

54<sup>th</sup> ASRC International Workshop Sakura-2019:  
Nuclear Fission and Structure of Exotic Nuclei

JAEA, Tokai, Japan, 25-27 March 2019

*Fission studies with microscopic level densities  
I: Fragment mass distributions and kinetic energies*

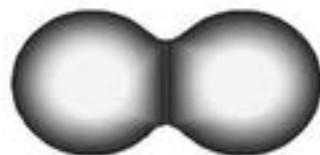
Jørgen Randrup, Lawrence Berkeley National Laboratory, California, USA

in collaboration with

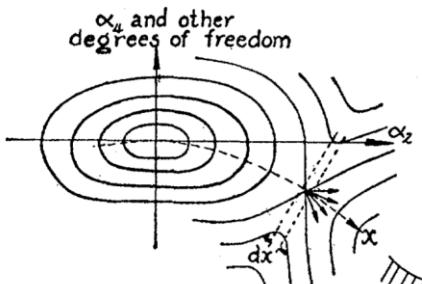
Martin Albertsson, Gillis Carlsson, David Ward, Sven Åberg, Mathematical Physics, Lund University, Sweden

Thomas Døssing, Niels Bohr Institute, Copenhagen University, Denmark

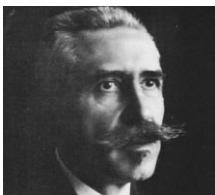
Peter Möller, Los Alamos National Laboratory, New Mexico, USA



# Fission shape dynamics



Bohr & Wheeler 1939

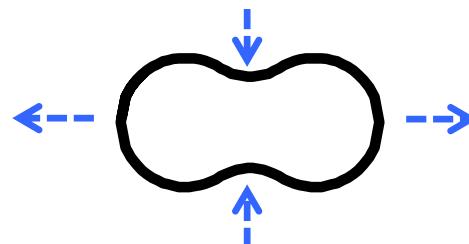


Paul Langevin  
(1872-1946)



Marian Smoluchowski  
(1872 – 1917)

Time evolution of the nuclear shape parameters,  $\chi(t)$ :



3QS (Nix 1969):  
Elongation  $Q_2$   
Neck radius  $c$   
Defs  $\varepsilon_{f1}$  &  $\varepsilon_{f2}$   
Asymmetry  $\alpha$

Langevin equation:  $d\pi/dt = \mathbf{F}^{cons} + \mathbf{F}^{diss}$

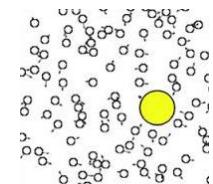
$M(\chi)$	$U(\chi)$	$\gamma(\chi)$
inertial tensor	potential energy	dissipation tensor

The shape motion is *highly dissipative* => Brownian motion

Smoluchowski equation:  $\mathbf{0} = \mathbf{F}^{cons} + \mathbf{F}^{diss}$

[1<sup>st</sup> application: Kramers 1940]

$U(\chi)$	$\gamma(\chi)$
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*Phase space dominance:* Metropolis walk on  $U(\chi)$

Macroscopic:

Wall formula

$$\dot{Q} = \rho \bar{v} \oint V^2 d\sigma$$

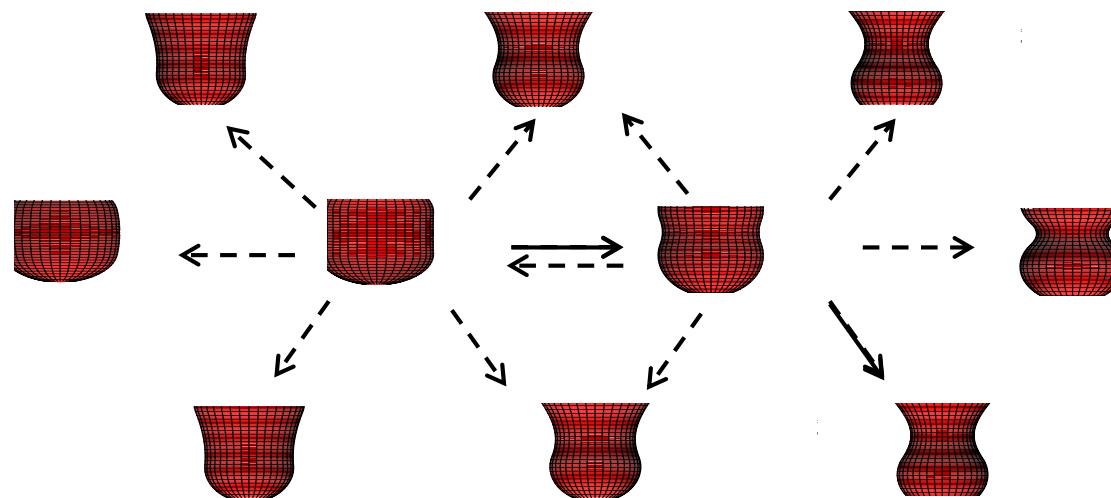
$$t_{\text{damp}} \approx 20-40 \text{ fm/c}$$

The nuclear *shape evolution* ...... is *highly dissipative* ...... so it is akin to *Brownian motion*

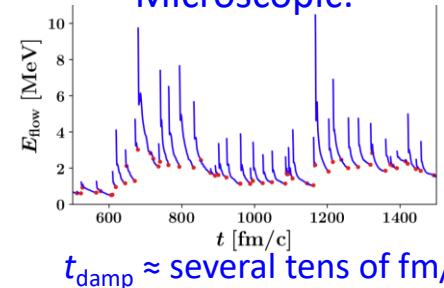
Phase space dominance:

$$\nu(\chi \rightarrow \chi^\square) \sim \rho(\chi^\square)$$

=&gt;

The fission evolution can be simulated by *random walks* among the shapes ...... generating a *ensemble of paths* through the space of shapes

Microscopic:



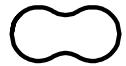
Metropolis:

$$P(\rho' \geq \rho) = 1$$

$$P(\rho' \leq \rho) = \rho'/\rho$$

# Calculation of the shape-dependent microscopic nuclear level density

H. Uhrenholt, S. Åberg, A. Dobrowolski, Th. Døssing, T. Ichikawa, P. Möller: NPA **913** (2013) 127



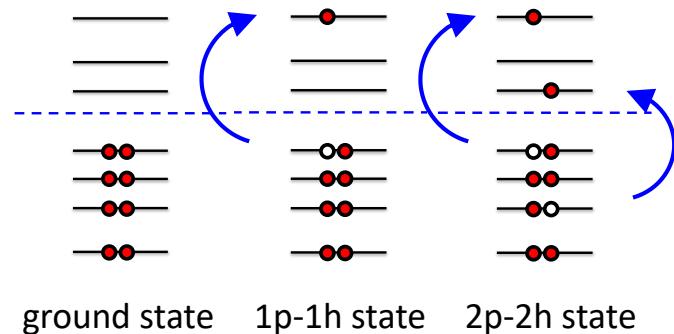
$$\rho(E, I, \pi) = \frac{1}{\Delta E} \int_E^{E+\Delta E} \sum_i \delta(E' - E_i(I, \pi)) dE' \quad \Rightarrow \quad \rho(E) = \sum_{I\pi} \rho(E, I, \pi)$$

$$E = E_p + E_n + E_{\text{rot}}$$

*Intrinsic states:*

Consider all multiple p-h excitations for protons and neutrons separately

$$|i\rangle = \prod_{\alpha=1}^n a_{\nu_\alpha}^+ a_{\nu'_\alpha} |0\rangle,$$



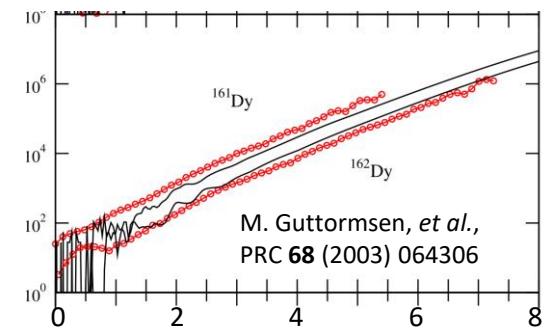
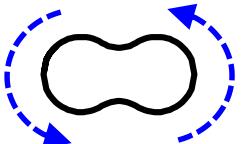
*Pairing*

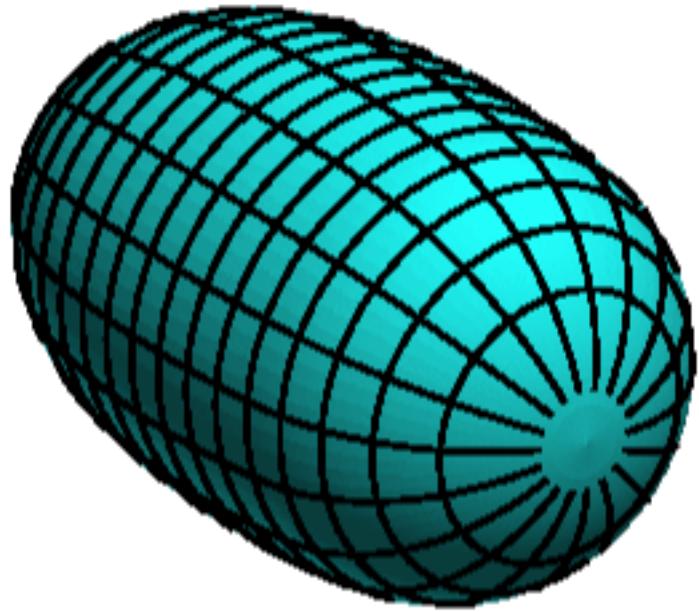
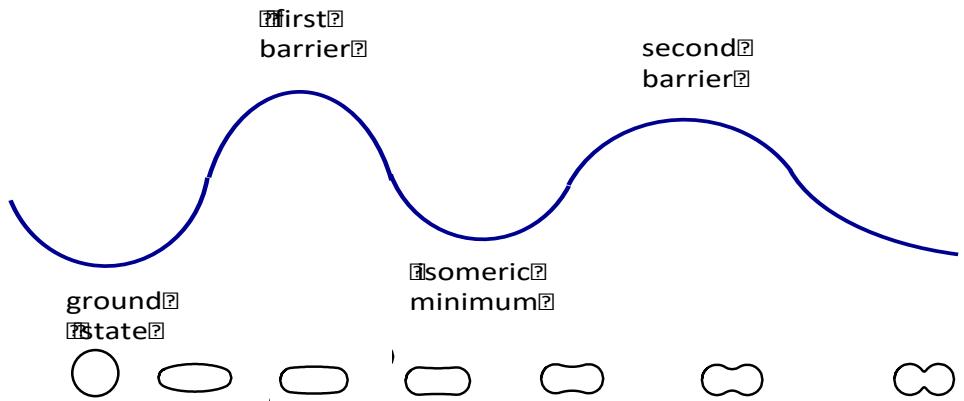
Calculate BCS pairing for each state:  $E_p$  &  $E_n$

*Rotation*

Rotational band built on each intrinsic state:  $E_{\text{rot}}$

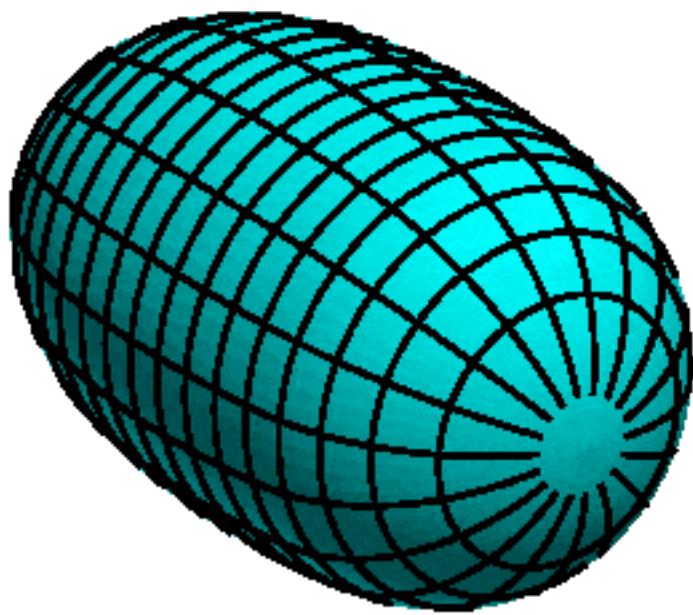
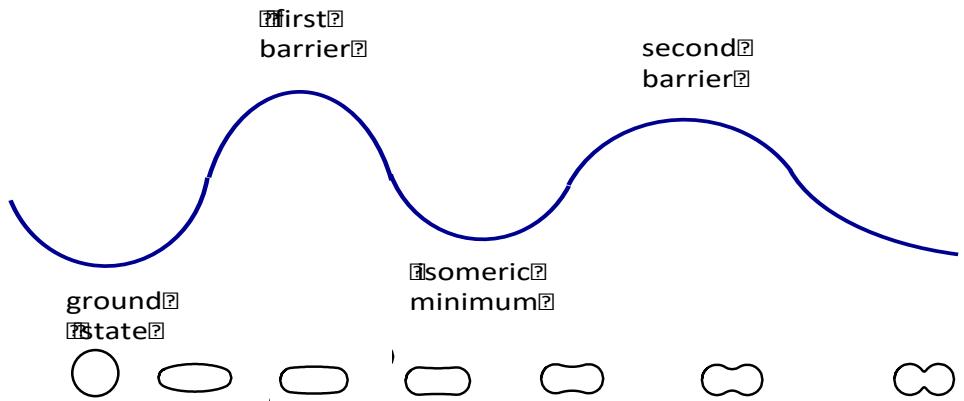
$$E_{\text{rot}}(I, K) = [I(I+1) - K^2]/2 \mathbb{I}_{\text{perp}}(\chi, \Delta_p, \Delta_n)$$





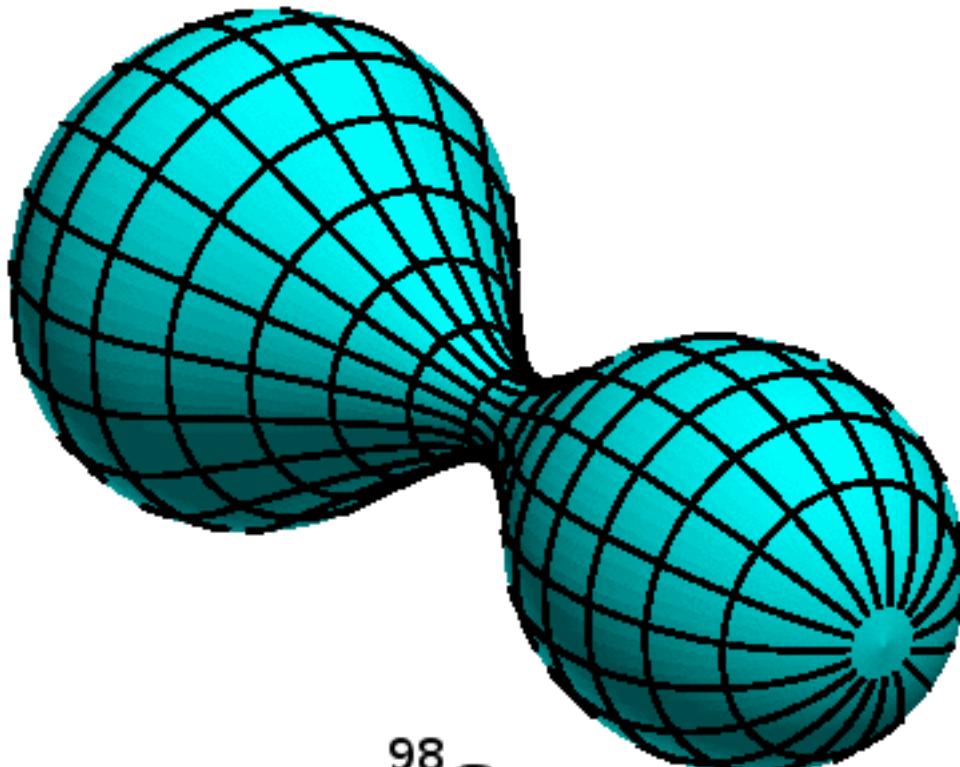
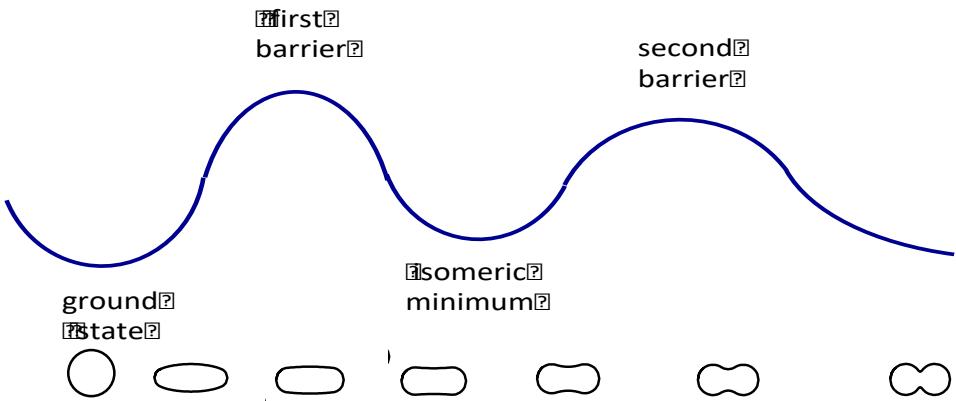
$$\nu(\chi \rightarrow \chi') \sim \rho(\chi')$$

Animation by Martin Albertsson



$$\nu(\chi \rightarrow \chi') \sim \rho(\chi')$$

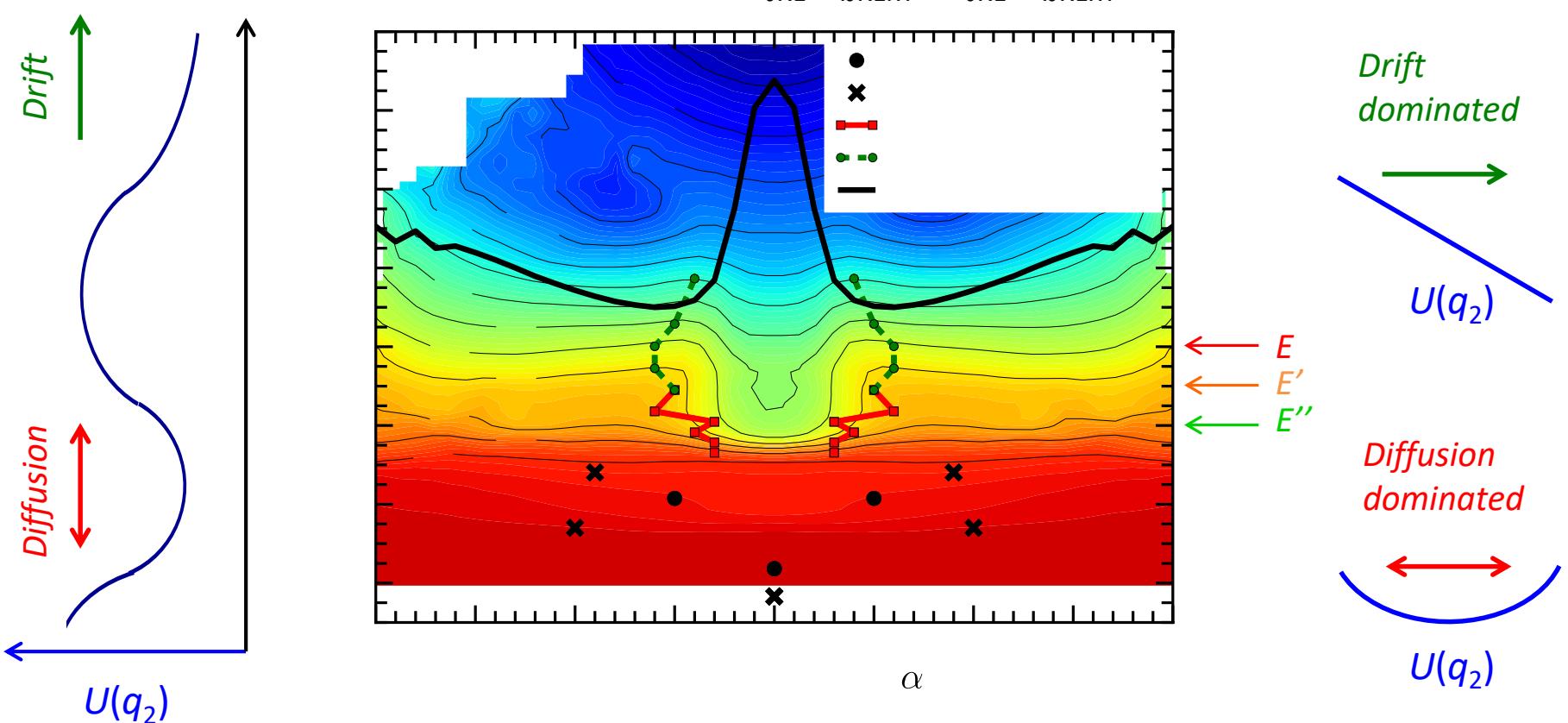
Animation by Martin Albertsson



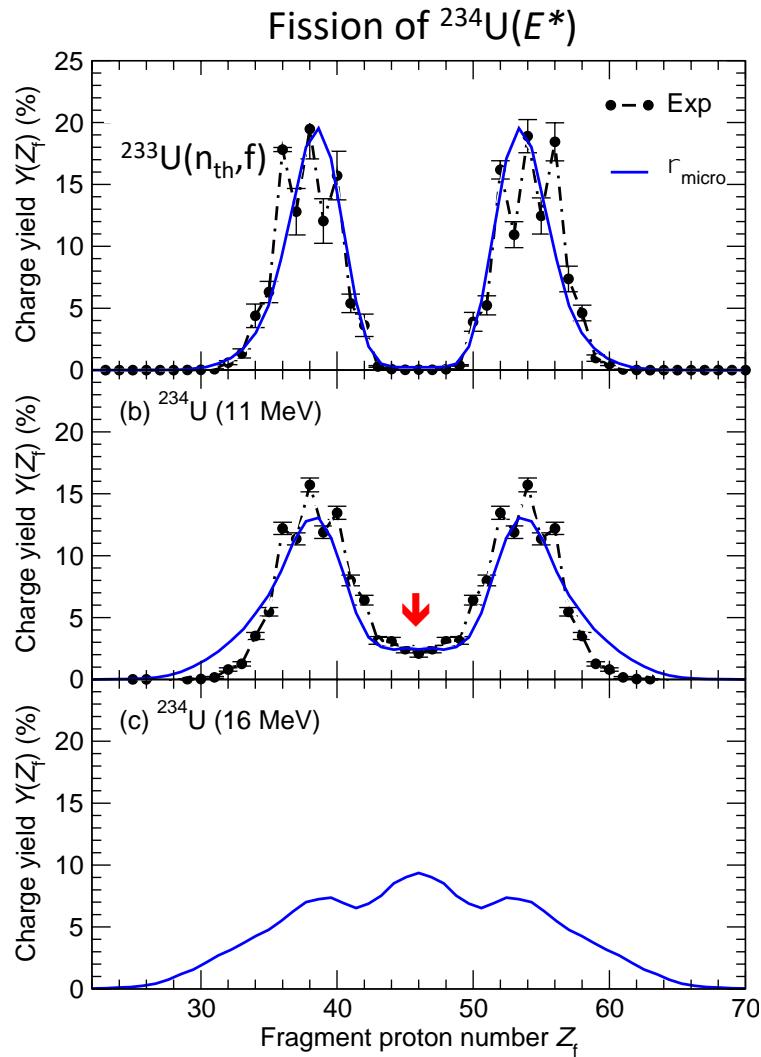
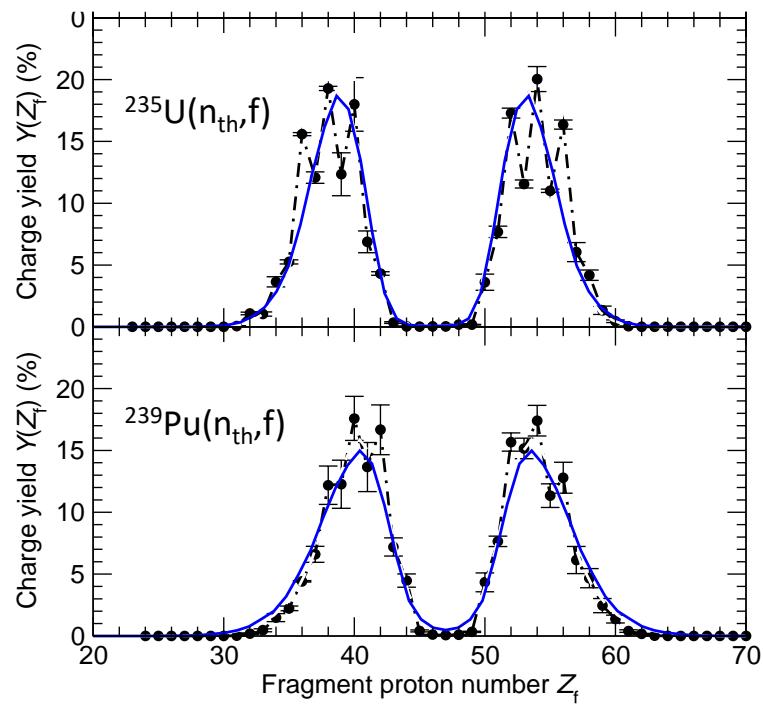
$$\nu(\chi \rightarrow \chi') \sim \rho(\chi')$$

Animation by Martin Albertsson

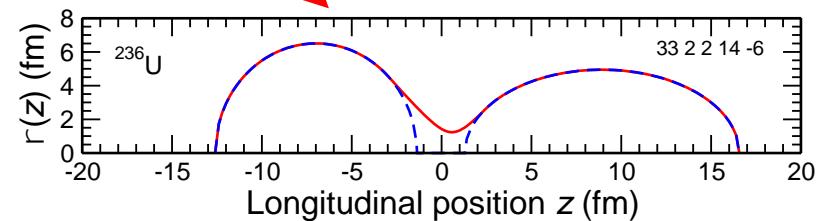
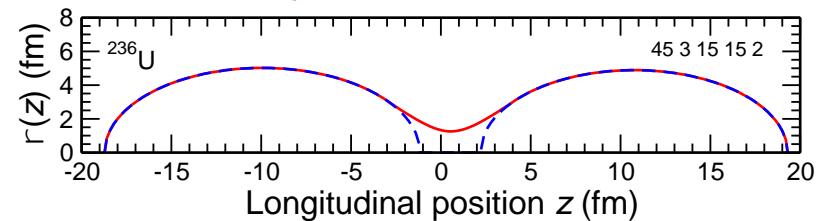
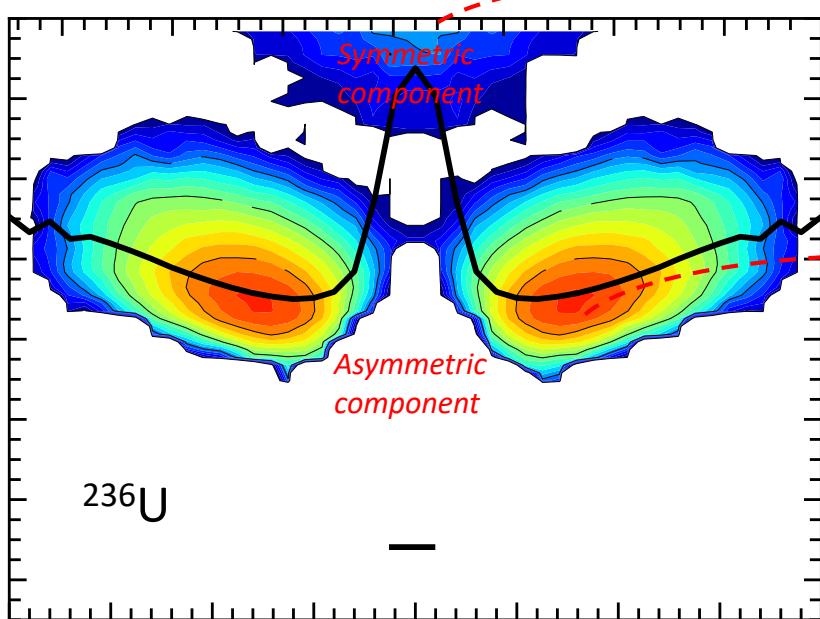
# Fission of $^{236}\text{U}$ : Equilibration during the random walk



# Mass yields using microscopic level densities

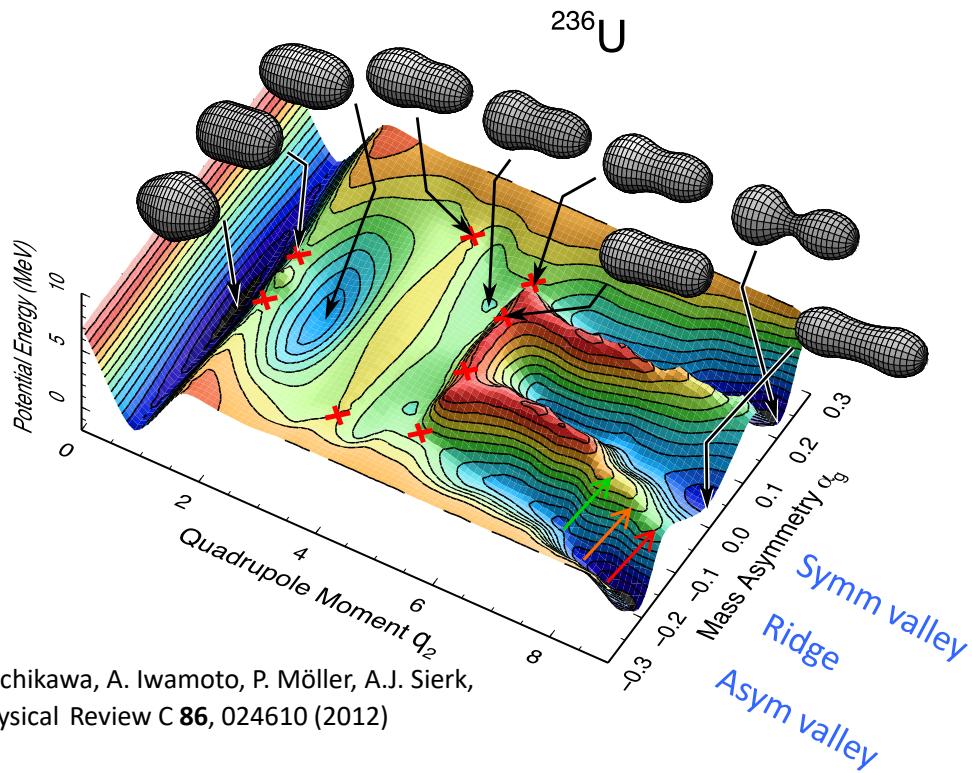


# Scission shapes ( $c_0 = 1.5$ fm)



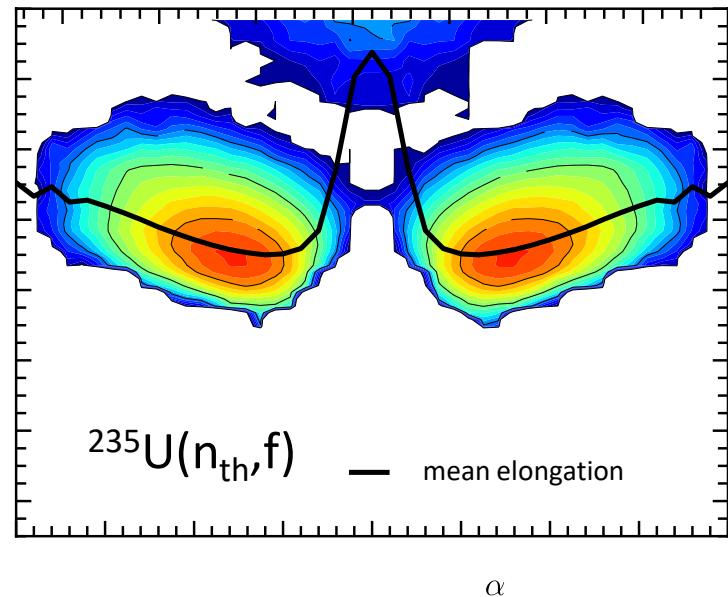
# Fission of $^{236}\text{U}$ :

Potential-energy surface



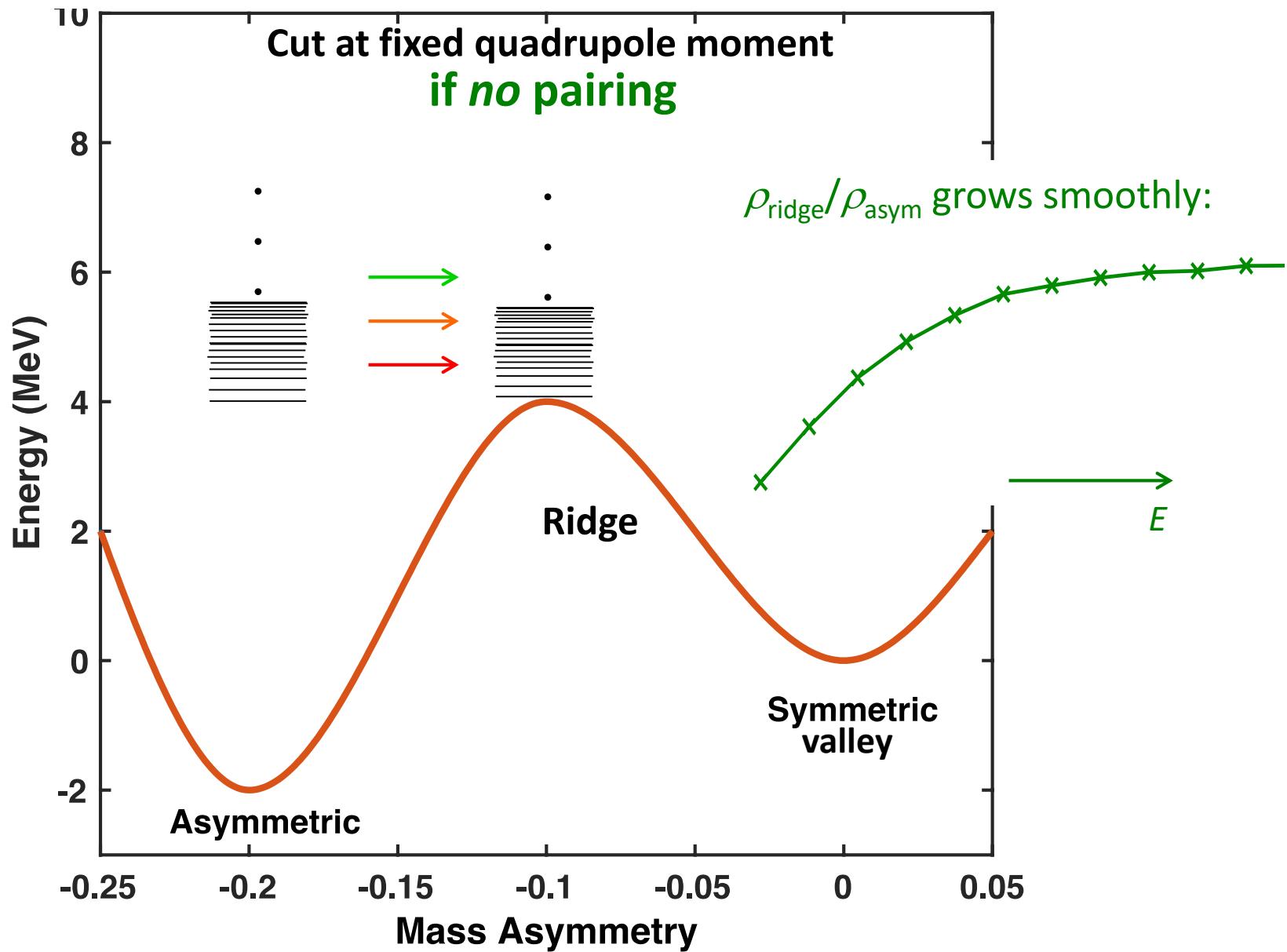
T. Ichikawa, A. Iwamoto, P. Möller, A.J. Sierk,  
Physical Review C **86**, 024610 (2012)

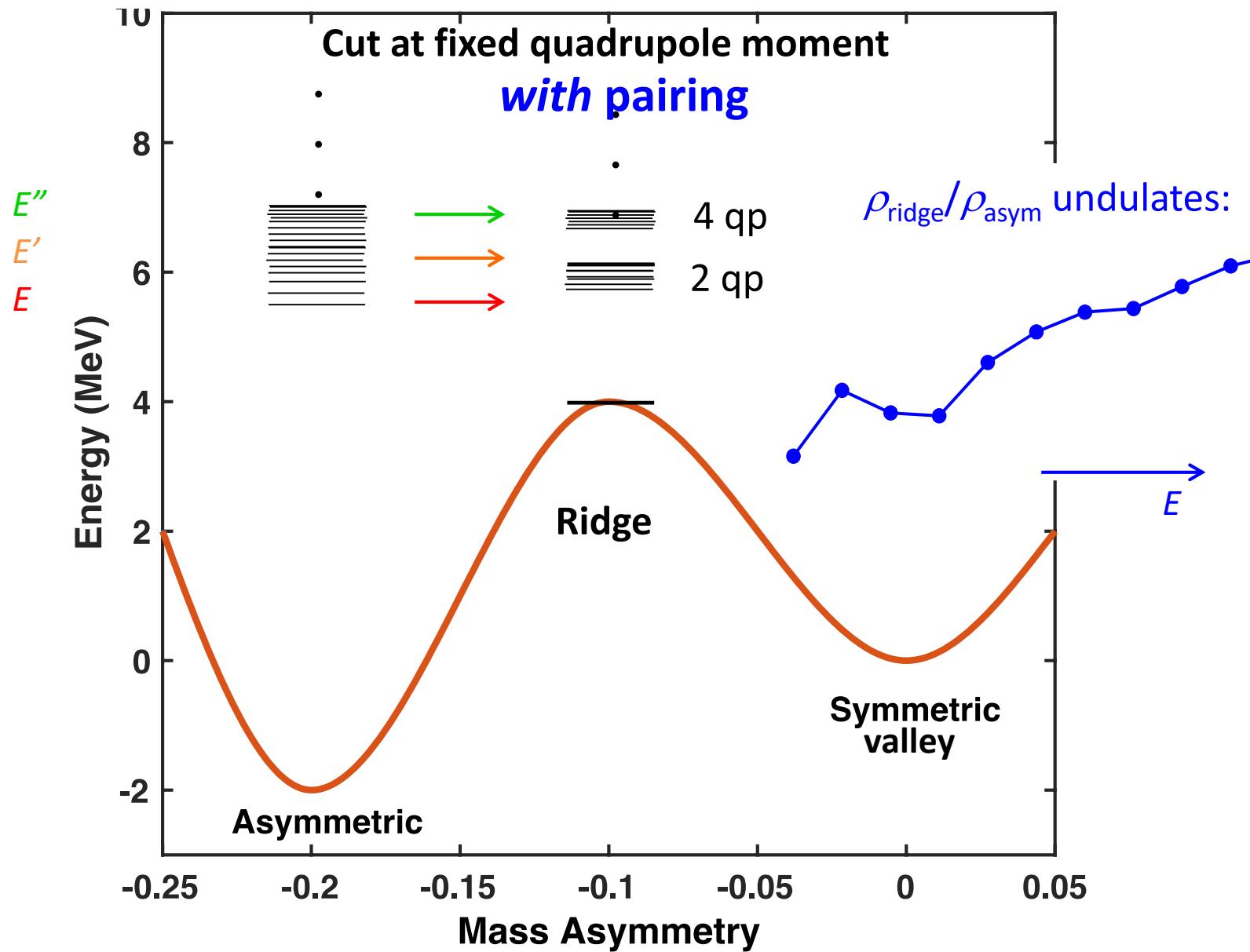
Scission shapes



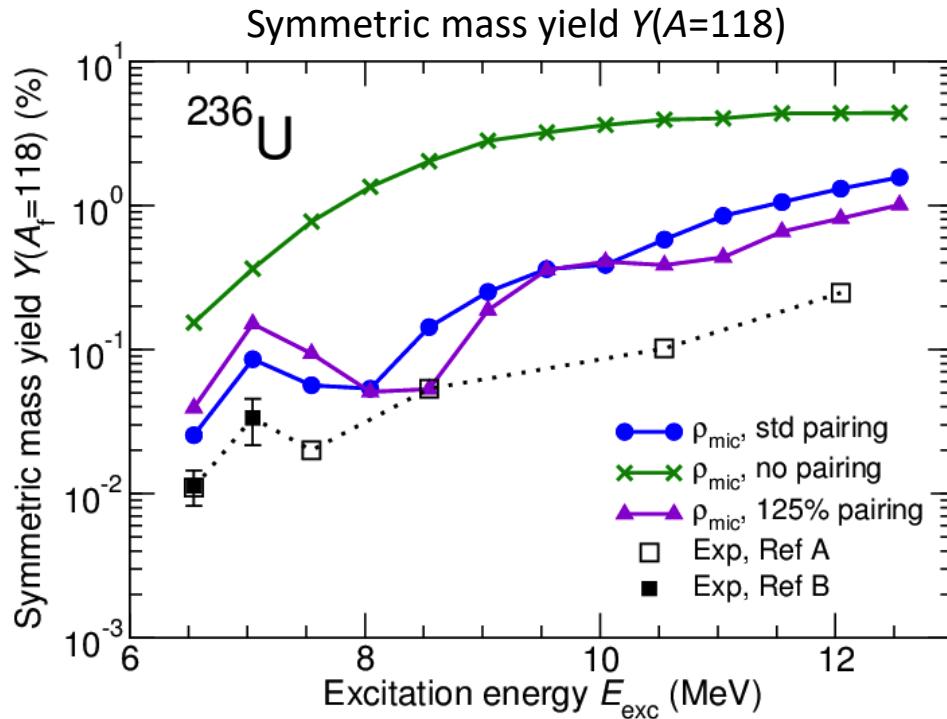
Shapes in the *valleys* have dilute single-particle spectra,  
hence negative shell energies and small pairing gaps

Shapes along the *ridge* have dense single-particle spectra,  
hence positive shell energies and large pairing gaps





# Energy dependence of fission yields: Non-monotonic behavior of the symmetric yield



The behavior is  
non-monotonic!

- A) L. E. Glendenin *et al.*, Phys. Rev. C **24**, 2600 (1981)
- B) M. B. Chadwick *et al.*, Nucl. Data Sheets **112**, 2887 (2011)
- C) M. E. Gooden *et al.*, Nucl. Data Sheets **131**, 319 (2016):  
*Non-monotonic* energy dependence of  $Y(A_f)$  from  $^{240}\text{Pu}^*$

OBS:

# Scission shapes

Extended  
shape  
lattice

Standard

→

$^{233}\text{U}(n_{\text{th}}, f)$

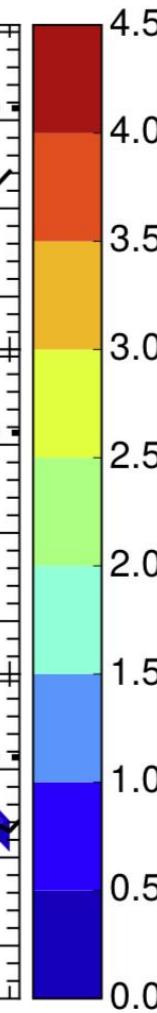
$c_0 = 1.5 \text{ fm}$

$c_0 = 2.5 \text{ fm}$

$^{235}\text{U}(n_{\text{th}}, f)$

Elongation  $q_2$

$^{239}\text{Pu}(n_{\text{th}}, f)$



(a)

(b)

(c)

(d)

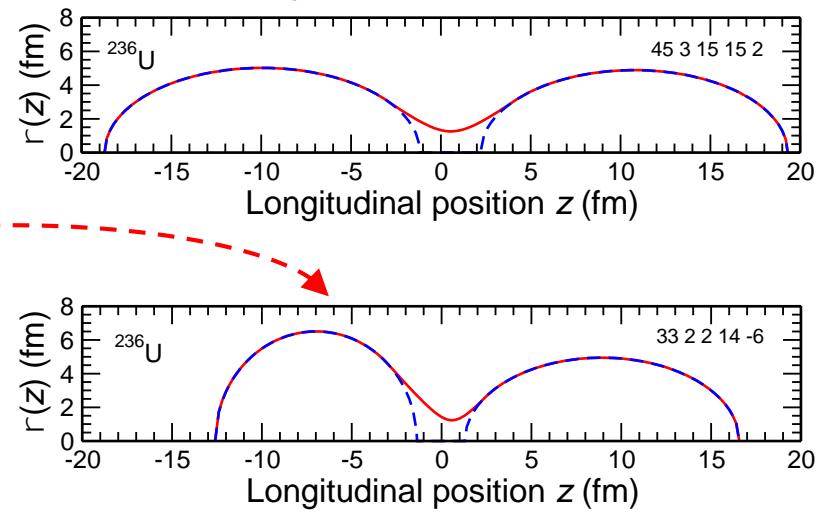
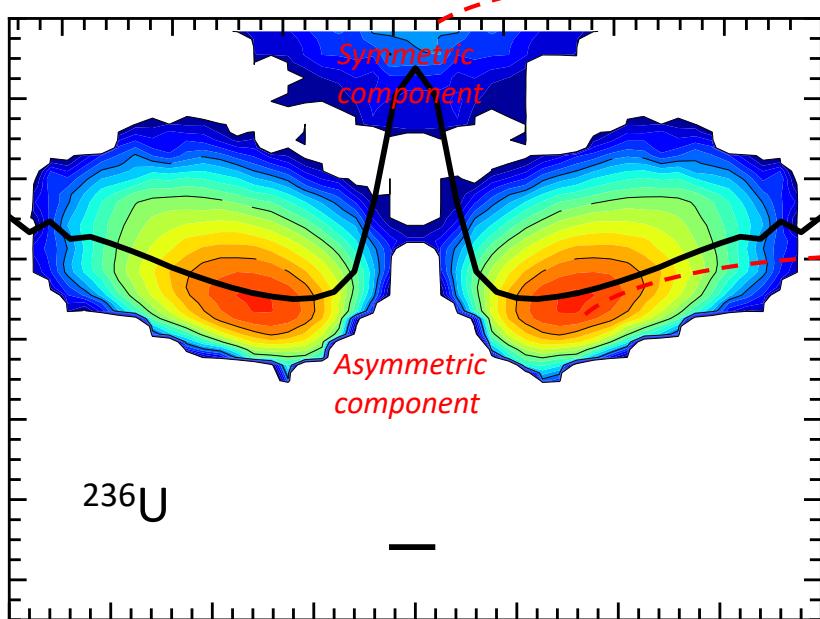
(e)

(f)

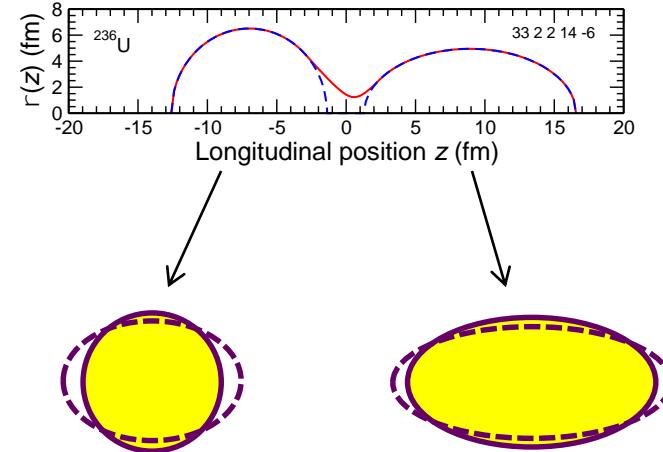
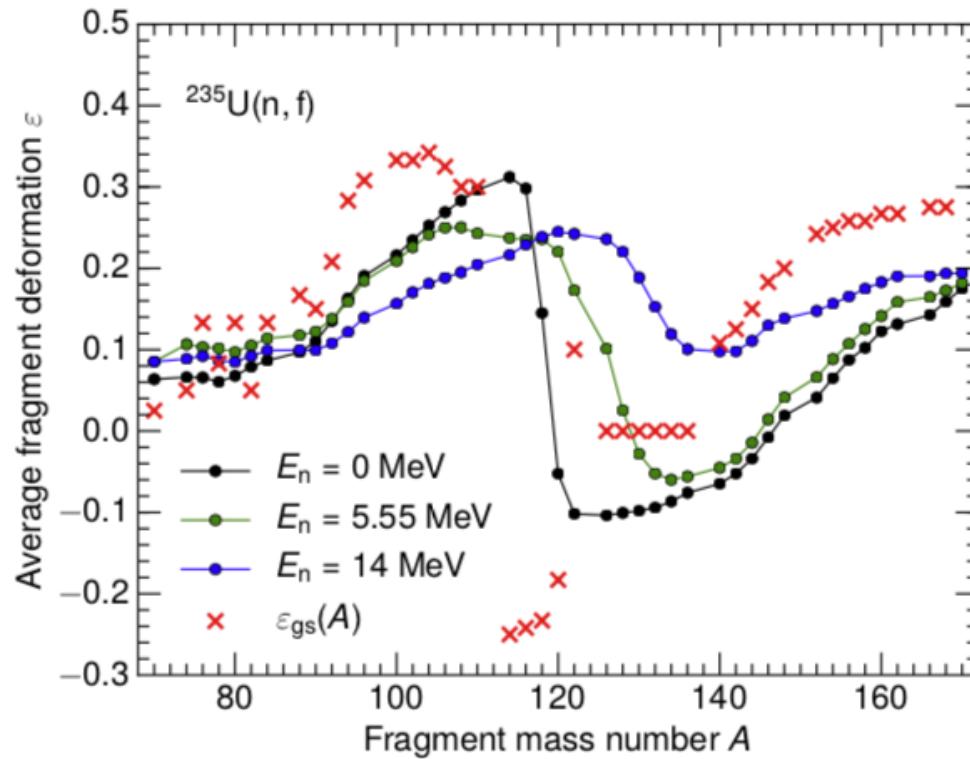
-0.4 -0.2 0.0 0.2 0.4

Mass asymmetry  $\alpha_g$

# Scission shapes ( $c_0 = 1.5$ fm)

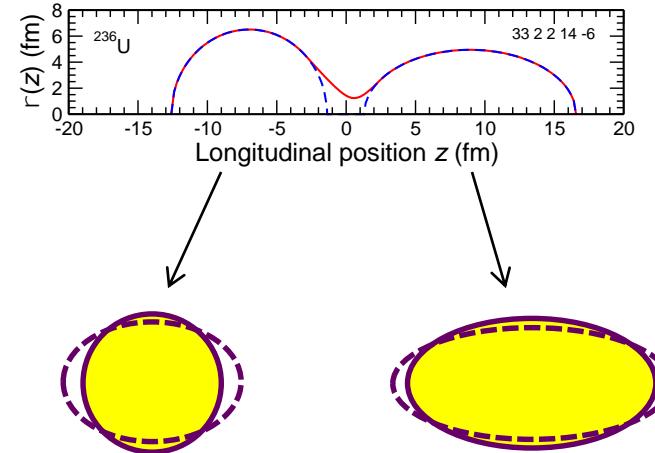
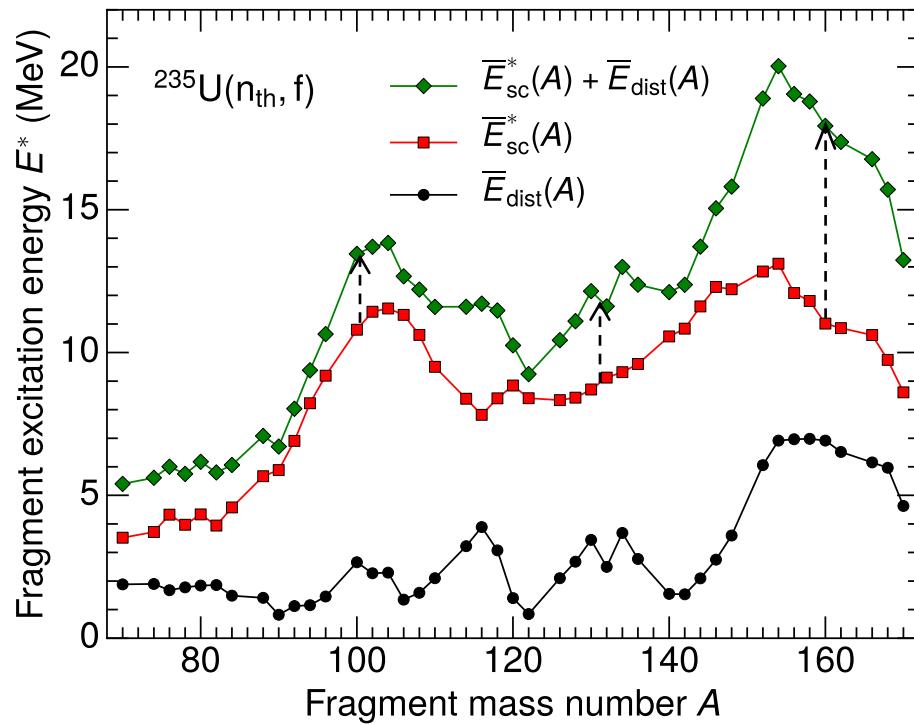


## Fragment shapes at scission: *distortion energy*



After separation the shapes relax to their equilibrium forms:  $\varepsilon_{\text{sc}} \rightarrow \varepsilon_{\text{gs}}$   
 and the associated fragment distortion energy  $E_{\text{dist}} = M(\varepsilon_{\text{sc}}) - M(\varepsilon_{\text{gs}})$   
 is converted to additional statistical excitation energy

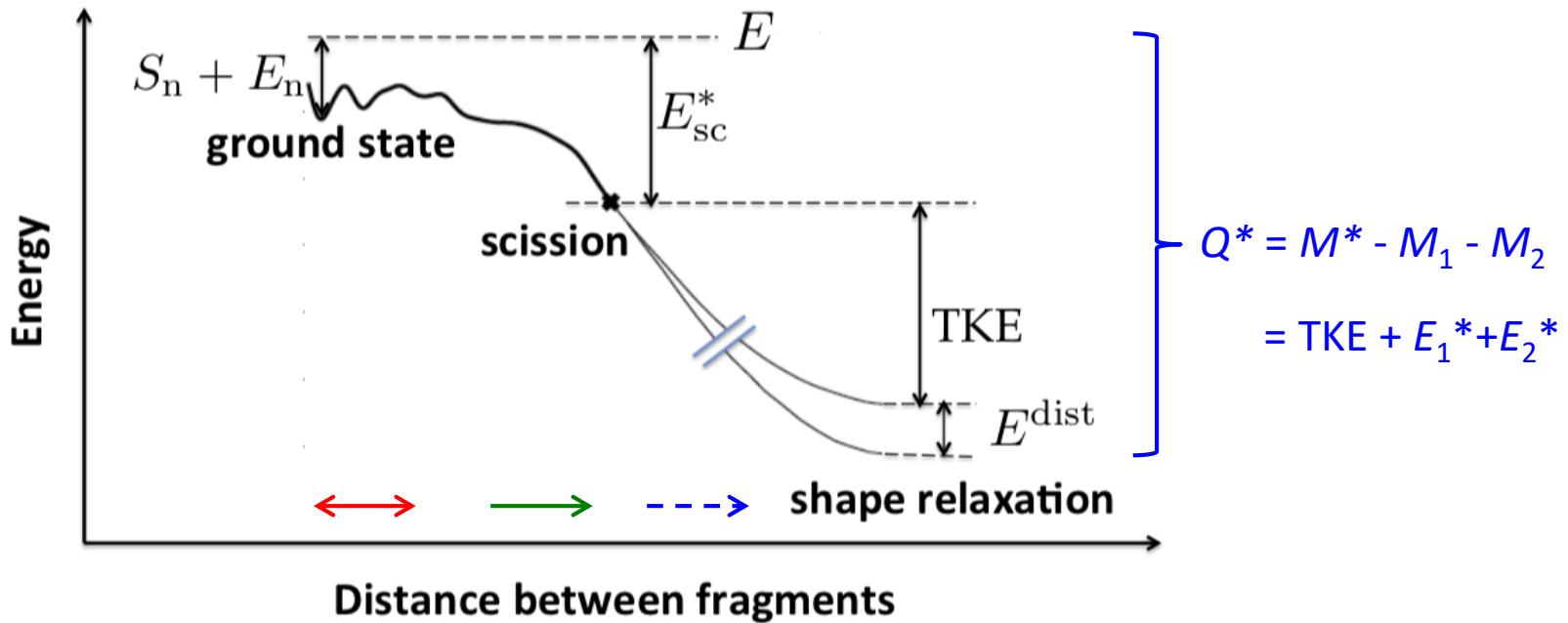
# Fragment shapes at scission: distortion energy



After separation the shapes relax to their equilibrium forms:  $\varepsilon_{sc} \rightarrow \varepsilon_{gs}$   
 and the associated fragment distortion energy  $E_{dist} = M(\varepsilon_{sc}) - M(\varepsilon_{gs})$   
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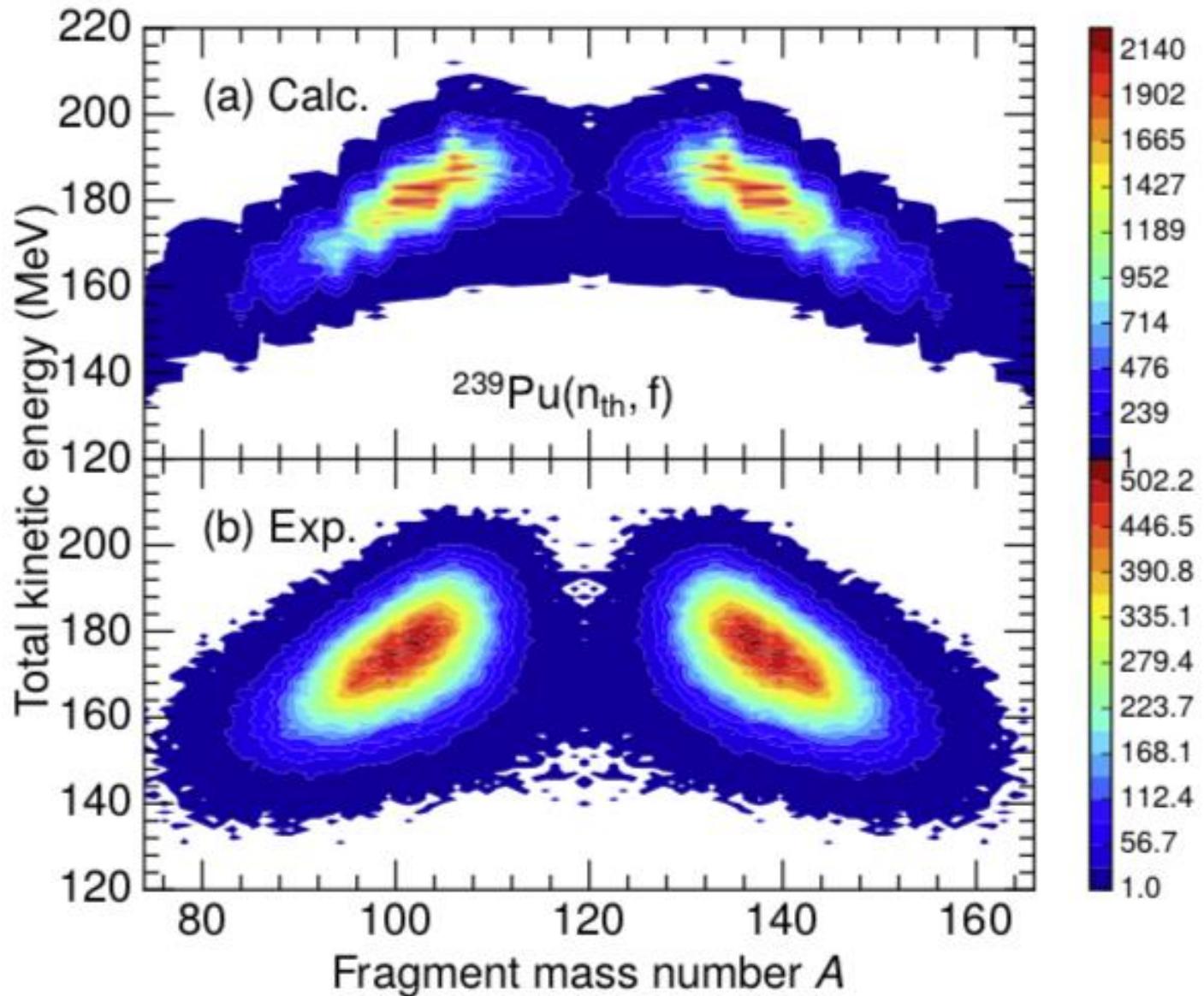
Tentative

✓ Extraction of TKE from the scission configuration

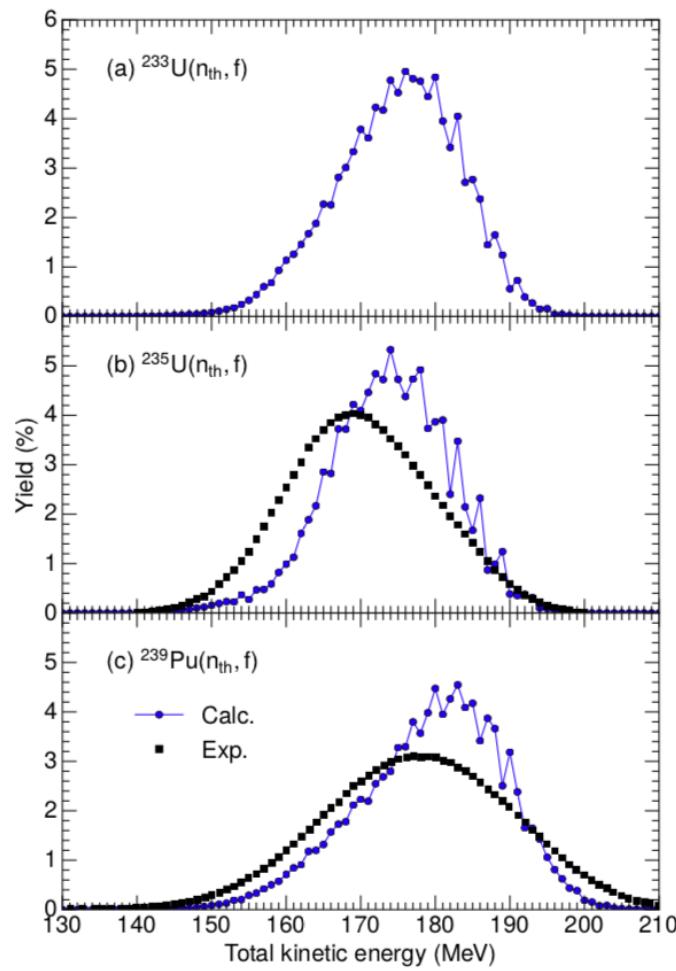
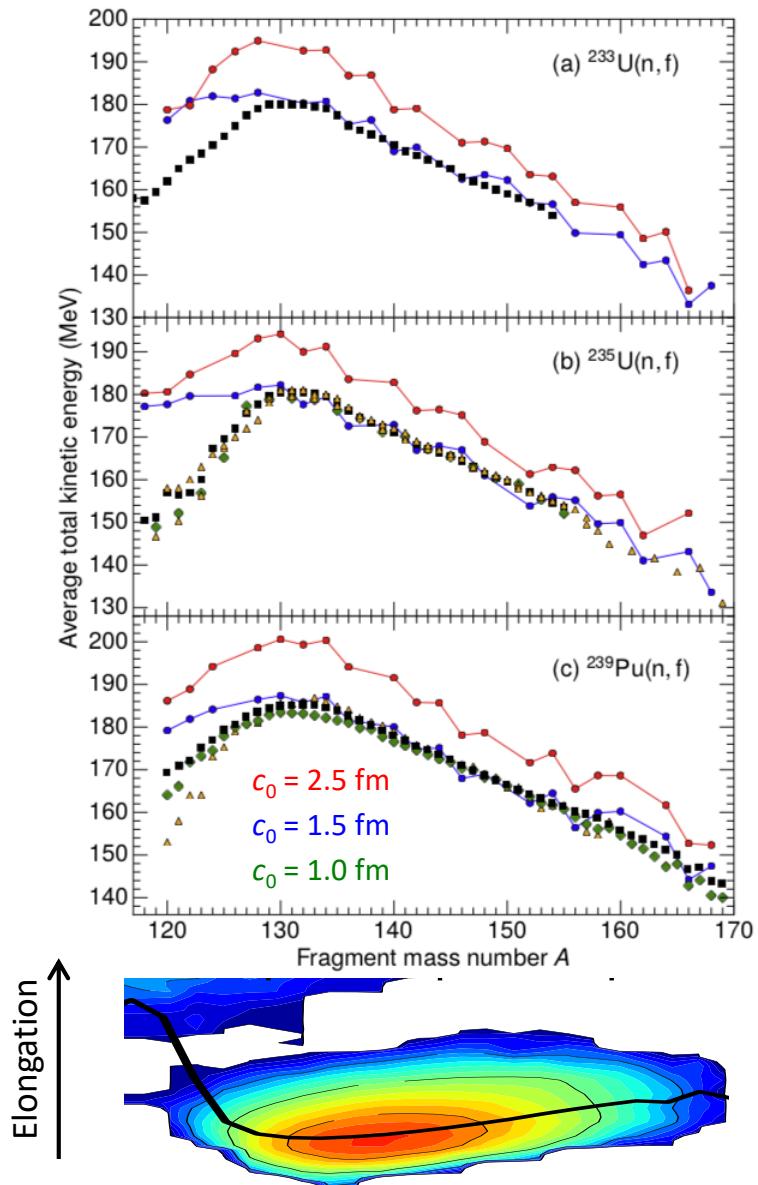


$$\text{TKE} = Q^* - E_{sciss}^* - E_{1,\text{dist}} - E_{2,\text{dist}}$$

# Joint ( $A$ ,TKE) distribution



# Total fragment kinetic energy



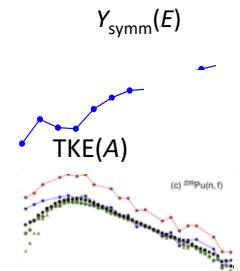
# *Fission studies with microscopic level densities*

## *I: Fragment mass distributions and kinetic energies*

The Brownian evolution of the nuclear shape can be simulated as a level-density guided Metropolis walk on the multi-dimensional deformation-energy surface

$$\frac{\nu(\chi \rightarrow \chi')}{\nu(\chi' \rightarrow \chi)} = \frac{\rho(\chi')}{\rho(\chi)}$$

The use of microscopic shape-dependent level densities provides a parameter-free model for calculating the intricate energy dependence of the fragment mass distribution and, perhaps, the total fragment kinetic energy distribution as well



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*Fission studies with microscopic level densities*

- ✓ *I: Fragment mass distributions and kinetic energies*
- *II: Microscopic treatment of energy partition*

*Martin Albertsson, Gillis Carlsson, David Ward, Sven Åberg, Mathematical Physics, Lund University, Sweden*

*Thomas Døssing, Niels Bohr Institute, Copenhagen University, Denmark*

*Peter Möller, Los Alamos National Laboratory, New Mexico, USA*

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