANA
 JOURNAL OF THE ITALIAN ASTRONOMICAL SOCIETY

ON THE ORIGIN OF THE CA-TI-CR ISOTOPIC ANOMALIES
 IN THE INCLUSION EK-1-4-1 OF THE
 ALLENDE-METEORITE

K.-L. KRATZ¹, W. BÖHMER¹, C. FREIBURGHANUS², P. MÖLLER³, B. PFEIFFER¹,
 T. RAUSCHER², F.-K. THIELEMANN²

Updates:

Experiments at **CERN / ISOLDE** & **GANIL / LISE**
 Theoretical n-capture rates (CN + DC)
 Astrophysical network calculations

Two astrophysical scenarios:

(i) **α -process**

(ii) **weak r-process**

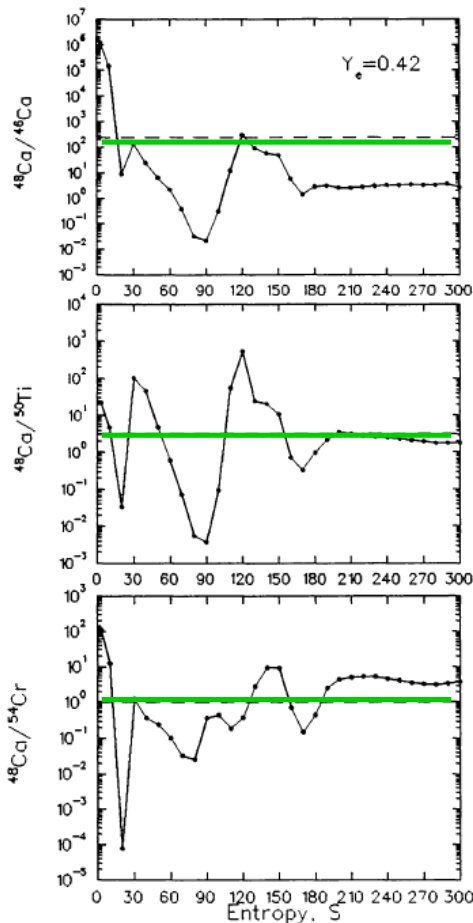
α -rich freezeout from

Si-QSE

Fe-QSE + n-capt.

↻ SN Ia

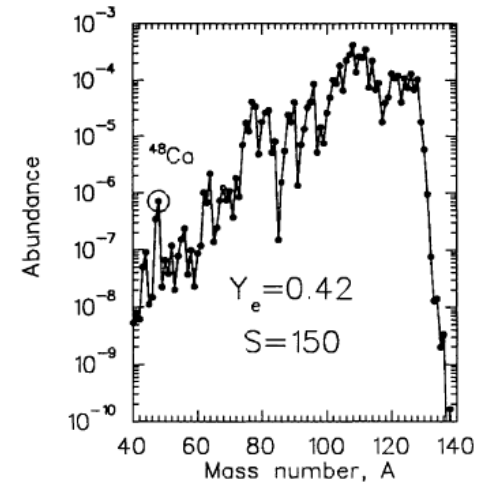
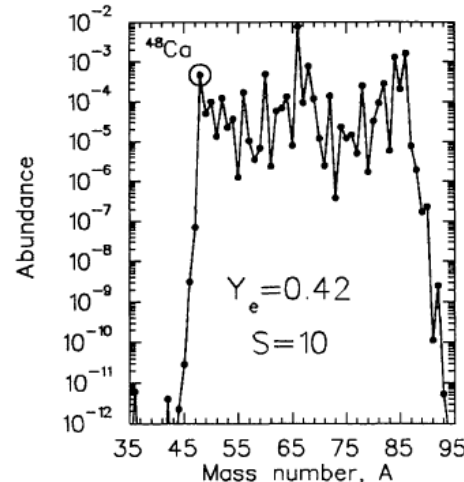
↻ SN II



EK-1-4-1

EK-1-4-1

EK-1-4-1

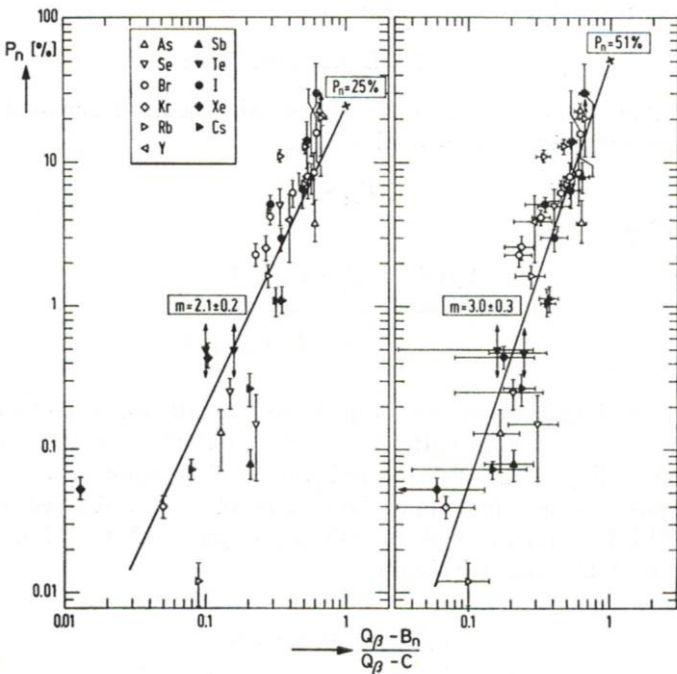


Z. Phys. 263, 435 (1973)

Systematics of Neutron Emission Probabilities from Delayed Neutron Precursors

K.-L. Kratz* and G. Herrmann

Institut für Kernchemie der Universität Mainz



Theoretically,

the two **gross / integral** β -decay quantities, $T_{1/2}$ and P_n , are interrelated via their traditional definition in terms of the so-called **β -strength function** $[S_\beta(E)]$

“Experimental” definition (Duke et al., 1970)

$$S_\beta(E) = \frac{b(E)}{f(Z, Q_\beta - E) \cdot T_{1/2}} \quad [s^{-1} \text{MeV}^{-1}]$$

$b(E)$ absolute β -feeding per MeV,
 $f(Z, Q_\beta - E)$ Fermi function,
 $T_{1/2}$ β -decay half-life.

P_n as ratio of $S_\beta(E) \times f$ above B_n to total $S_\beta(E) \times f$ within Q_β $T_{1/2}$ as reciprocal $S_\beta(E) \times f$

assuming $S_\beta(E) = \text{const.}$; $f \sim (Q_\beta - E)^5$ and cut-off energy C



The “Kratz-Herrmann Formula”

$$P_n = a[(Q_\beta - B_n)/(Q_\beta - C)]^b$$

$$T_{1/2} = c[1/(Q_\beta - C)]^d$$

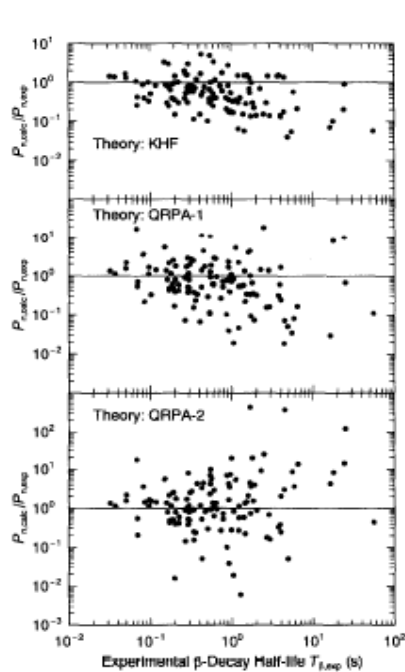
From KHF to MPK (I)

Progr. Nucl. Energy 41, 39 (2002)

STATUS OF DELAYED-NEUTRON PRECURSOR DATA: HALF-LIVES AND NEUTRON EMISSION PROBABILITIES

Bernd Pfeiffer¹ and Karl-Ludwig Kratz
Institut für Kernchemie, Universität Mainz, Germany

Peter Möller
Theoretical Division, Los Alamos National Laboratory, Los Alamos, NM 87545

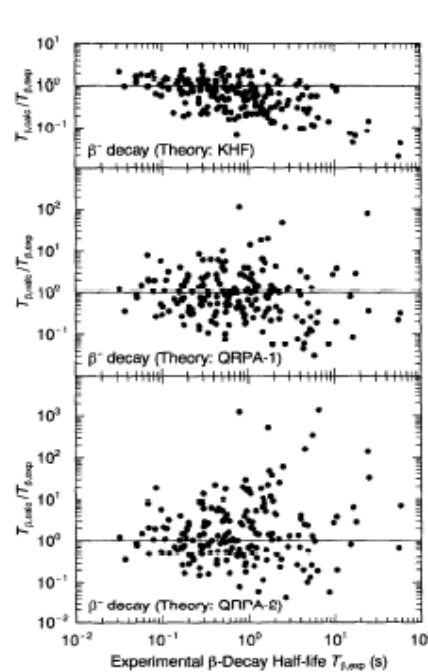


σ_{rms}^2

2.8

4.1

5.6

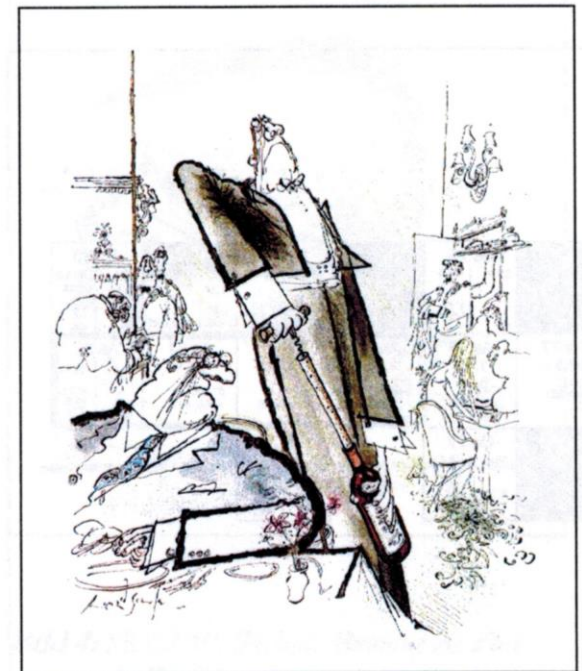
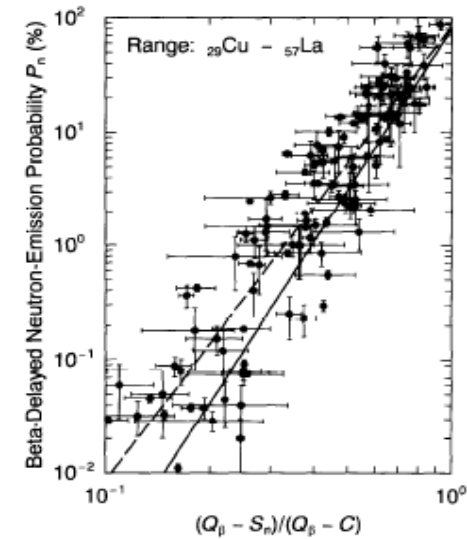


σ_{rms}^2

2.6

3.9

5.8



...and Peter was not amused !

Ratios of calculated to experimental P_n (left) and $T_{1/2}$ (right):
upper part: Kratz - Herrmann **Formula** (KHF)
lower parts: QRPA(GT); Pfeiffer – Kratz – Möller **Model** (MPK)

...a simple **formula** cannot be better than a microscopic **model** !

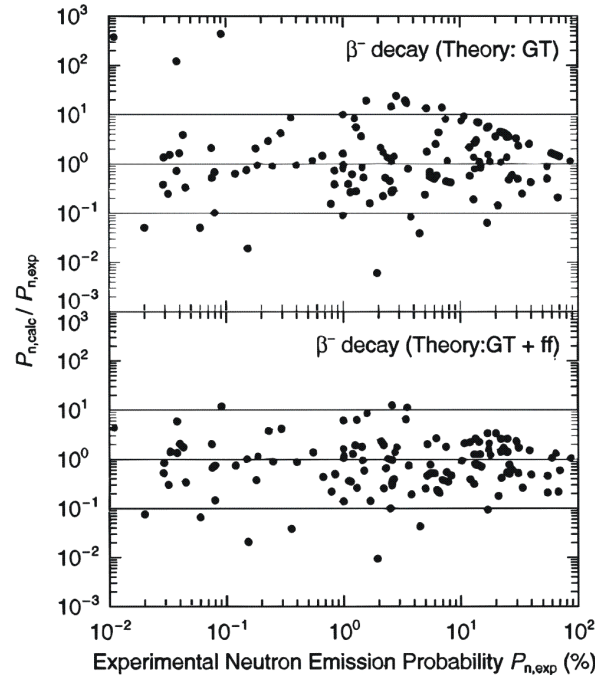
Let's think about possible improvements...



Phys. Rev. C67 (2003)

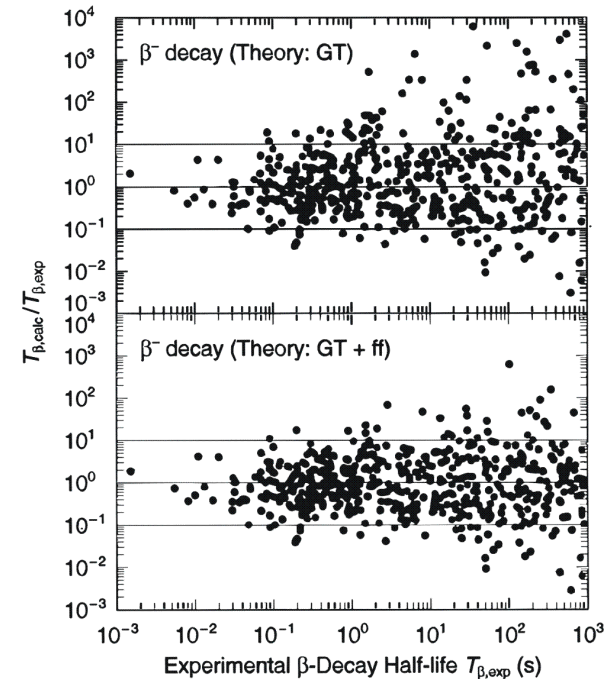
New calculations of gross β -decay properties for astrophysical applications:
Speeding-up the classical r-process

P. Möller, B. Pfeiffer, K.-L. Kratz



Total error 3.5

$T_{1/2}(\text{GT+ff})$ become shorter
 \curvearrowright faster r-process matter flow



Total error 3.1

- add ff-part of Gross Theory to GT-strength
- include empirical spreading of GT-strength

$T_{1/2}$ and P_n calculations in 3 steps:

(1) Mass model FRDM

$\sim Q_\beta, S_n, \varepsilon_2$
Folded-Yukawa wave fcts.

SP shell model QRPA (pure GT)
with input from FRDM
potential: Folded Yukawa
pairing-model: Lipkin-Nogami

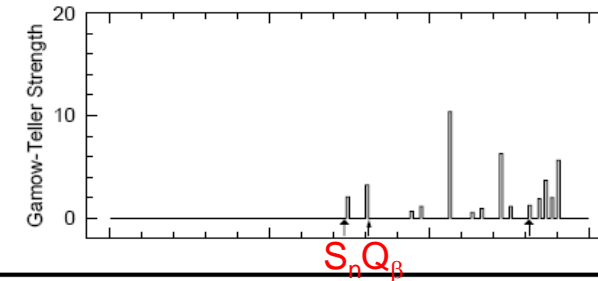
(2) as in (1) with empirical spreading of SP transition strength, as shown in experimental $S_\beta(E)$

(3) as in (2) with addition of first-forbidden strength from Gross Theory

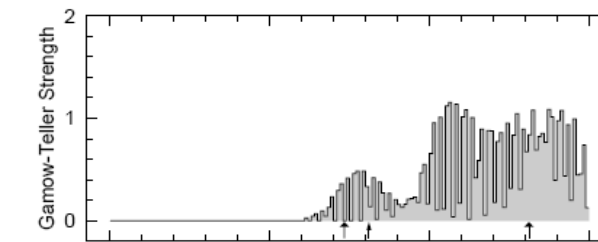
“Typical spherical example”:

β -Decay of ^{92}Rb in 3 Successively Improved Models

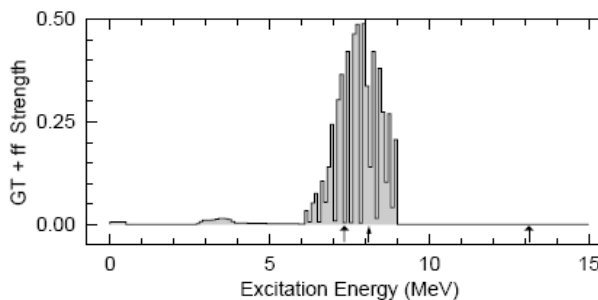
(Exp.: $T_{1/2} = 4.49$ s $P_n = 0.01$ %)



QRPA
Pure GT strength
Pure q.p. levels
 $P_n = 82.81$ %
 $T_{1/2} = 1.30$ (h)



QRPA
Pure GT strength
Spreading of
q.p. levels
 $P_n = 4.14$ %
 $T_{1/2} = 12.05$ (min)



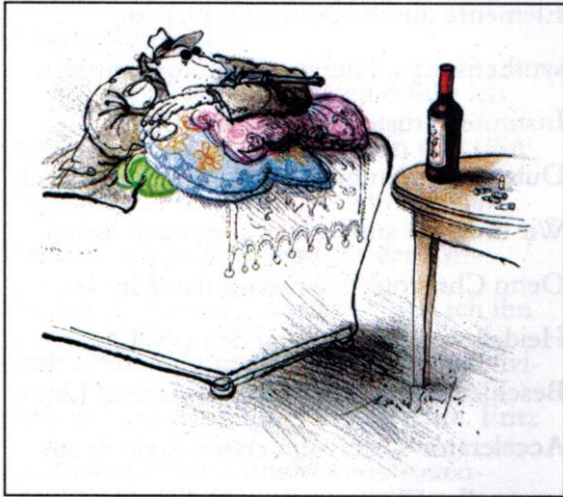
QRPA (GT)
+
Gross Th. (ff)
with spreading
 $P_n = 0.05$ %
 $T_{1/2} = 8.47$ (s)

note: effects on $T_{1/2}$ and P_n !

The Astrophys. Journal 792 (2014)

A High-Entropy-Wind r-Process Study Based on Nuclear-Structure Quantities from the New Finite-Range Droplet Model FRDM(2012)

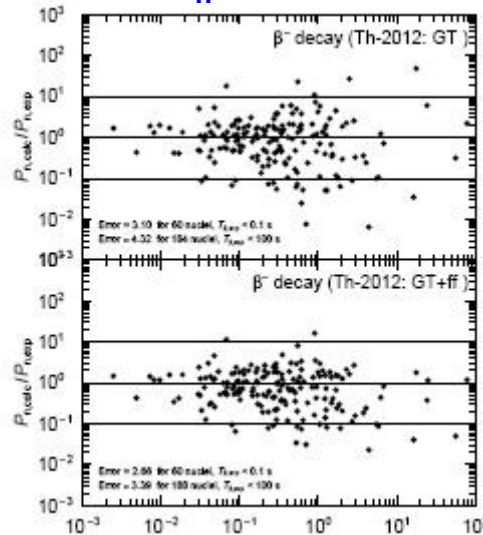
K.-L. Kratz, K. Farouqi, **P. Möller**



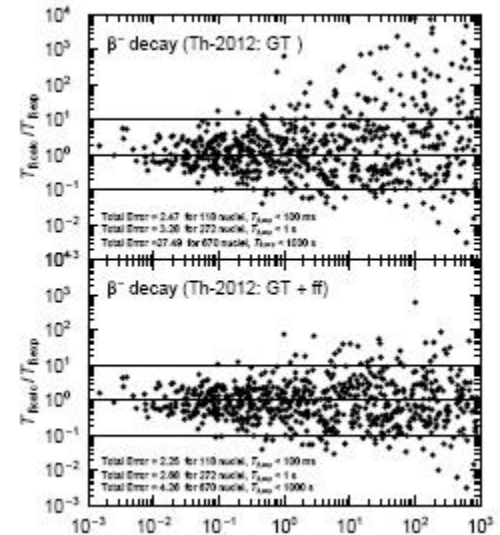
Patience is required...

...consistent nuclear-data input for masses and β -decay properties from FRDM(2012) and QRPA(2012)

P_n values



Half-lives



$T_{1/2}(\text{calc})/T_{1/2}(\text{exp})$

2002 3.9

2003 3.1

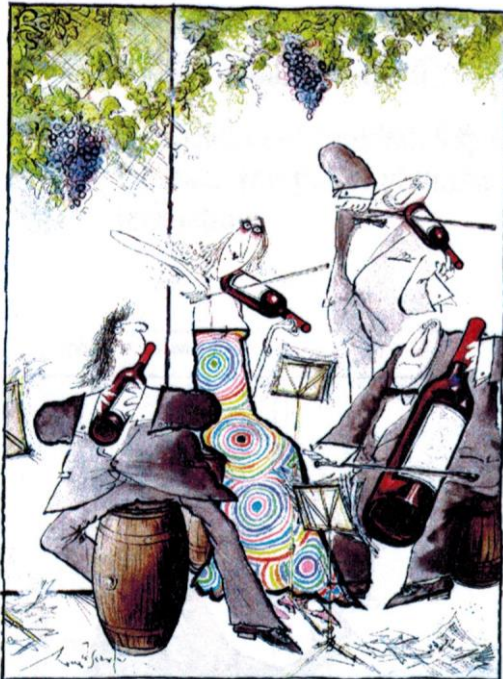
2012 2.7

Solar system r-process abundances

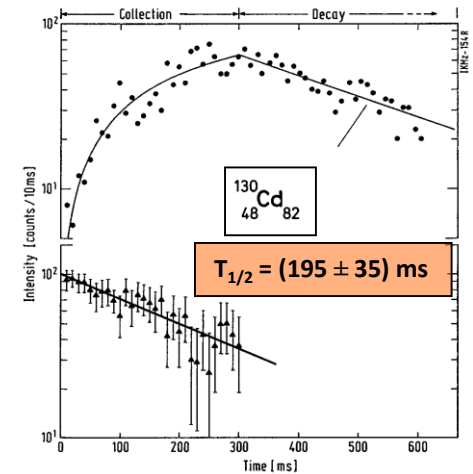
Inst. Phys. Conf. Ser. No. 88 / J. Phys. G: Nucl. Phys. 14 Suppl. (1988)

Constraints on r-process conditions from beta-decay properties far off stability and r-abundances

K-L Kratz, F-K Thielemann, W Hillebrandt, P Möller, V Harms, A Wöhr
and J W Truran



Concerted action



Model predictions at that time:
 $30 \text{ ms} \leq T_{1/2} \leq 1.2 \text{ s}$

The Astrophys. J. 403 (1993)

ISOTOPIC r -PROCESS ABUNDANCES AND NUCLEAR STRUCTURE FAR FROM STABILITY: IMPLICATIONS FOR THE r -PROCESS MECHANISM

KARL-LUDWIG KRATZ

Institut für Kernchemie, Universität Mainz, Fritz-Straßmann-Weg 2, D-W-6500 Mainz, Germany

JEAN-PHILIPPE BITOUZET¹ AND FRIEDRICH-KARL THIELEMANN

Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138

AND

PETER MÖLLER AND BERND PFEIFFER

Institut für Kernchemie, Universität Mainz, Fritz-Straßmann-Weg 2, D-W-6500, Mainz, Germany

Received 1991 August 13; accepted 1992 July 14

The FK²L waiting-point approach (I)

With the nuclear physics knowledge at that time...

J. Phys. G 24 (1988)

With $T_{1/2}(\text{exp})$ and $N_{r,0}$

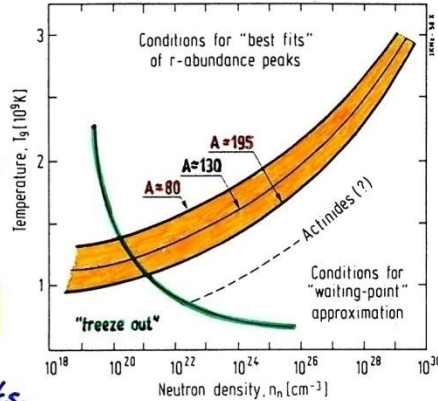
test of

"waiting-point concept":

correlate $T_{1/2} \Rightarrow N_{r,0}$

deduce $T_{1/2} - n_n$ band

conditions to fit the $A \approx 80$ and 130 $N_{r,0}$ -peaks.

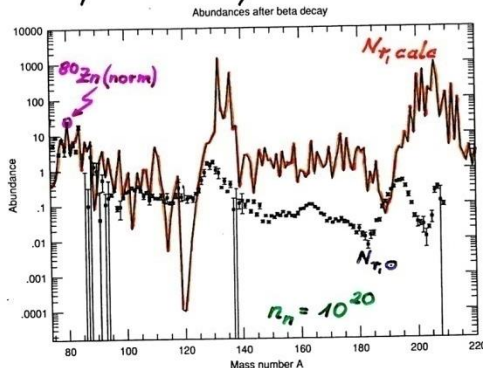


When we (Hillebrandt, Kratz, Möller, Thielemann) started in 1986, already numerous attempts existed to search for the site(s) of the r-process... but **none really successful**.

Ap. J. 403 (1993)

Static and time-dependent r-process calculations for "freeze-out" conditions.

Use of internally consistent nuclear-physics input (FRDM + QRPA).



Steady-flow NOT global

- wrong trend, increasing N_r with A
- $A \approx 130$ and 135 peaks shifted, too large;
- \Rightarrow indicates too low n_n

Therefore, our approach:

The "site-independent" waiting-point concept was utilised to deduce unique astrophysical conditions for an r-process.

For the first time, a "unified model" (FRDM & RPA) for all nuclear properties was used, aided by the first experimental nuclear data of r-process isotopes:

N=50 ⁷⁹Cu, ⁸⁰Zn, ⁸¹Ga,

N=56,57 ^{91,92}Br,

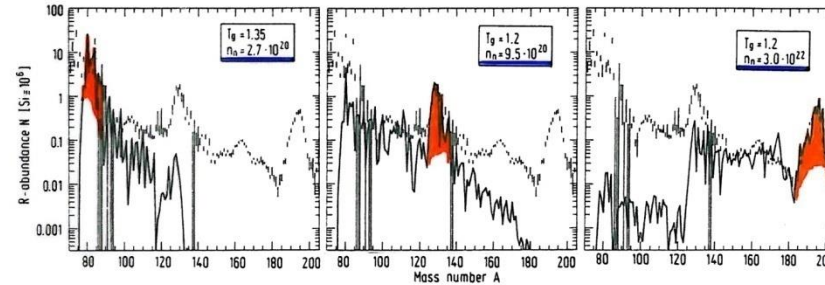
N=60-63 ⁹⁷⁻¹⁰⁰Rb,

N=82 ¹³⁰Cd, ¹³¹In

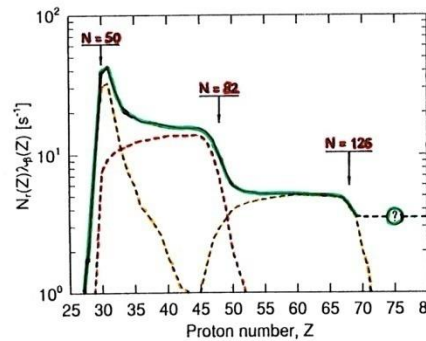
The FK²L waiting-point approach (II)

Consequences:

- r -process must have several components with different n -densities and different r -process paths.



- steady-flow is only local, breaks down at each N_T -peak \rightarrow at N_{magic}



analogous to s -process with $N_s G = const.$

empirical r -process picture with superposition of (minimum) 3 components

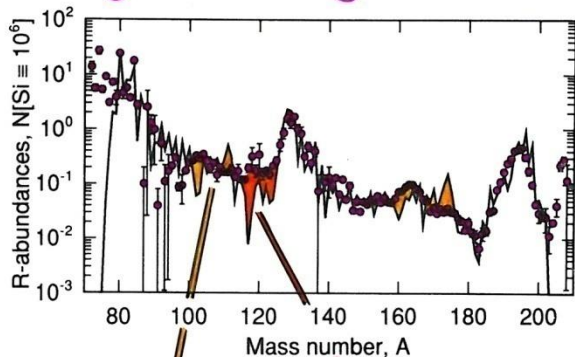
$$\boxed{N_T \lambda_T = const.}$$

weighting acc. to $T_{1/2}$ ($^{80}Zn, ^{130}Cd, ^{135}Tm$)

The FK²L waiting-point approach (III)

Superposition of 3 components

↗ good overall agreement with $N_{r,0}$



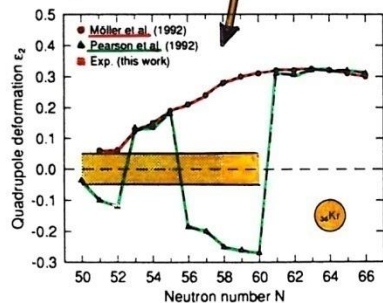
but...

Local deficiencies

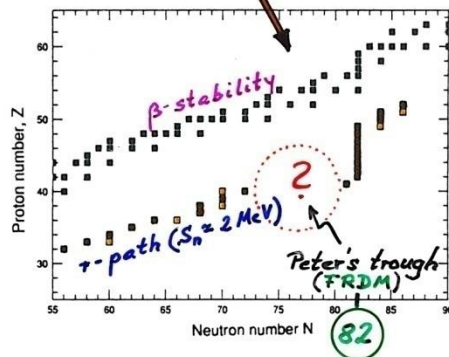
interpreted as nuclear-structure signatures of r-isotopes not accessible in terrestrial laboratories

⇒ model deficiencies!

(i) phase transitions
pn-residual interact.



(ii) shell corrections
 $N=50, 82, 126$



first euphoric statement by
W. Hillebrandt:
“...best $N_{r,0}$ fit so far...;
long-standing problem solved...”

↻ birth of $N=82$
“shell-quenching”
idea ...

...this catchword coined by
W. Nazarewicz later led to
numerous misinterpretations...

Consequences of the FK²L waiting-point approach

1. A unified approach

...for the first time all nuclear properties were studied in a selfconsistent way

2. Remaining deficiencies

...missing monopole and quadrupole p–n residual interactions

Consequences:

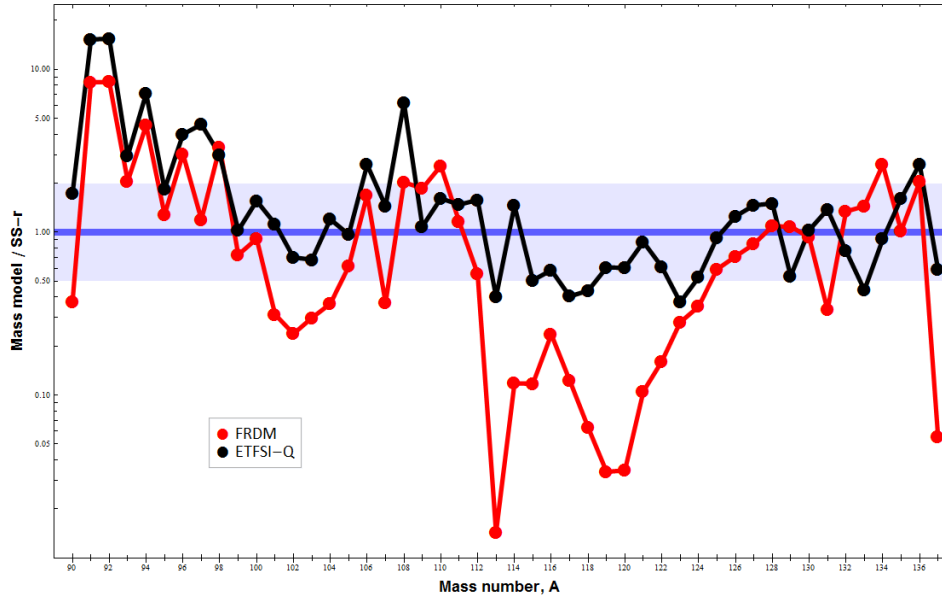
with the filling of the $1g_{9/2}$ proton shell, the $1g_{7/2}$ neutron orbitals are lowered, not contained in the QRPA:

the spherical N=56 subshell strength in the QRPA is underestimated, leading to a too early onset of deformation over a too wide Z-range;

overestimation of the Z=50 and N=82 shell strengths \Rightarrow

r-path moved from $(Z,N) = (40,72)$ to $(Z,N) = (41,81)$ leaving a gap in A of 10 units, where not a single waiting-point isotope will exist; this leads, for example, to the famous A ~ 120 r-abundance trough (“Peter’s trough”)

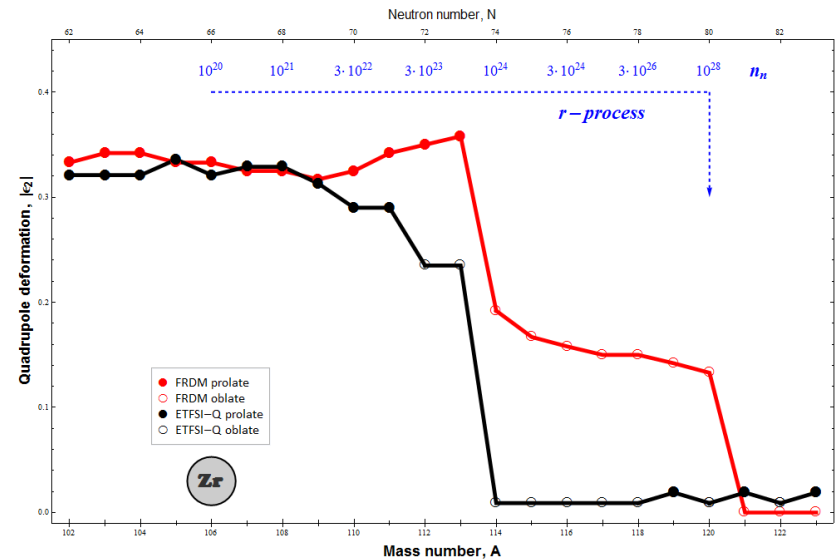
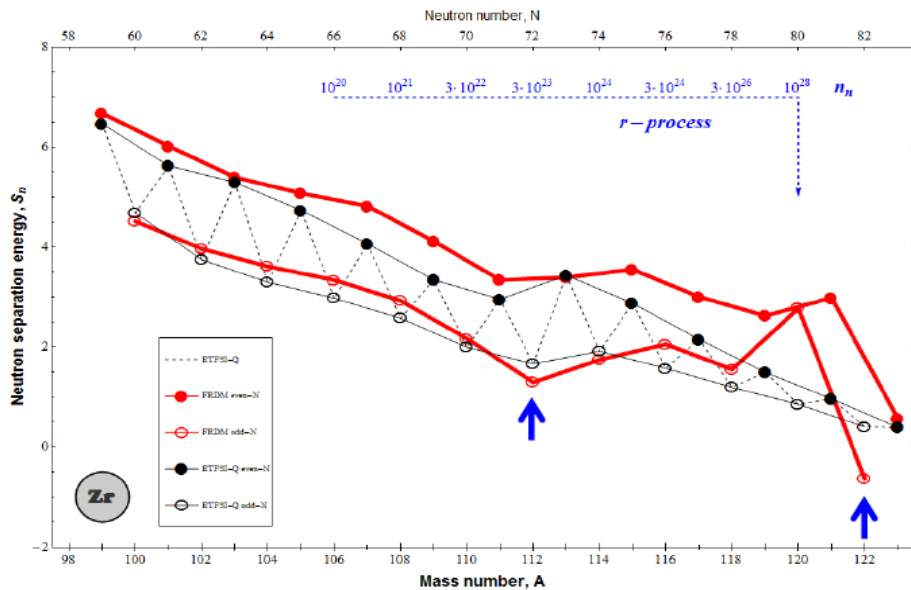
Deviation from SS-r: FRDM vs. ETFSI-Q



How to fill up the FRDM $A \approx 115$ “trough”?

e.g. “tampering” with the $T_{1/2}$ would require completely unrealistic long $T_{1/2}$ (w.-p.) of the order 1 – 20 s

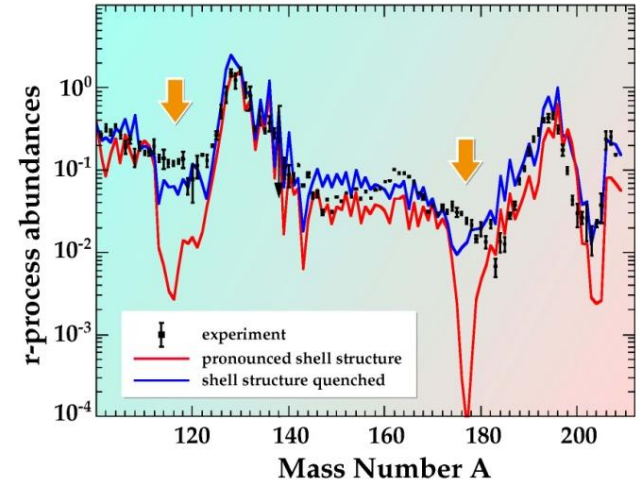
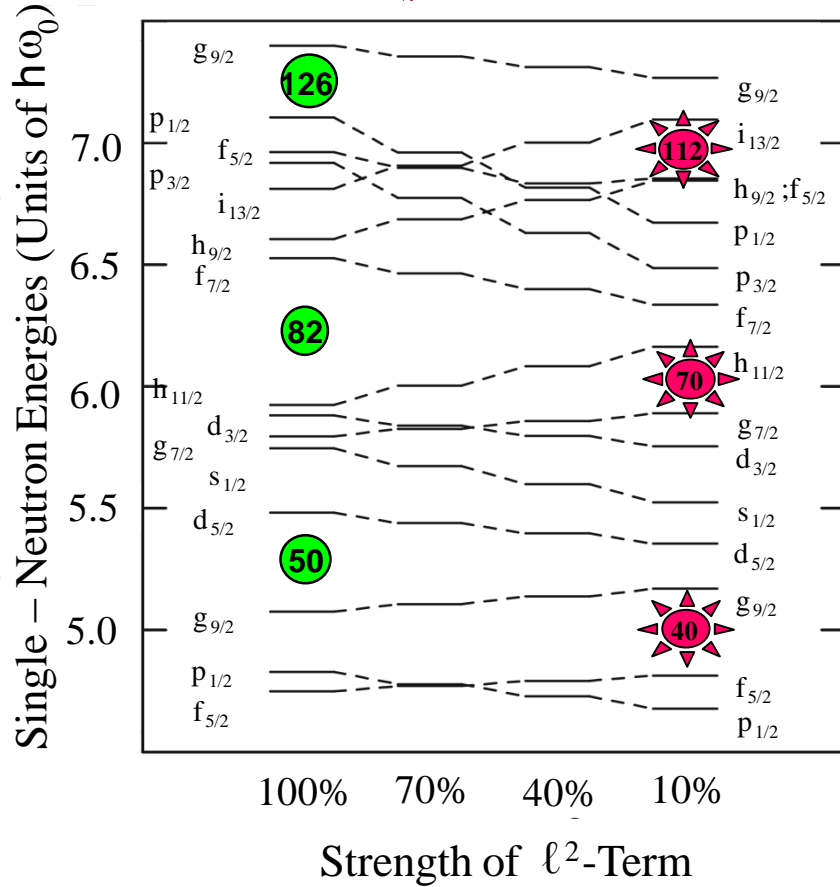
⇒ hence, the solution **MUST** lie in the masses (S_n) and the correlated trend of deformation in this shape-transition region before $N=82$



Effects of N=82 "shell quenching"

...reduction of the spin-orbit coupling strength; caused by interaction between bound and continuum states.

$N/Z \Rightarrow$



change of

- shell-gaps
- deformation
- r-process path (S_n)
- r-matter flow (τ_n)

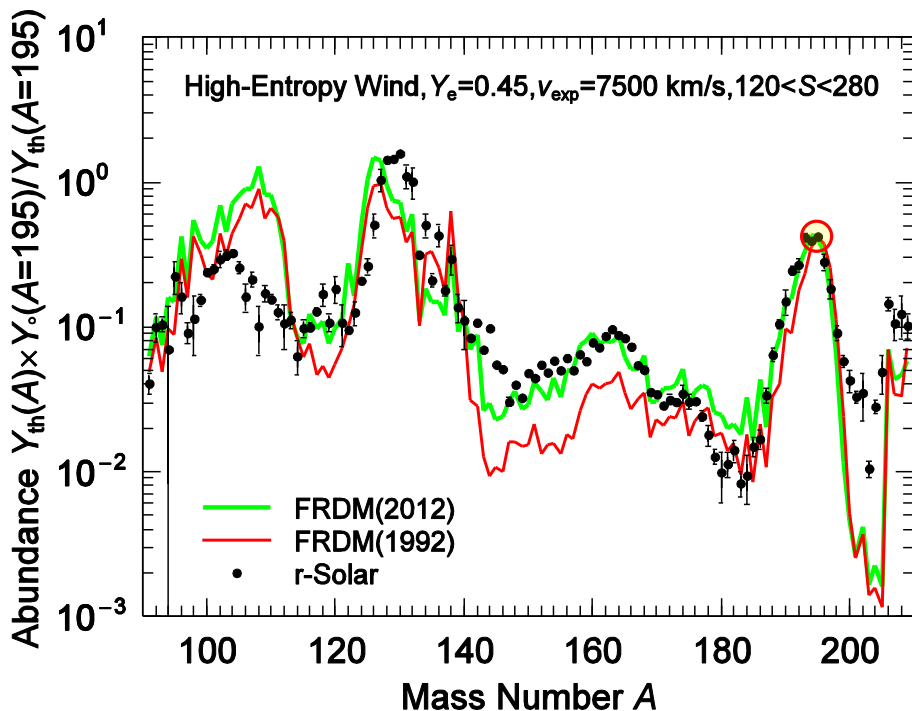
- high-j orbitals \Uparrow (e.g. $vg_{9/2}$, $vh_{11/2}$)
- low-j orbitals \Downarrow (e.g. $vd_{5/2}$, $vf_{7/2}$)
- new "magic" numbers (e.g. 40, 70, 112)

...controversial discussions until today !

First HEW calculations with FRDM(2012) and QRPA(2012)

Good news at the end...

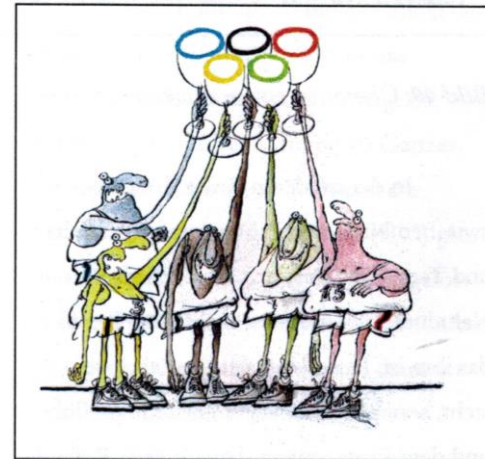
Comparison between $N_{r,\odot}$ and $N_{r,calc}$:
FRDM(1992) and **FRDM(2012)**



Improvements and remaining deficiencies:

- still overabundances in $80 \leq A \leq 110$ region
- “abundance trough” at $A \approx 120$ removed
- 2nd r-peak slightly improved
- N=82 bottle-neck behavior improved
- REE “pygmy-peak” well reproduced
- shape of 3rd r-peak well reproduced
- shape-transition region above N=126 still imperfect \Rightarrow deep trough
- Pb, Bi too low \Rightarrow contribution from α -backday not yet included

Success !!!



Conclusion nuclear-structure models

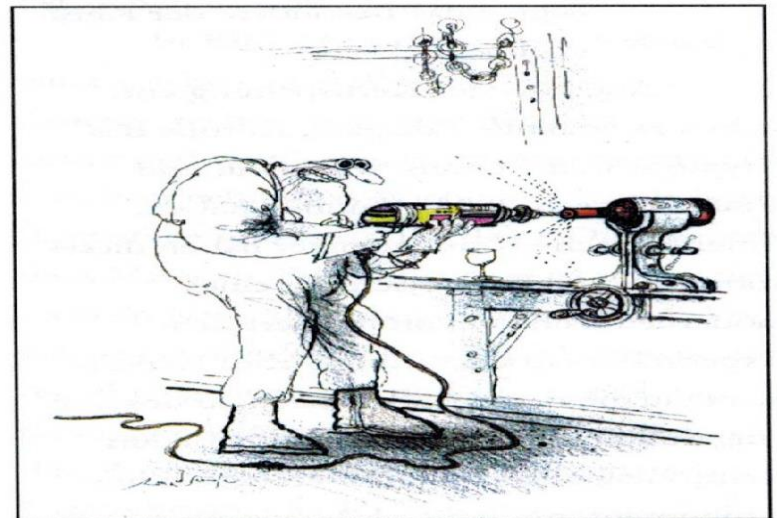
P. Möller, B. Pfeiffer & K.-L. Kratz, Phys. Rev. C67 (2003)

“...let us emphasize that **there is no “correct” model in nuclear physics**. Any modeling of nuclear-structure properties involves **approximations**... with the goal to obtain a formulation that can be solved in practice, but that **“retains the essential features”** of the true system under study, so that one can still learn something.

It may well turn out, that when proceeding from a simplistic, macroscopic approach to a more microscopic model the first overall result may be worse just in terms of agreement between calculated and measured data. However, the **disagreements may now be understood** more easily, and further microscopic-based, realistic improvements will become possible.”

...drill to depth

... but still, Nature sometimes disagrees with my nice model.



Summary and conclusion

...considerable progress during our 3 decades of collaboration,

but still a lot remains to be done
for the coming 30 years
in all interrelated fields...



However, right now
let's first celebrate
Peter's 70th birthday



Happy birthday !

Bernd, Khalil, Oliver, K.-L.
and Gisela