

# *Surrogate reactions for $(n,f)$ and $(n,g)$ cross sections*

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Nuclear Fission and Structure of Exotic Nuclei  
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STARLiTeR Collaboration

 Lawrence Livermore  
National Laboratory



LLNL-PRES-XXXXXX

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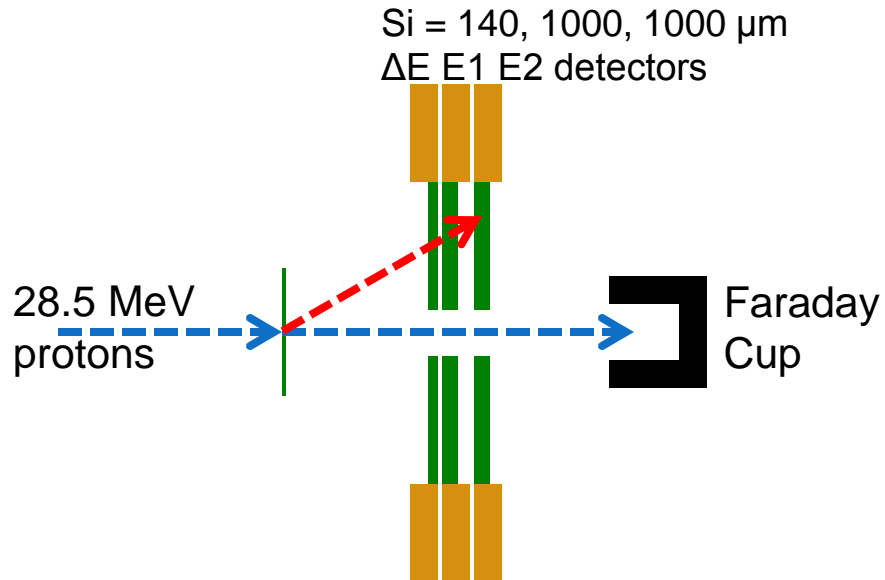
# Overview

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- Brief discussion of experimental setup – STARLiTeR
- Surrogate ratio measurements of (n,f) cross sections
  - Am24X
  - Pu23X
- Absolute surrogate measurement of  $^{87}\text{Y}(n,g)$
- Summary



# StarLiTeR at Texas A&M



We record the total energy, angle and determine the particle type

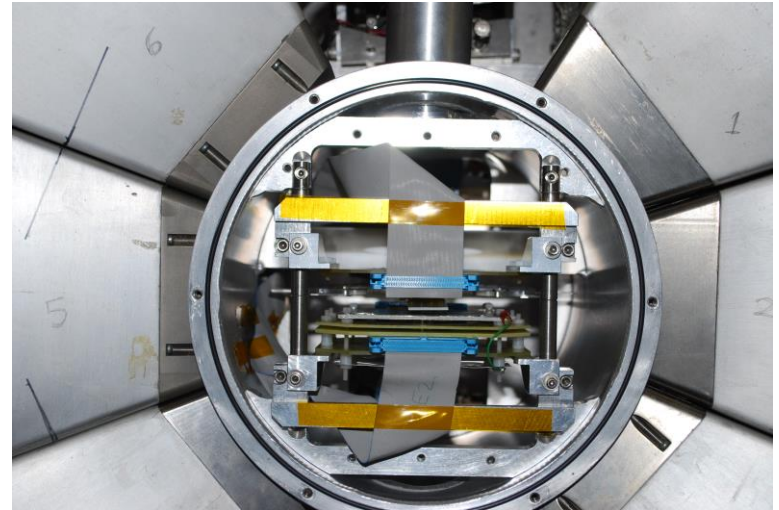
We have run 10 experiments totaling 11 weeks of beam time since April 2012.

Six more experiments planned through summer 2014.

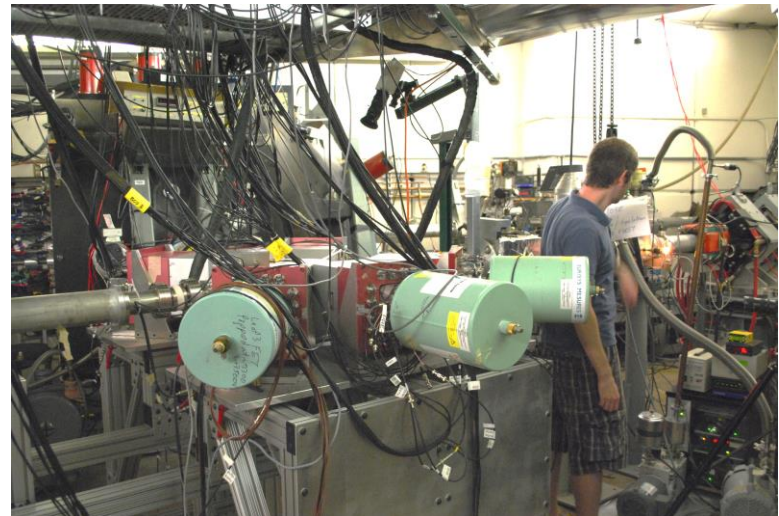
Typical values:

- Si energy resolution 75 keV one sigma
- Gamma energy resolution 1.7 keV one sigma
- Angle range 30 to 60 degrees

Top view of the STARLiTeR array. Beam enters from top and exits bottom. HPGe surrounds chamber Si telescope in center.



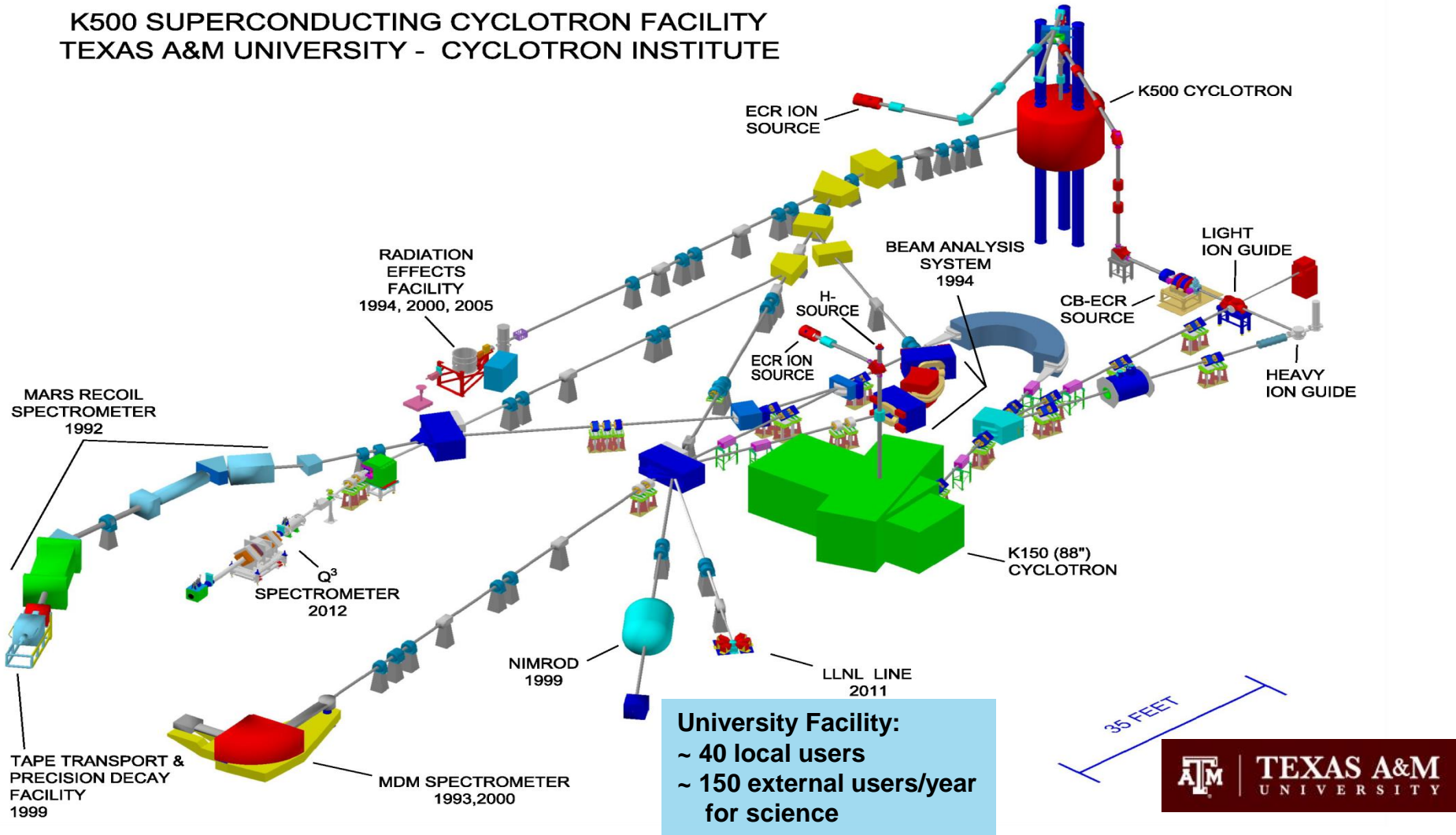
STARLiTeR array closed up around the chamber. Cyclotron beam comes in from the right and exits on the left.



# Texas A&M Cyclotron Institute, College Station, Texas

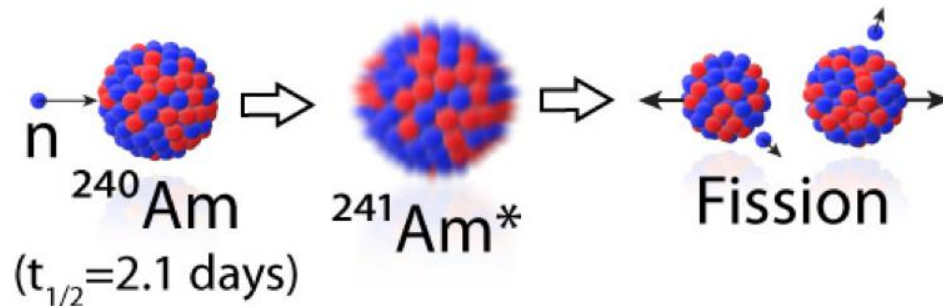
Two machines: K150 and K500 Cyclotrons stable beams available and radioactive beams under development

## K500 SUPERCONDUCTING CYCLOTRON FACILITY TEXAS A&M UNIVERSITY - CYCLOTRON INSTITUTE



# The desired reaction: $^{240}\text{Am}(n,f)$

This reaction has never been measured.  $^{240}\text{Am}$  has a half-life of 2.1 days, which makes it unreasonable to use as a target.



$^{239}\text{Am}$ 12 hours	$^{240}\text{Am}$ 2.1 days	$^{241}\text{Am}$ 433 years	$^{242}\text{Am}$ 16 hours	$^{243}\text{Am}$ 7370 years	$^{244}\text{Am}$ 10 hours
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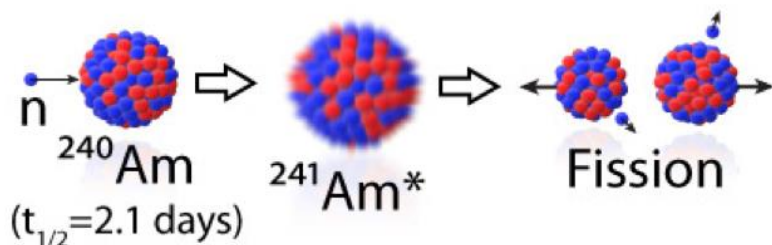
Desired ← Surrogate

The surrogate ratio technique can be used, by populating the same compound nucleus using a longer-lived target.

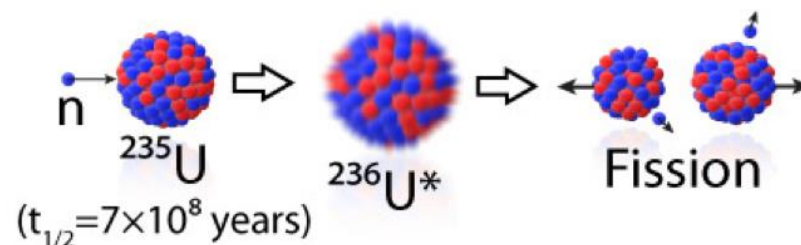


# The Surrogate Reaction

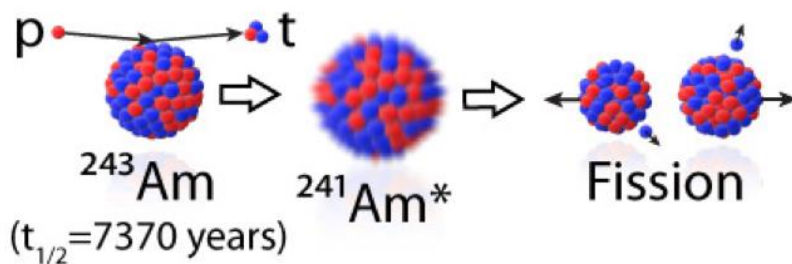
**Desired Reaction:**  $^{240}\text{Am}(n,f)$



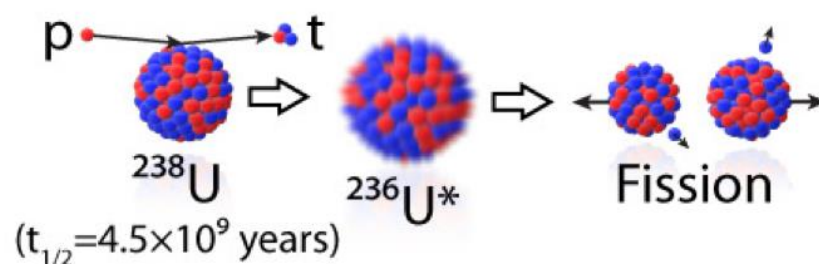
**Reference Reaction:**  $^{235}\text{U}(n,f)$



**Surrogate Reaction:**  $^{243}\text{Am}(p,tf)$



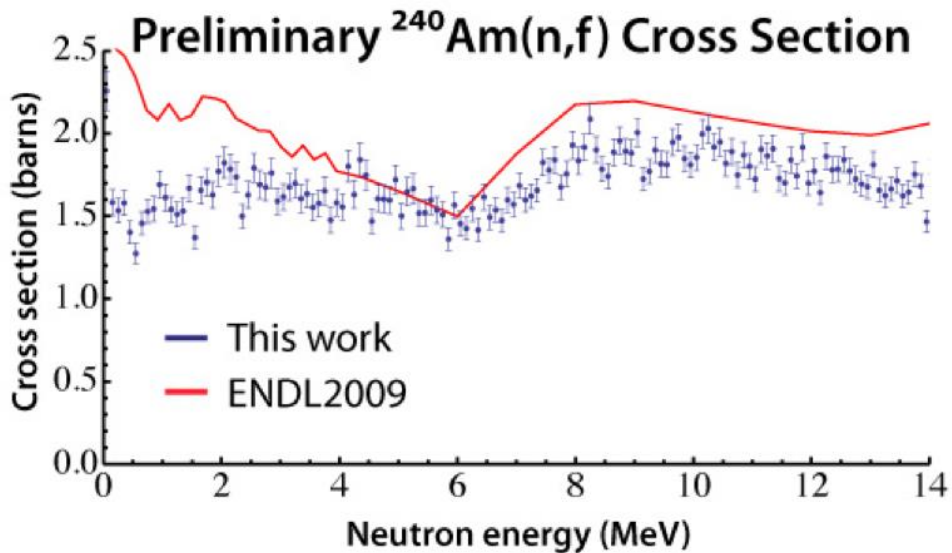
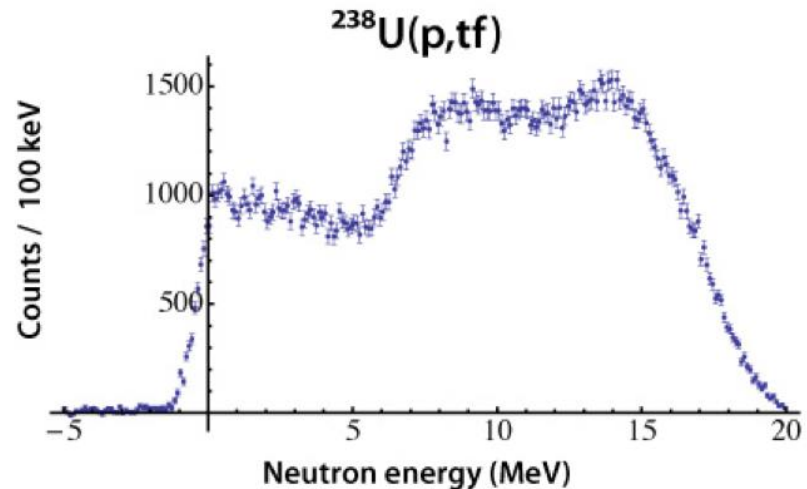
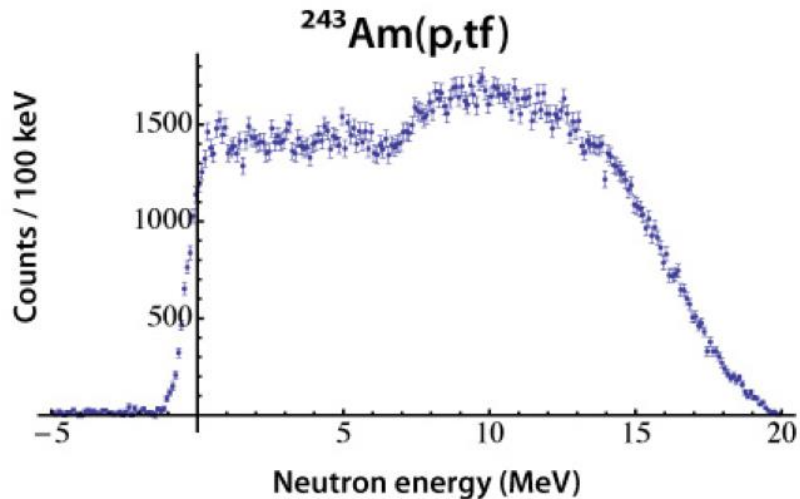
**Ratio Reaction:**  $^{238}\text{U}(p,tf)$



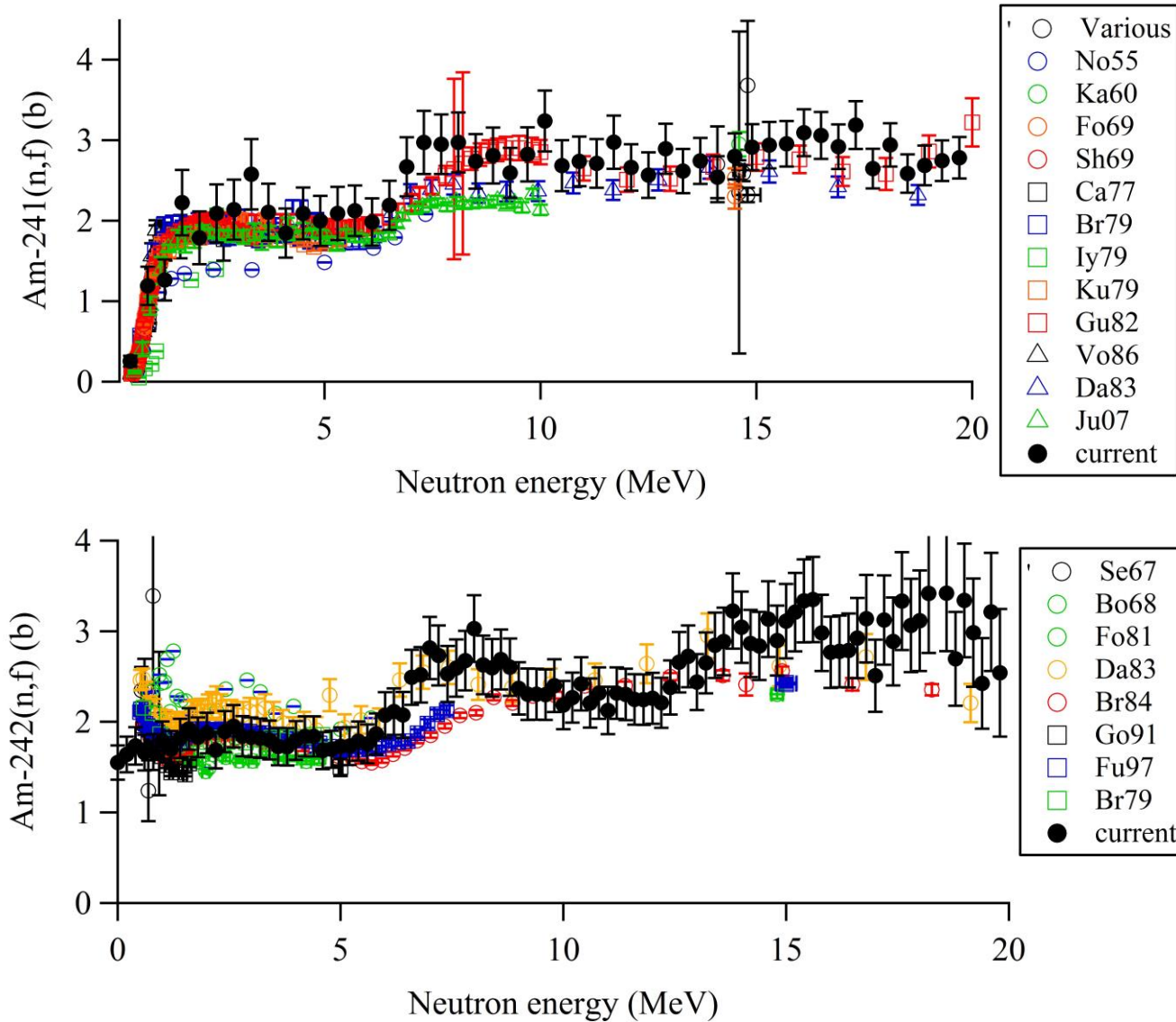
$$\sigma(^{240}\text{Am}(n, f), E) = \frac{N(^{243}\text{Am}(p, tf), E)}{N(^{238}\text{U}(p, tf), E)} \times \frac{\sigma_{CN}(^{240}\text{Am}(n, tot)^{241}\text{Am}^*, E)}{\sigma_{CN}(^{235}\text{U}(n, tot)^{236}\text{U}^*, E)} \times \sigma(^{235}\text{U}(n, f), E)$$



# Particle spectra and $^{240}\text{Am}(n,f)$



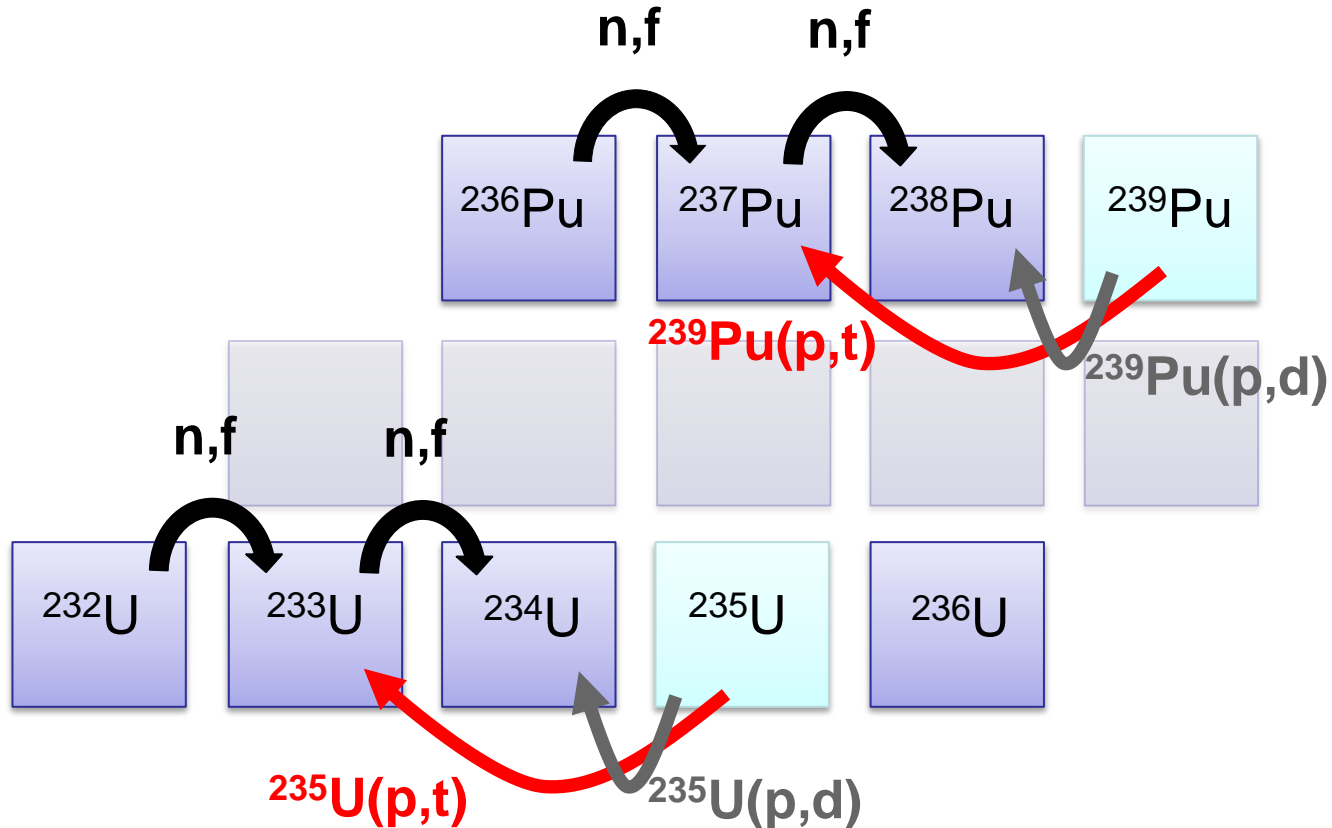
**$^{242}\text{Am}(n,f)$  and  $^{241}\text{Am}(n,f)$  cross section measurements using the surrogate ratio  $^{243}\text{Am}(^3\text{He},^3\text{He}')$  and  $^{243}\text{Am}(^3\text{He},^4\text{He})$  direct reactions and  $^{236}\text{U}(^3\text{He},^3,^4\text{He})$**   
 Courtesy Jo Ressler





# $^{23x}\text{Pu}$ cross section measurements using the Surrogate Ratio

Courtesy Richard Hughes



$$S(^{236}\text{Pu}(n, f), E) = \frac{N(^{239}\text{Pu}((p, tf), E))}{N(^{235}\text{U}((p, tf), E))} \cdot S(^{232}\text{U}(n, f), E)$$

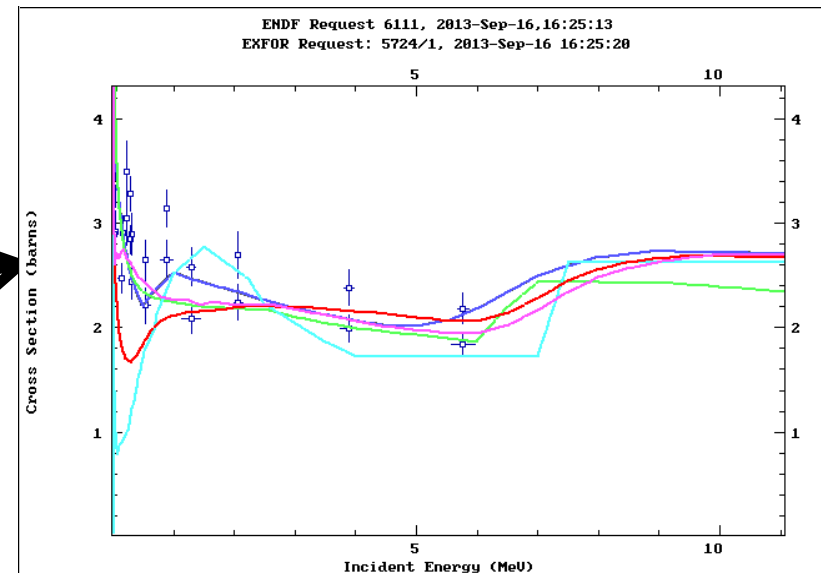


# Investigating $^{236}\text{Pu}$ (n,f) cross sections

Courtesy Richard Hughes

- Half lives rapidly fall off for isotopes lighter than  $^{239}\text{Pu}$ :

Isotope	$T_{1/2}$
$^{239}\text{Pu}$	24110 yrs
$^{238}\text{Pu}$	87.7 yrs
$^{237}\text{Pu}$	45.64 days
$^{236}\text{Pu}$	2.86 yrs

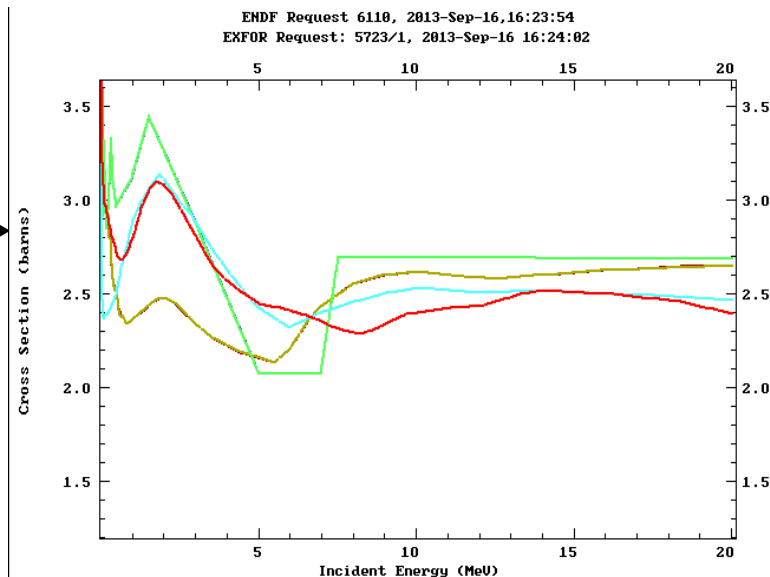
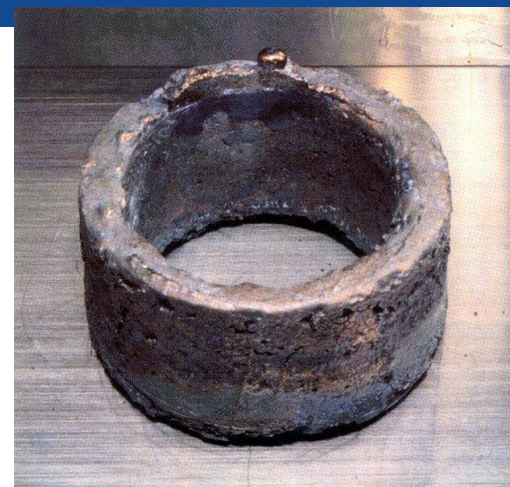


# Investigating $^{237}\text{Pu}$ (n,f) cross sections

Courtesy Richard Hughes

- Half lives rapidly fall off for isotopes lighter than  $^{239}\text{Pu}$ :

Isotope	$T_{1/2}$
$^{239}\text{Pu}$	24110 yrs
$^{238}\text{Pu}$	87.7 yrs
$^{237}\text{Pu}$	45.64 days
$^{236}\text{Pu}$	2.86 yrs

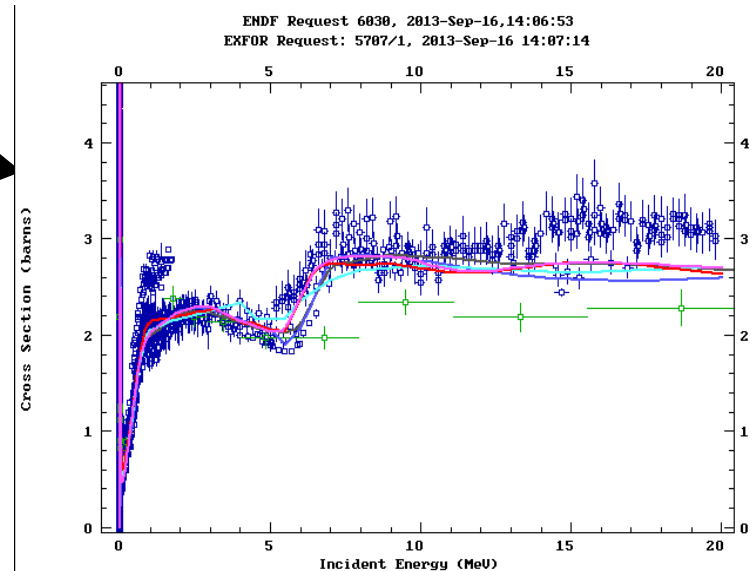


# Investigating $^{238}\text{Pu}(n,f)$ cross sections

Courtesy Richard Hughes

- Half lives rapidly fall off for isotopes lighter than  $^{239}\text{Pu}$ :

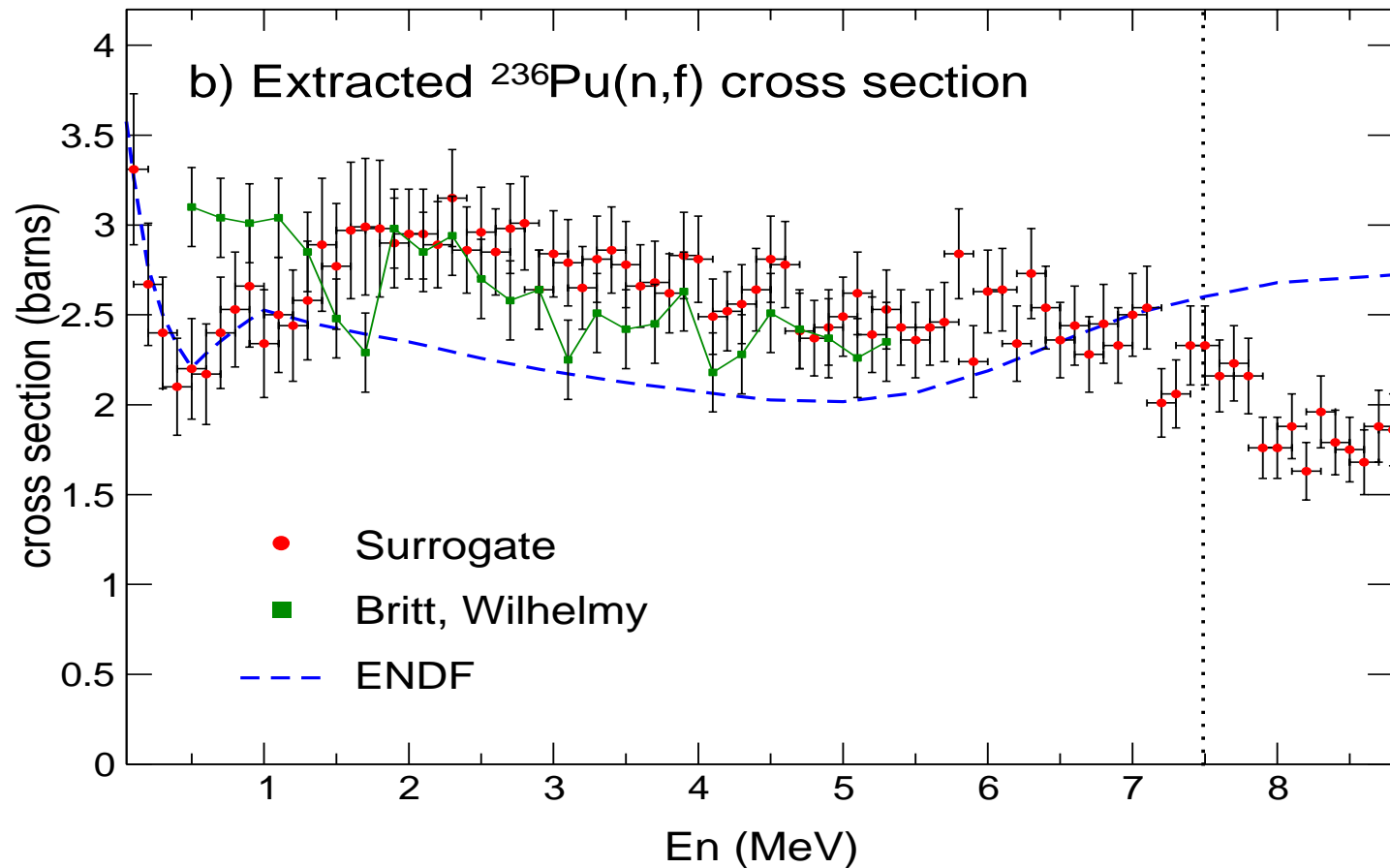
Isotope	$T_{1/2}$
$^{239}\text{Pu}$	24110 yrs
$^{238}\text{Pu}$	87.7 yrs
$^{237}\text{Pu}$	45.64 days
$^{236}\text{Pu}$	2.86 yrs



# $^{236}\text{Pu}(n,f)$ cross section from $^{239}\text{Pu}(p,t)^{237}\text{Pu}^*/^{235}\text{U}(p,t)^{233}\text{U}^*$ surrogate ratio

Courtesy Richard Hughes

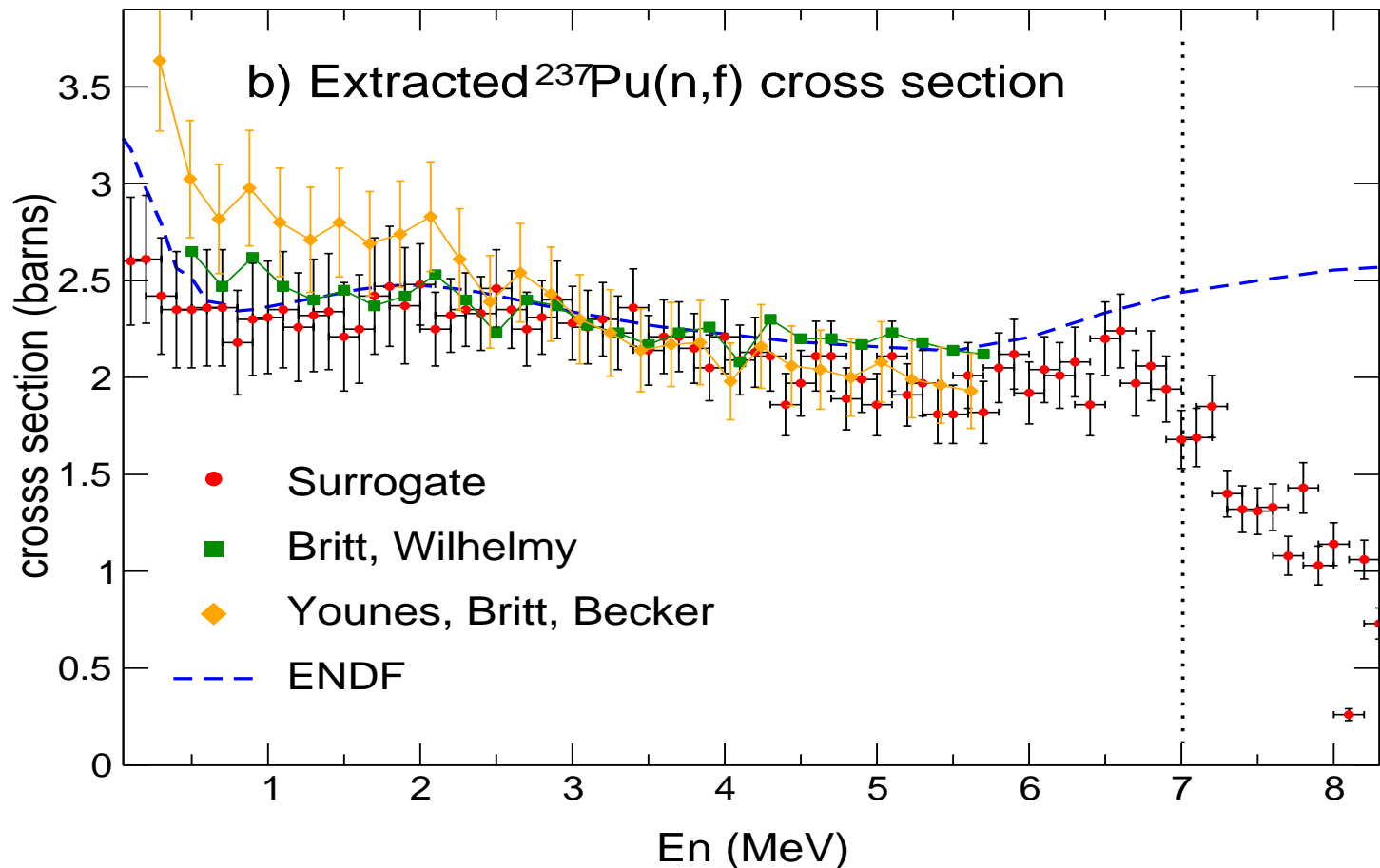
$$S(^{236}\text{Pu}(n,f), E) = \frac{N(^{239}\text{Pu}((p,tf), E))}{N(^{235}\text{U}((p,tf), E))} \cdot S(^{232}\text{U}(n,f), E)$$



# $^{237}\text{Pu}(n,f)$ cross section from $^{239}\text{Pu}(p,d)^{238}\text{Pu}^*/^{235}\text{U}(p,d)^{234}\text{U}^*$ surrogate ratio

Courtesy Richard Hughes

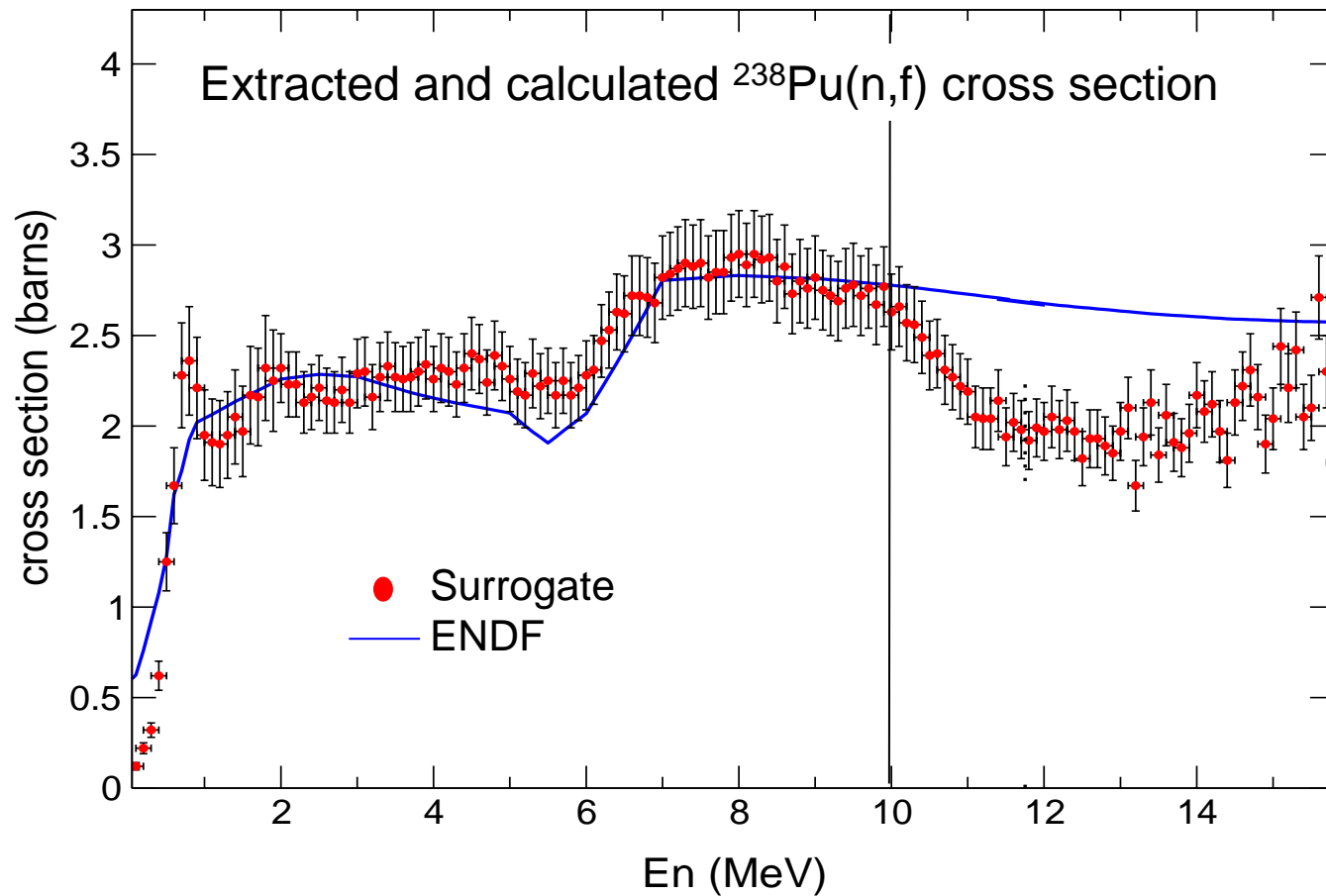
$$S(^{237}\text{Pu}(n,f), E) = \frac{N(^{239}\text{Pu}(p,df), E)}{N(^{235}\text{U}(p,df), E)} \cdot S(^{233}\text{U}(n,f), E)$$



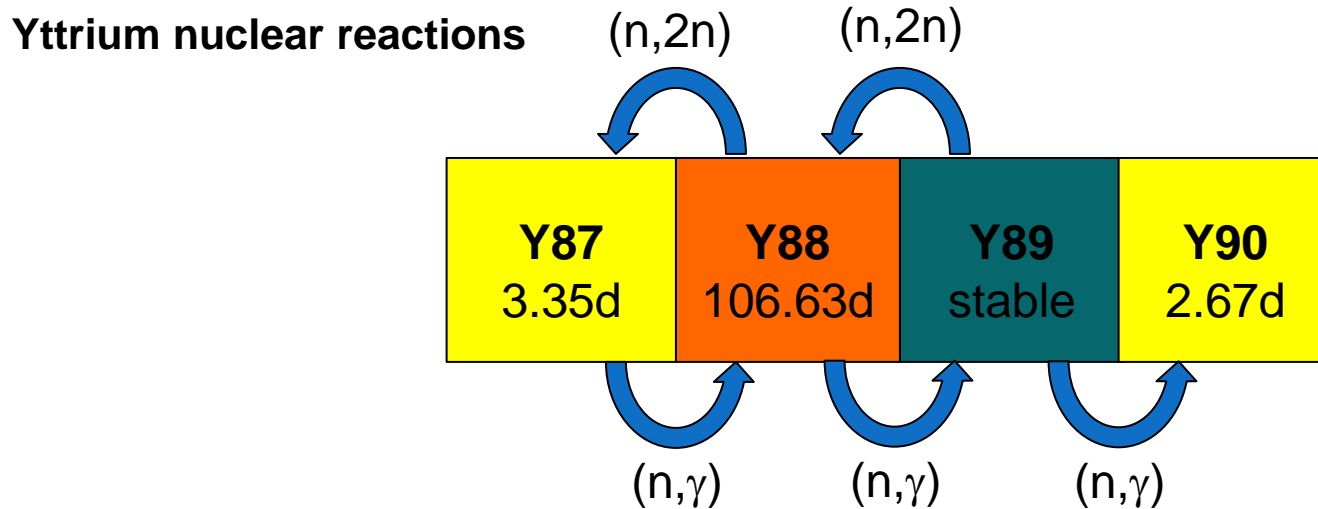
# $^{238}\text{Pu}(n,f)$ cross section from $^{239}\text{Pu}(p,p')^{239}\text{Pu}^*/^{235}\text{U}(p,p')^{235}\text{U}^*$ surrogate ratio

Courtesy Richard Hughes

$$S(^{238}\text{Pu}(n,f), E) = \frac{N(^{239}\text{Pu}(p,p'), E)}{N(^{235}\text{U}(p,p'), E)} \cdot S(^{234}\text{U}(n,f), E)$$



# Where's the experimental challenge in determining $^{87}\text{Y}(n,\gamma)$ ?



Y89 is the only stable Yttrium isotope.

Nearly impossible to make and irradiate a Y87 target to get an energy dependent cross section.

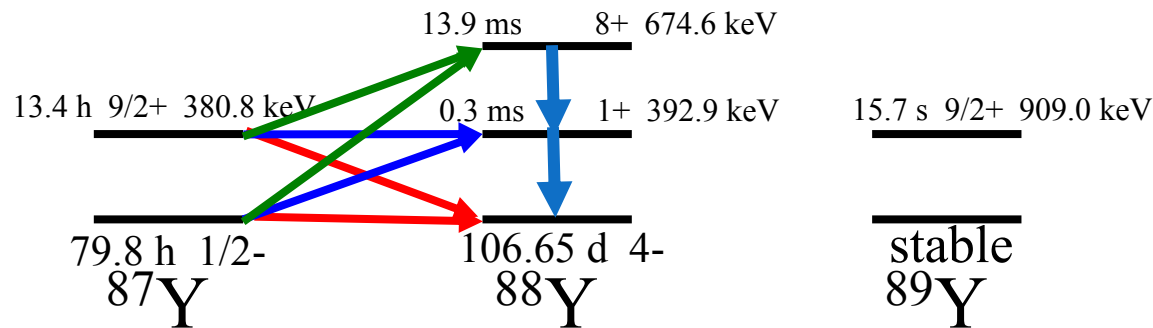
Ex: 100 micrograms of Y87 equals  $1.7 \times 10^{12}$  becquerels or 44.8 Curies

Ex: 100 micrograms of Y88 equals  $5.2 \times 10^{10}$  becquerels or 1.5 Curies

Yttrium isotopes also contain multiple isomeric states...



# Yttrium isotopes and isomer levels



We ran for 13 days (39 shifts) starting June 24, 2012. The experiment involved ~20 students, post-docs, and researchers from >5 institutions. We collected 3 TeraBytes of data on multiple targets focusing on Y89 (also Zr9x).

- Si energy resolution 75 keV one sigma
- Gamma energy resolution 1.7 keV one sigma
- Particle angular range 30 to 60 degrees
- Y89 mono-isotopic 760 micrograms/cm<sup>2</sup>
- Zr9x ~1000 micrograms/cm<sup>2</sup>

# Strategies for constraining HF inputs

## I. Calculate $\gamma$ SF & LD:

- Theory challenging, not all nuclei covered, but progress is being made
- Experiments needed to verify theory & improve

## II. Constraints from neighbors:

- Measure  $(n,\gamma)$  cross sections in other nuclei & do regional fits
- Extrapolations required

## III. Constraints from surrogate observables:

- Measure quantities in **actual nuclei of interest**
- Theory needed to relate measurement to desired cross section

# The surrogate method

## Experiment

Produce CN  $^{88}\text{Y}$  via alternative reaction  $p + ^{89}\text{Y} \rightarrow d + ^{88}\text{Y}$  involving stable  $^{89}\text{Y}$   
 Measure outgoing surrogate particle **d** in coincidence with observables indicative of relevant decay channel  $\rightarrow P_{\delta\gamma}(E)$

Here:  $\gamma$  rays measured in coincidence with deuterons:

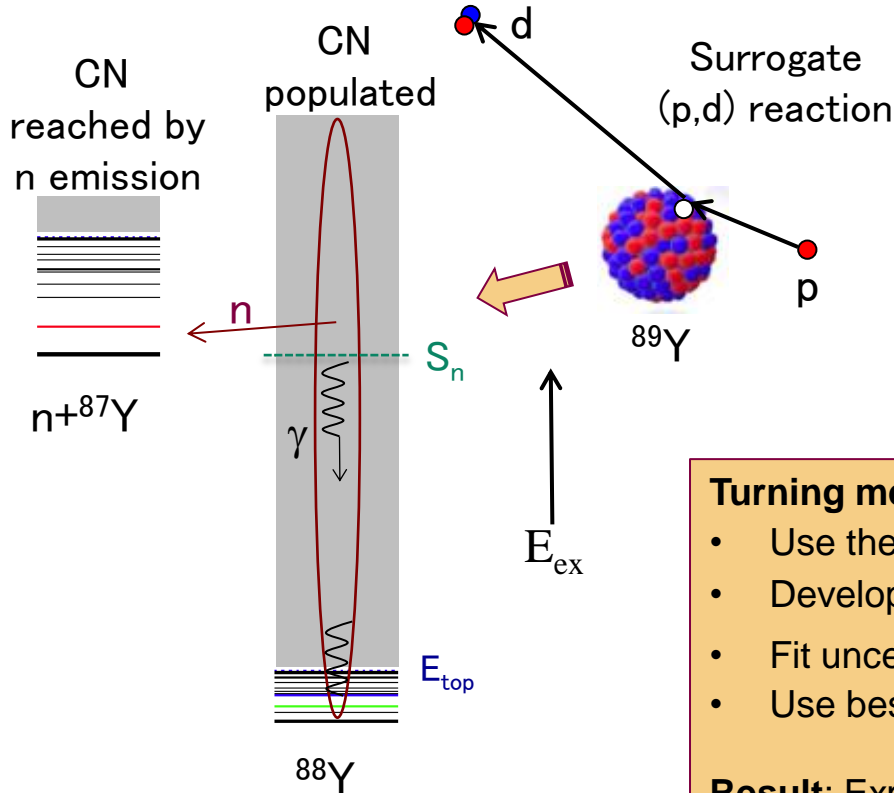
$$P_{\delta\gamma}(E) = \frac{N_{\delta-\gamma}(E_{\text{exc}})}{N_{\delta}(E_{\text{exc}})}$$

## Formalism

A Surrogate experiment gives

$$P_{\delta\gamma}(E) = \sum_{J,\pi} F_{\delta}^{\text{CN}}(E,J,\pi) \cdot G_{\gamma}^{\text{CN}}(E,J,\pi)$$

Hauser-Feshbach description of "desired" CN reaction

$$\sigma_{\alpha\gamma} = \sum_{J,\pi} \sigma_{\alpha}^{\text{CN}}(E,J,\pi) \cdot G_{\gamma}^{\text{CN}}(E,J,\pi)$$


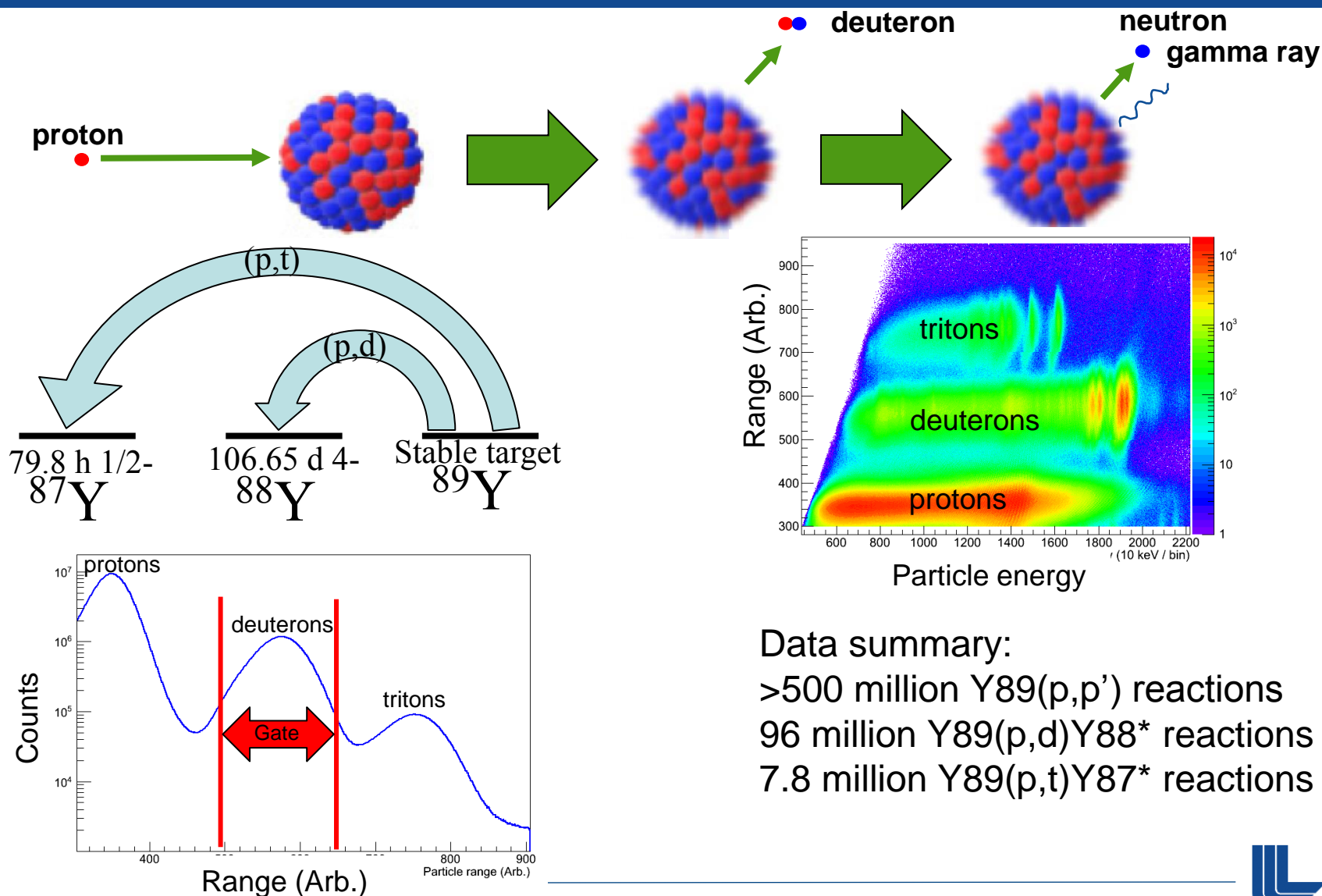
### Turning measurement into cross section

- Use theory to describe Surrogate reaction, predict  $F_{\delta}^{\text{CN}}$
- Develop rough decay model  $G_{\gamma}^{\text{CN}}$
- Fit uncertain parameters in  $G_{\gamma}^{\text{CN}}$  to reproduce  $P_{\delta\gamma}$
- Use best-fit parameters to calculate desired  $\sigma_{\alpha\gamma}$

**Result:** Experimentally constrained cross section.

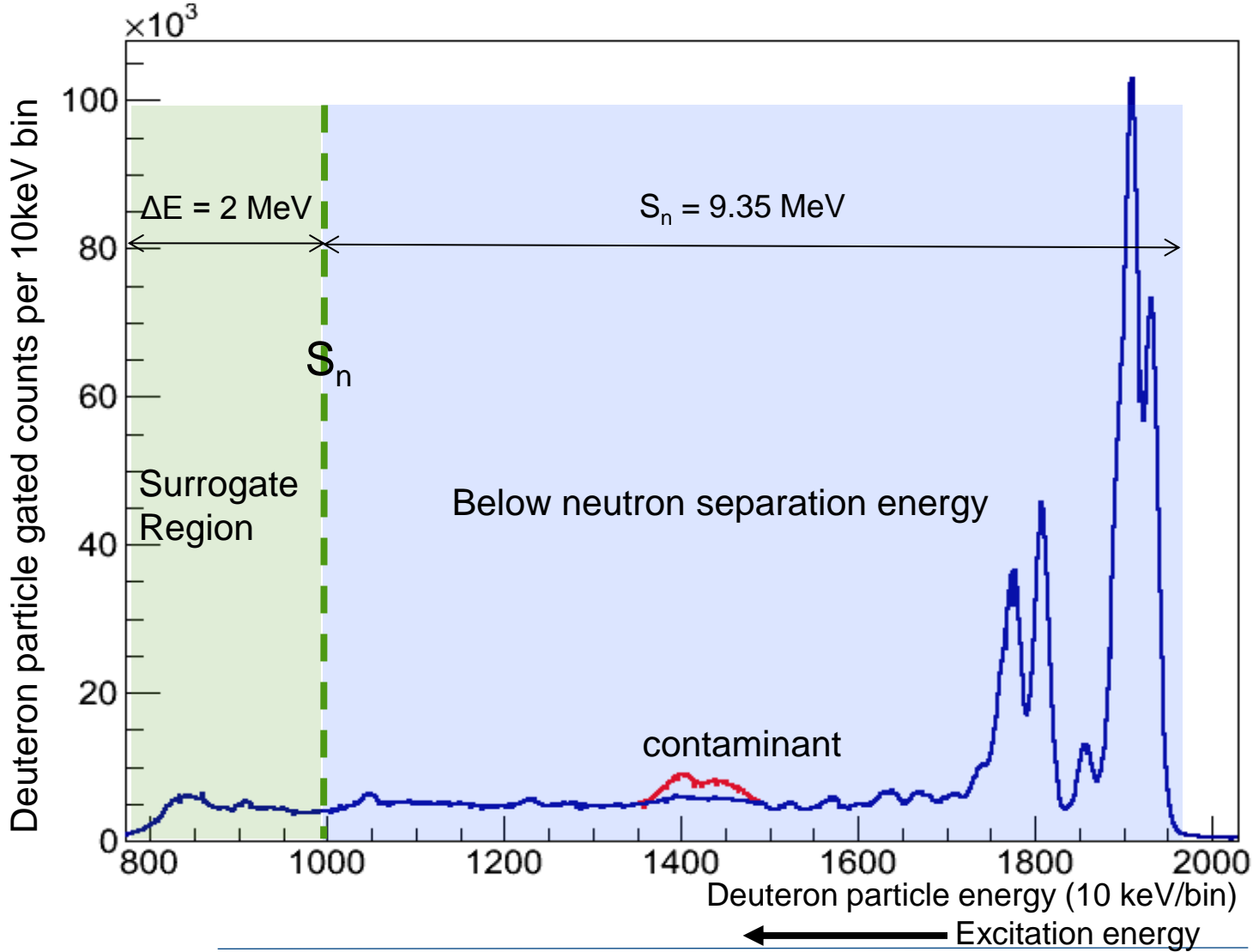
# Yttrium surrogate nuclear reactions

## Reactions populate Y89, Y88, Y87 nuclei



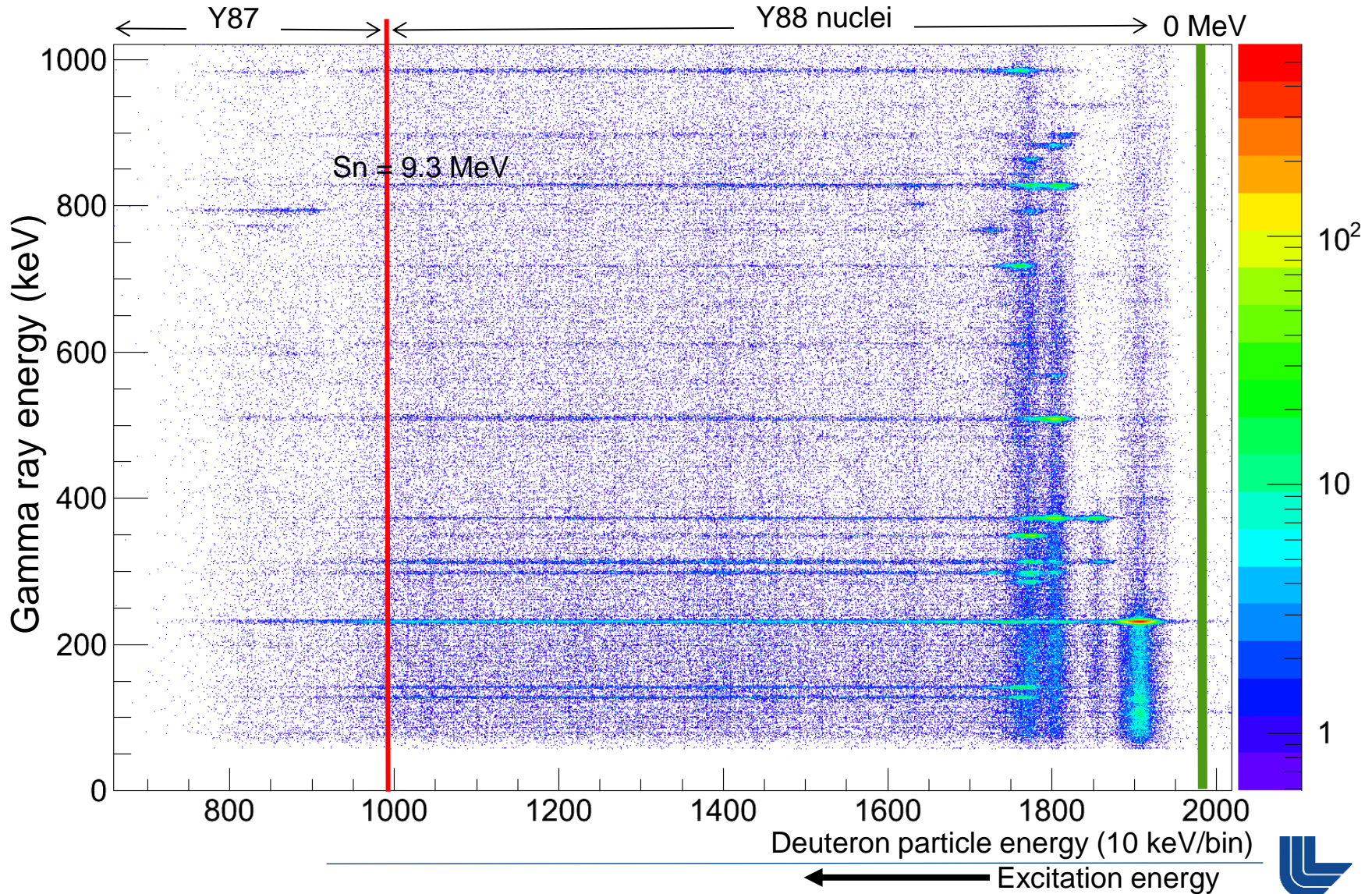
Particle singles data used to determine  $^{89}\text{Y}(n,\gamma)$  using the  $^{89}\text{Y}(p,d)^{88}\text{Y}^*$  reaction

$$P_{\delta\gamma}(E) = \frac{N_{\delta\gamma}(E_{\text{exc}})}{N_{\delta}(E_{\text{exc}})}$$



# $^{89}\text{Y}(p,d)^{88}\text{Y}$ Gamma production versus excitation energy

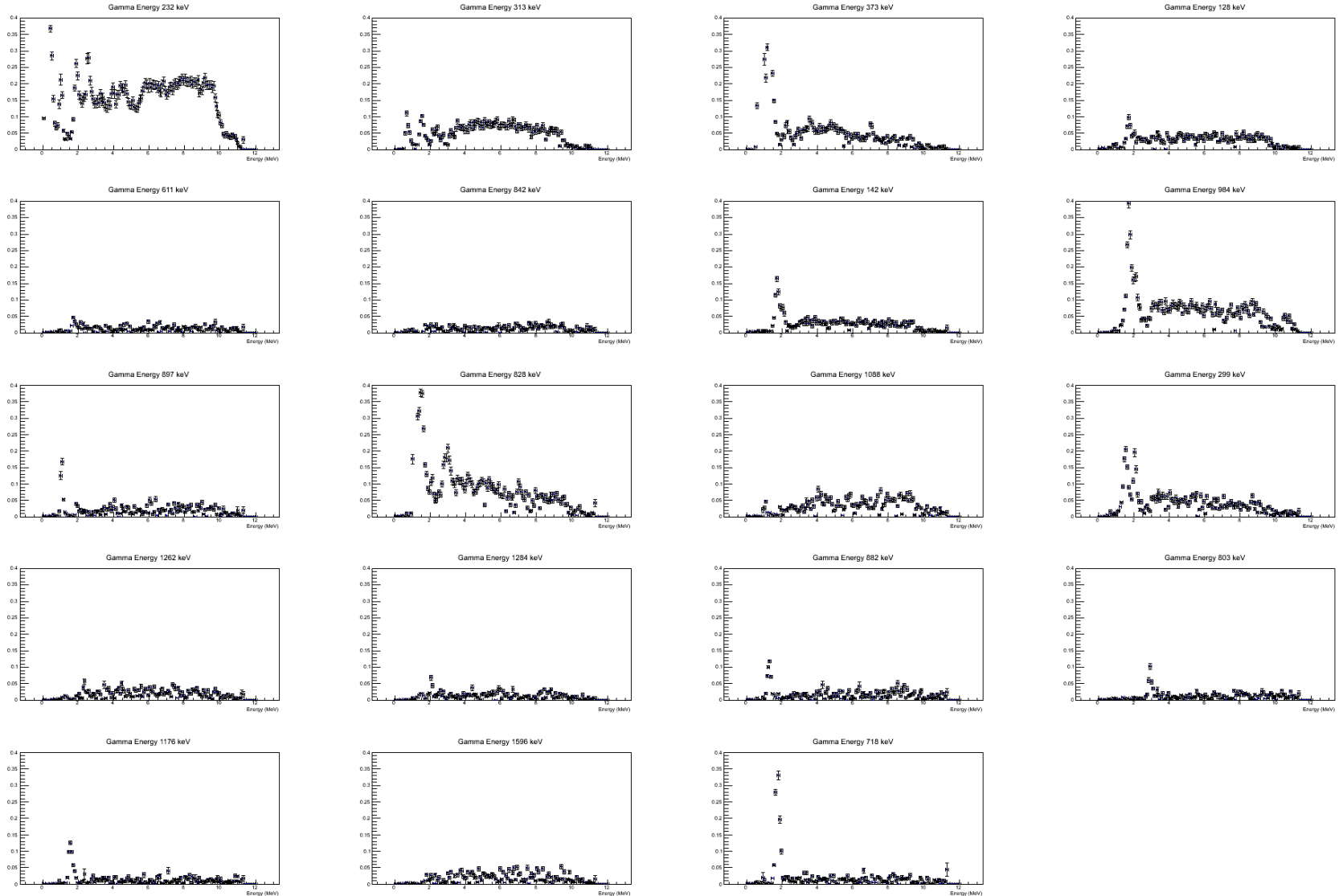
$$P_{\delta\gamma}(E) = \frac{N_{\delta\gamma}(E_{\text{exc}})}{N_{\delta}(E_{\text{exc}})}$$



# $^{89}\text{Y}(p,d)^{88}\text{Y}$ gamma decay probabilities

Many contributions used to determine the final parameters

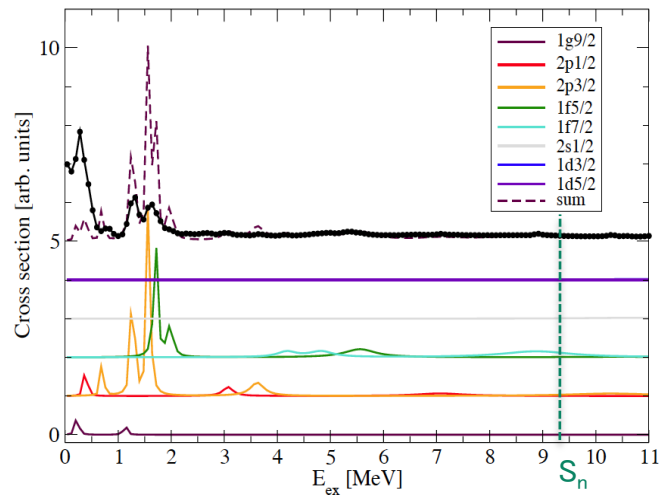
$$P_{\delta\gamma}(E) = \frac{N_{\delta\gamma}(E_{\text{exc}})}{N_{\delta}(E_{\text{exc}})}$$



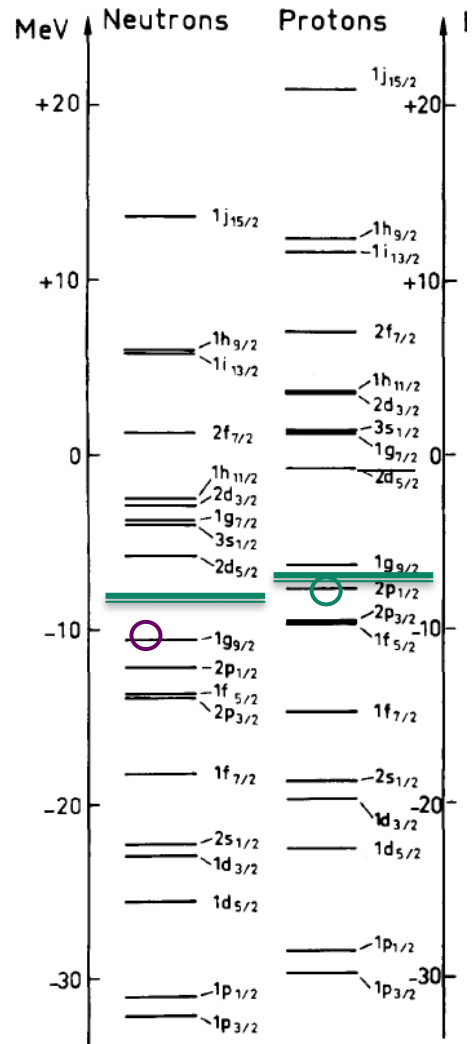
# From measurement to cross section: Theory step I

$$P_{\chi}(E) = \sum_{J,\pi} F_{\delta}^{CN}(E,J,\pi) \cdot G_{\chi}^{CN}(E,J,\pi)$$

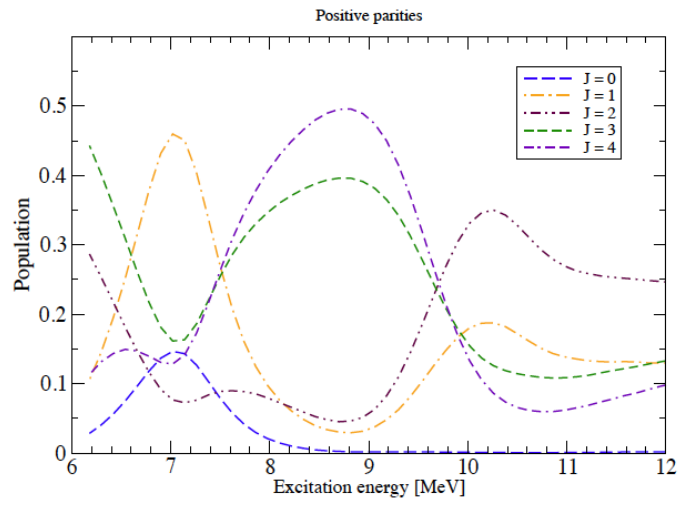
- Theoretical description of  $^{89}\text{Y}(p,d)^{88}\text{Y}^*$



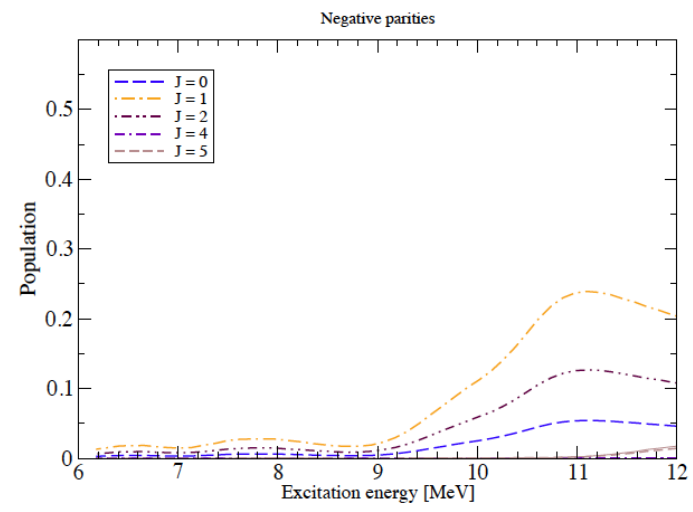
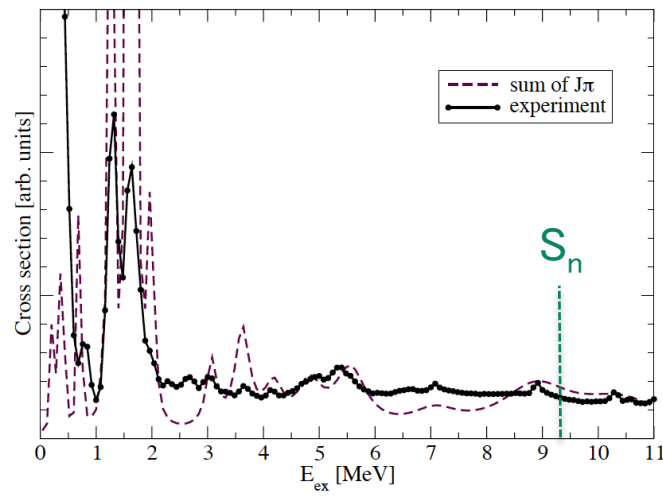
$^{89}\text{Y}$



- Extract spins, determine  $J^{\pi}$  distribution of  $^{88}\text{Y}$  as function of E.



- Closer comparison with experiment:



(do for proton holes  $2p_{1/2}$ ,  $2p_{3/2}$ ,  $1g_{9/2}$ )

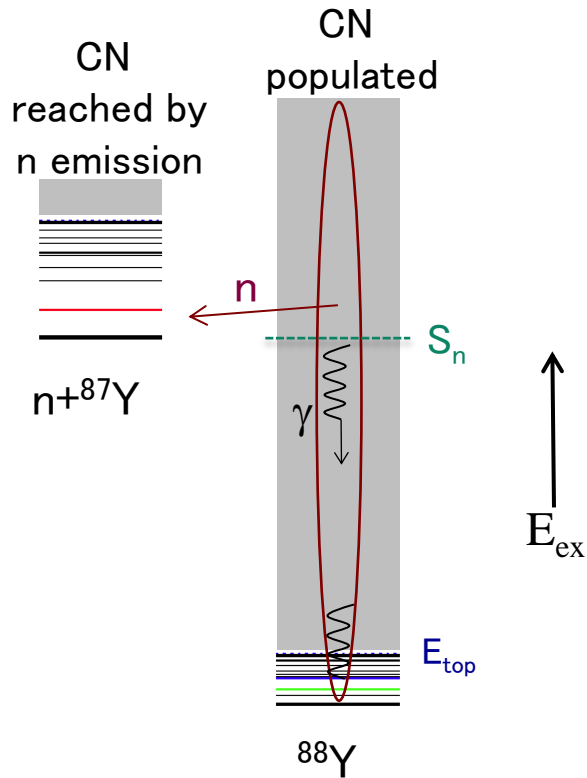
○ neutron hole  
made in reaction



# From measurement to cross section: Theory step II

## Develop rough decay model for $^{88}\text{Y}^*$ :

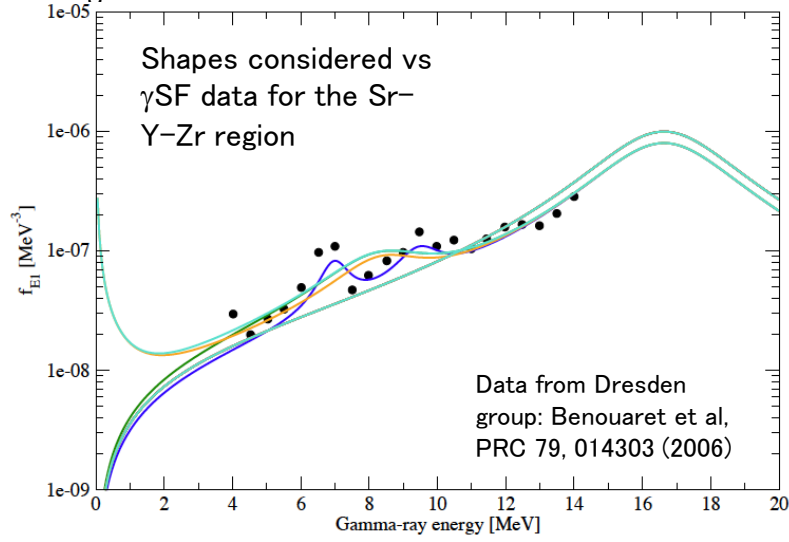
- Literature values for energies + spins of low-lying states
- Literature values for branchings between low states
- LDs of Gilbert-Cameron form w/variable parameters



measurements &

### Example $\gamma\text{SF}$ :

Considering 5 shapes, w and w/o upturn and/or extra strength around 8 MeV.



Decay model gives  $G_{\chi}^{\text{CN}}(E, J, \pi)$  for range of LD and  $\gamma\text{SF}$  parameters.

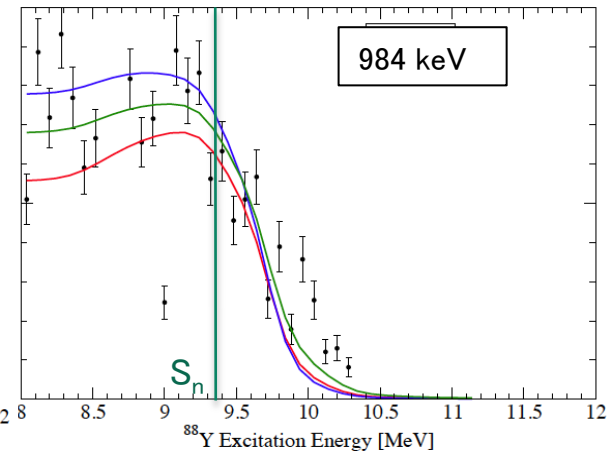
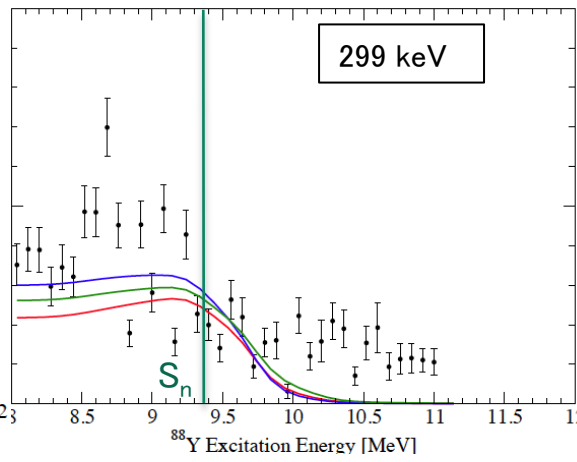
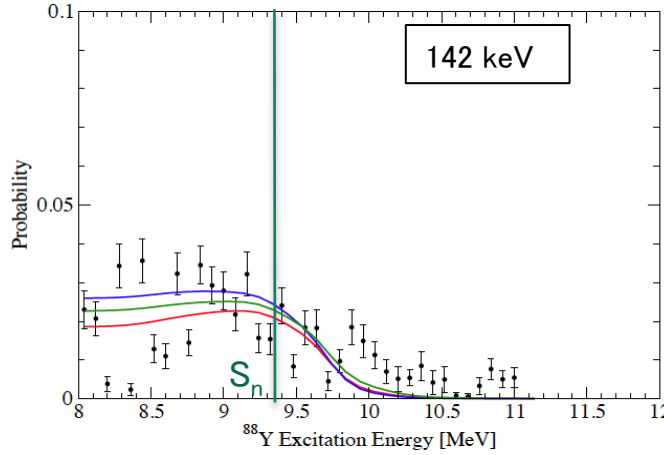
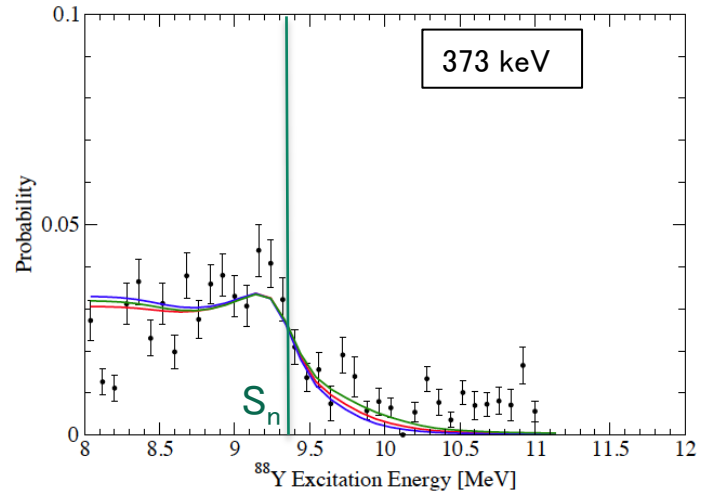
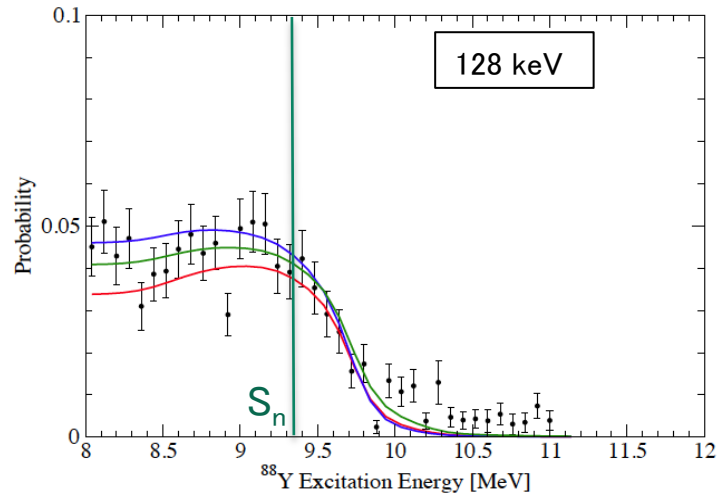
$$P_{\chi}(E) = \sum_{J, \pi} F_{\delta}^{\text{CN}}(E, J, \pi) \cdot G_{\chi}^{\text{CN}}(E, J, \pi)$$

### Getting ready for fitting:

- Study of sensitivity to HF inputs was used to isolate relevant parameters for fit.
- Checks of  $\gamma$  decay against surrogate data used to identify incorrect branchings, etc.

# From measurement to cross section: Theory step III

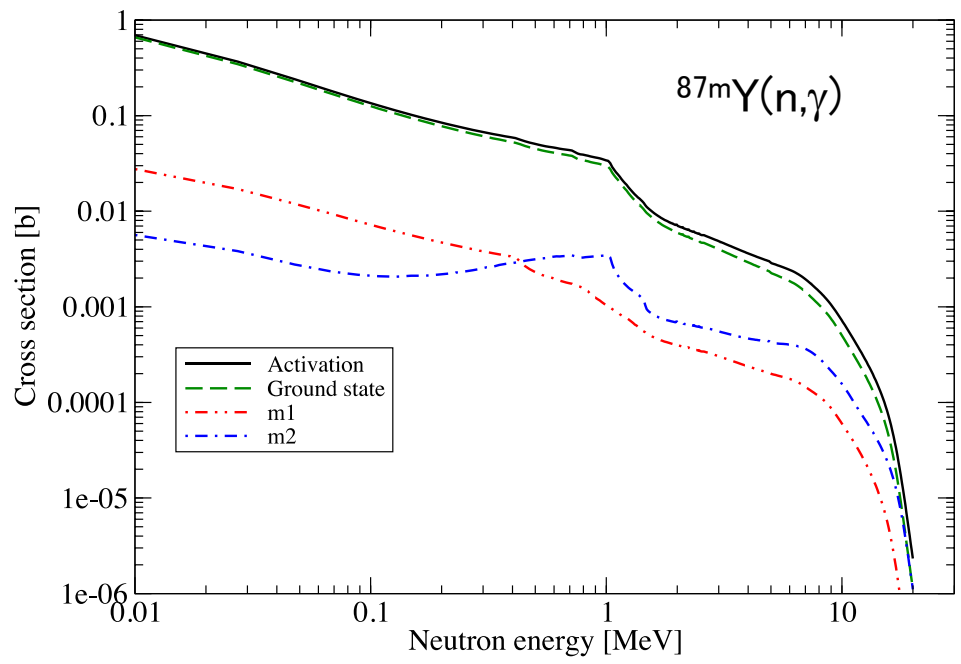
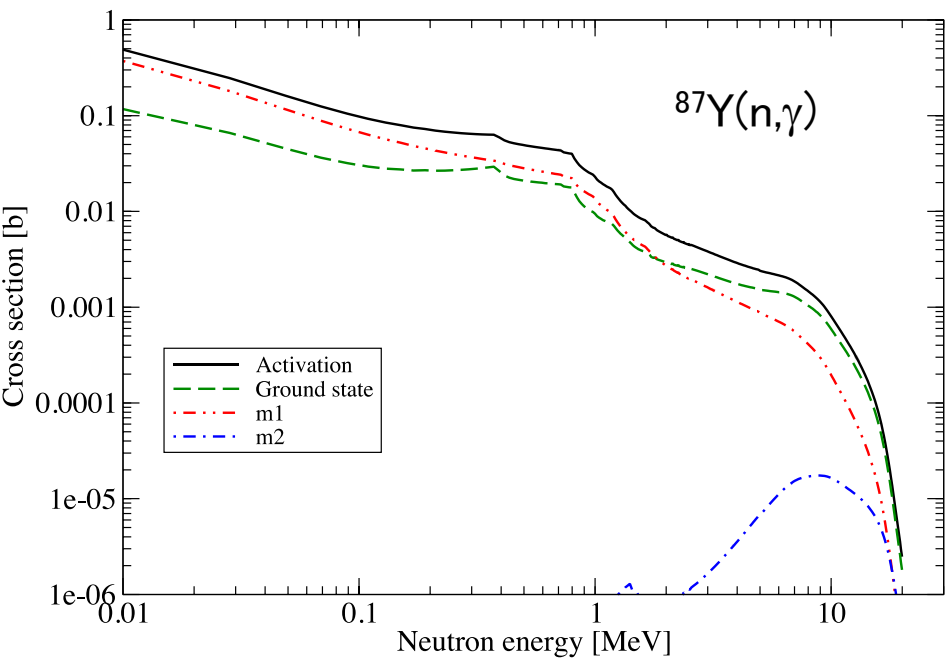
Fitting HF inputs to reproduce surrogate observables



Fit yields 'best' set of parameters & uncertainty estimate.

$$P_{\chi}(E) = \sum_{J,\pi} F_{\delta}^{CN}(E,J,\pi) \cdot G_{\chi}^{CN}(E,J,\pi)$$

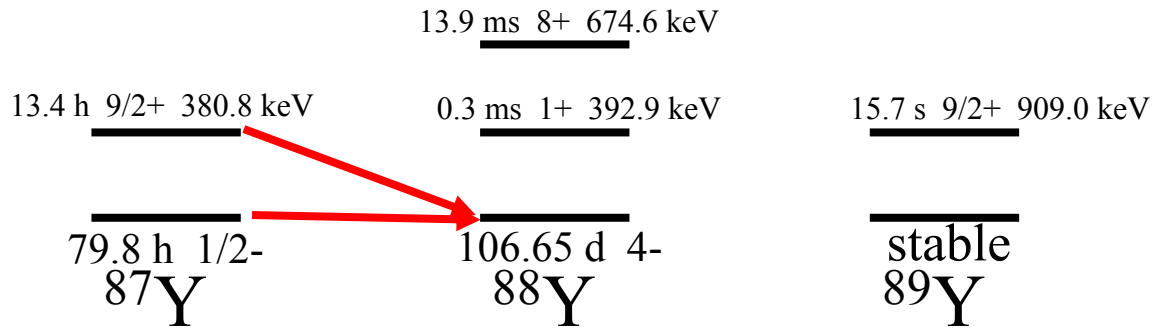
# From measurement to cross section: Theory step IV



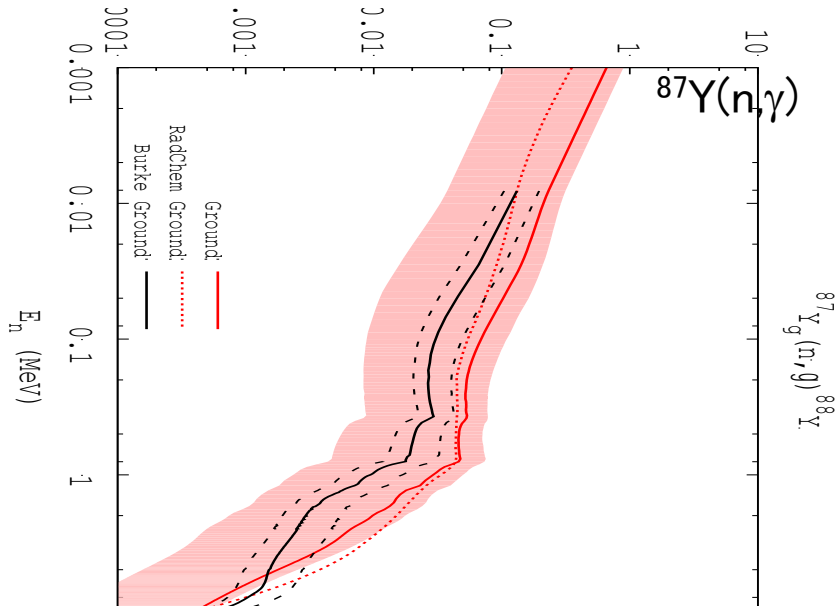
Using 'best' set of parameters to calculate  $^{87}\text{Y}(n,\gamma)$  and  $^{87\text{m}}\text{Y}(n,\gamma)$ , including cross sections for gs and isomers in  $^{88}\text{Y}$

$$\sigma_{\alpha\chi} = \sum_{J,\pi} \sigma_{\alpha}^{\text{CN}}(E,J,\pi) \cdot G_{\chi}^{\text{CN}}(E,J,\pi)$$

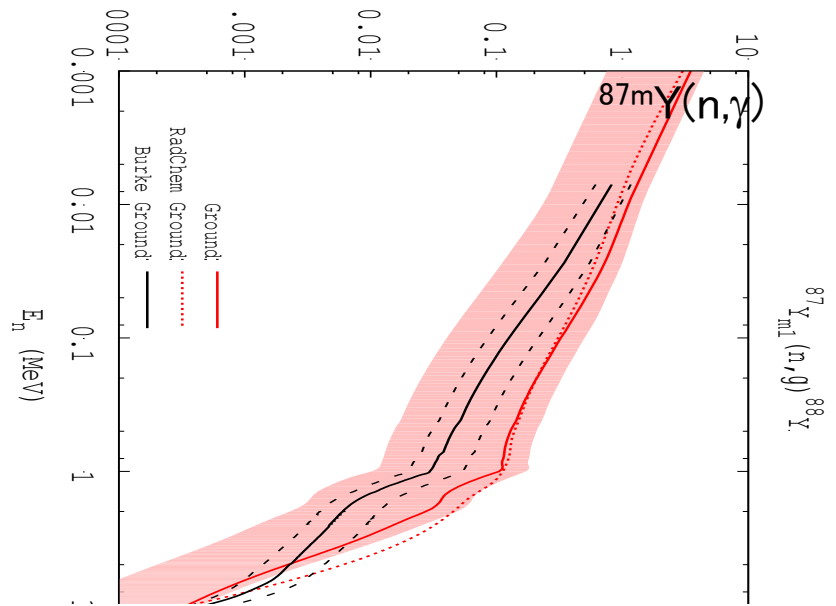
# Results – to GS



$^{87}\text{Y}(n,\gamma)$  to  $^{88}\text{Y}$  ground state ( $4^-$  at 0 keV)



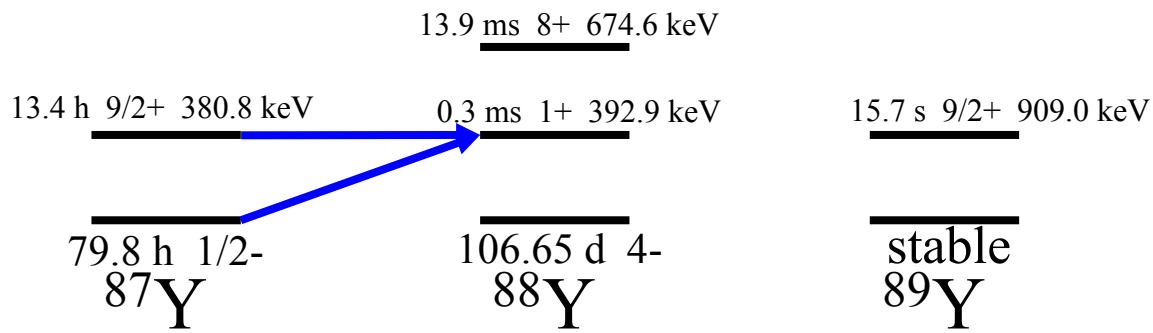
Target  $^{87}\text{Y}$  in gs ( $1/2^-$  at 0 keV)



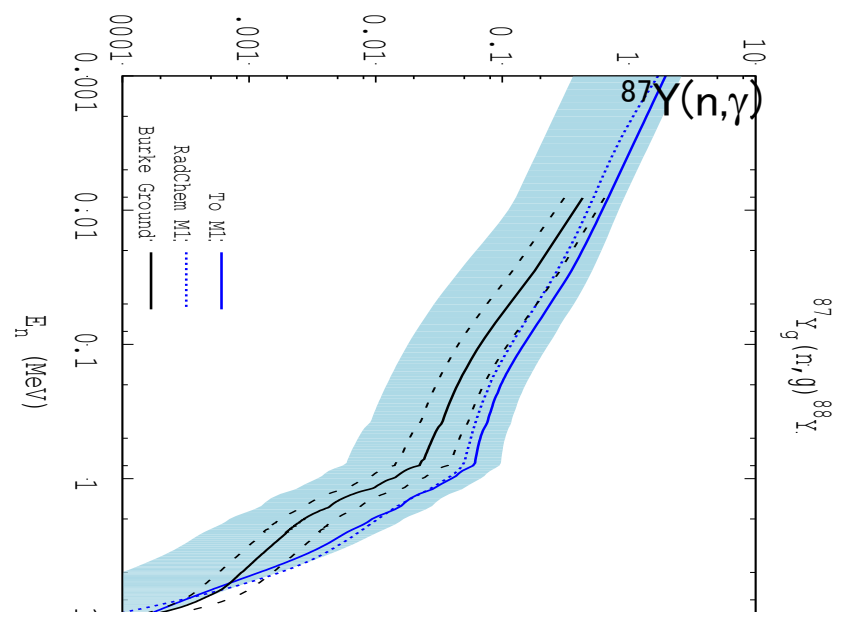
Target  $^{87}\text{Y}$  in M1 ( $9/2^+$  at 381 keV)

Plots courtesy of R.D. Hoffman

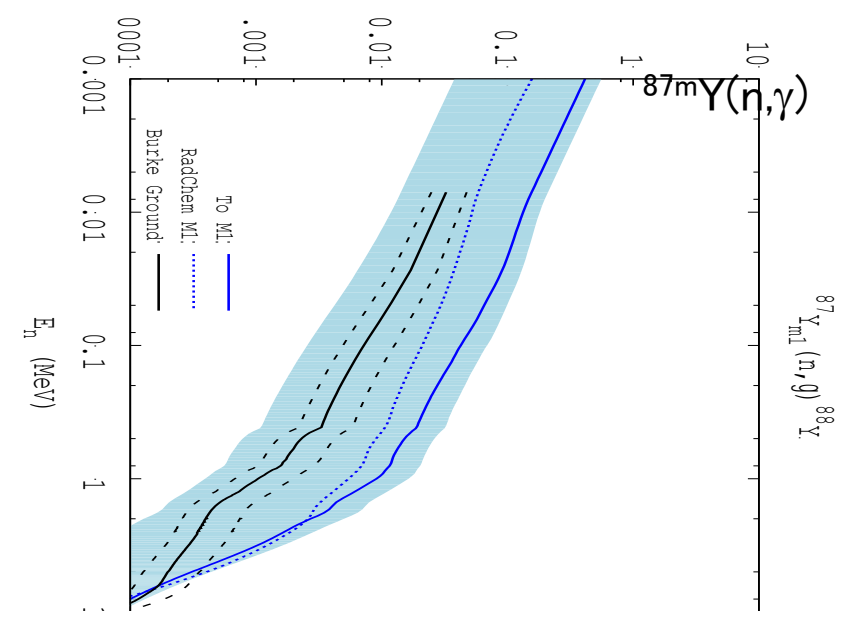
# Results – to M1



## $^{87}\text{Y}(n,\gamma)$ to $^{88}\text{Y}$ M1 isomer (1+ at 393 keV)



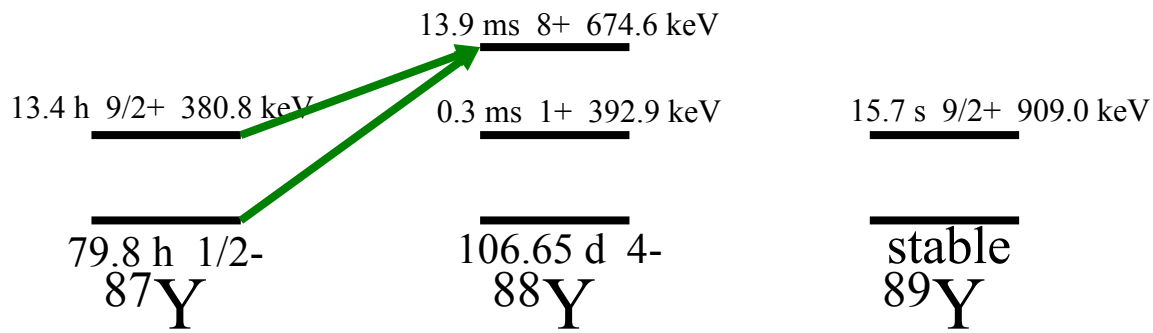
Target  $^{87}\text{Y}$  in gs (1/2- at 0 keV)



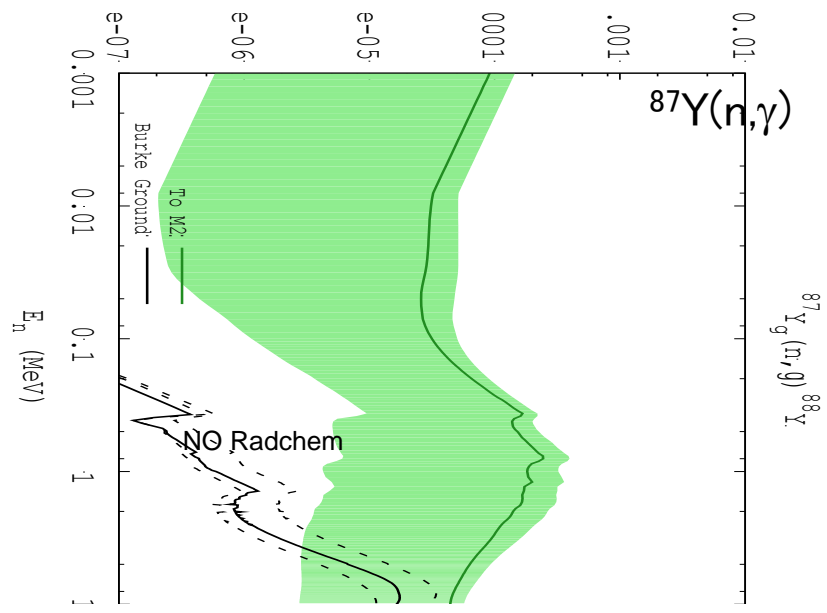
Target  $^{87}\text{Y}$  in M1 (9/2+ at 381 keV)

Plots courtesy of R.D. Hoffman

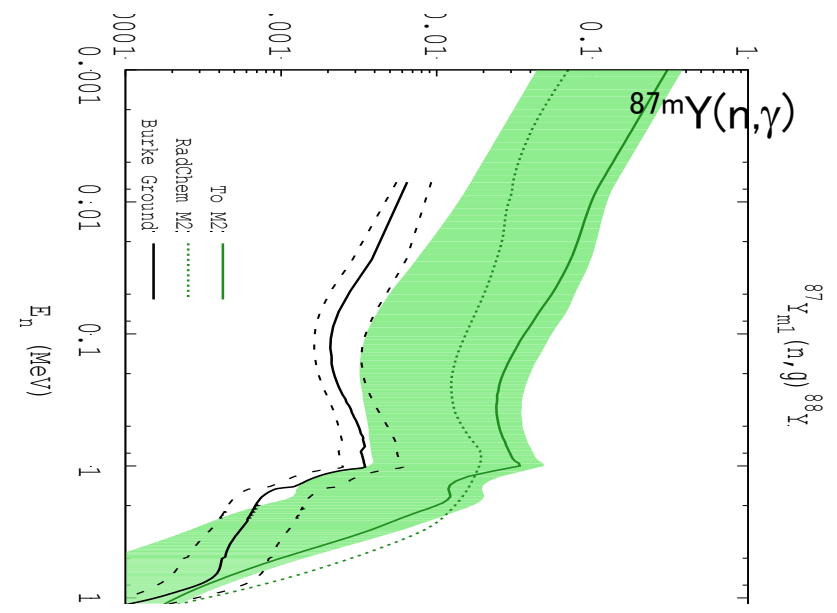
# Results - to M2



$^{87}\text{Y}(n,\gamma)$  to  $^{88}\text{Y}$  M2 isomer (8+ at 675 keV)



Target  $^{87}\text{Y}$  in gs (1/2- at 0 keV)



Target  $^{87}\text{Y}$  in M1 (9/2+ at 381 keV)

Plots courtesy of R.D. Hoffman

# Concerns we are addressing

- I. Improve nuclear structure in  $^{88}\text{Y}$  and  $^{87}\text{Y}$ :**
  - Unknown or incorrect spins & parities in literature
  - Poorly known or incorrect  $\gamma$  branching ratios in literature
  - Improvements through more detailed study of present data, coupled with theory: e.g. compare angular distributions of singles to fix spins; study  $\gamma\gamma$  coincidences to fix branching ratios
  
- II. Improve calculations for  $^{89}\text{Y}(p,d)$ :**
  - Use results of improvements above to revisit predicted  $J^\pi$  distributions
  
- III. Improved experimental analysis:**
  - Revisit experimental coincidence probabilities to reduce data scatter
  
- IV. Improved fitting of theory to experimental results:**
  - Use improved experimental coincidence probabilities in fits
  - Carry out more detailed parameter variation in fit
  - Carry out more detailed uncertainty analysis
  
- V. High-E region:**
  - Include direct-semidirect cross section contributions
  
- VI. Benchmark:**
  - Use approach to determine known  $^{90}\text{Zr}(n,\gamma)$  cross section: underway
  - $^{91}\text{Zr}(p,d)$ ,  $^{92}\text{Zr}(p,d)$ ,  $^{94}\text{Zr}(p,d)$ ,  $^{96}\text{Zr}(p,d)$  were also measured



# Summary

1) We have measured the following cross sections using the surrogate ratio method in the past couple of years.

- $^{236}\text{Pu}(n,f)$
- $^{237}\text{Pu}(n,f)$
- $^{238}\text{Pu}(n,f)$
- $^{240}\text{Am}(n,f)$
- $^{241}\text{Am}(n,f)$
- $^{242}\text{Am}(n,f)$

2) We have a path forward and preliminary results for measuring (n,g) reactions for ground states and isomeric states for spherical nuclei

$^{87}\text{Y}(n,g)$  GS  $\rightarrow$  GS

$^{87}\text{Y}(n,g)$  M1  $\rightarrow$  GS

$^{87}\text{Y}(n,g)$  GS  $\rightarrow$  M1

$^{87}\text{Y}(n,g)$  M1  $\rightarrow$  M1

$^{87}\text{Y}(n,g)$  GS  $\rightarrow$  M2

$^{87}\text{Y}(n,g)$  M1  $\rightarrow$  M2

3) **The burden of proof is on us.** We are benchmarking the surrogate (n,g) using  $^{90}\text{Zr}$ ,  $^{91}\text{Zr}$ ,  $^{92}\text{Zr}$ ,  $^{94}\text{Zr}$ ,  $^{96}\text{Zr}$  (p,d) data from the same experiment.  $^{90}\text{Zr}(n,g)$  is the key benchmark from the surrogate reaction  $^{92}\text{Zr}(p,d)$ .





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# Yttrium surrogate nuclear reactions

