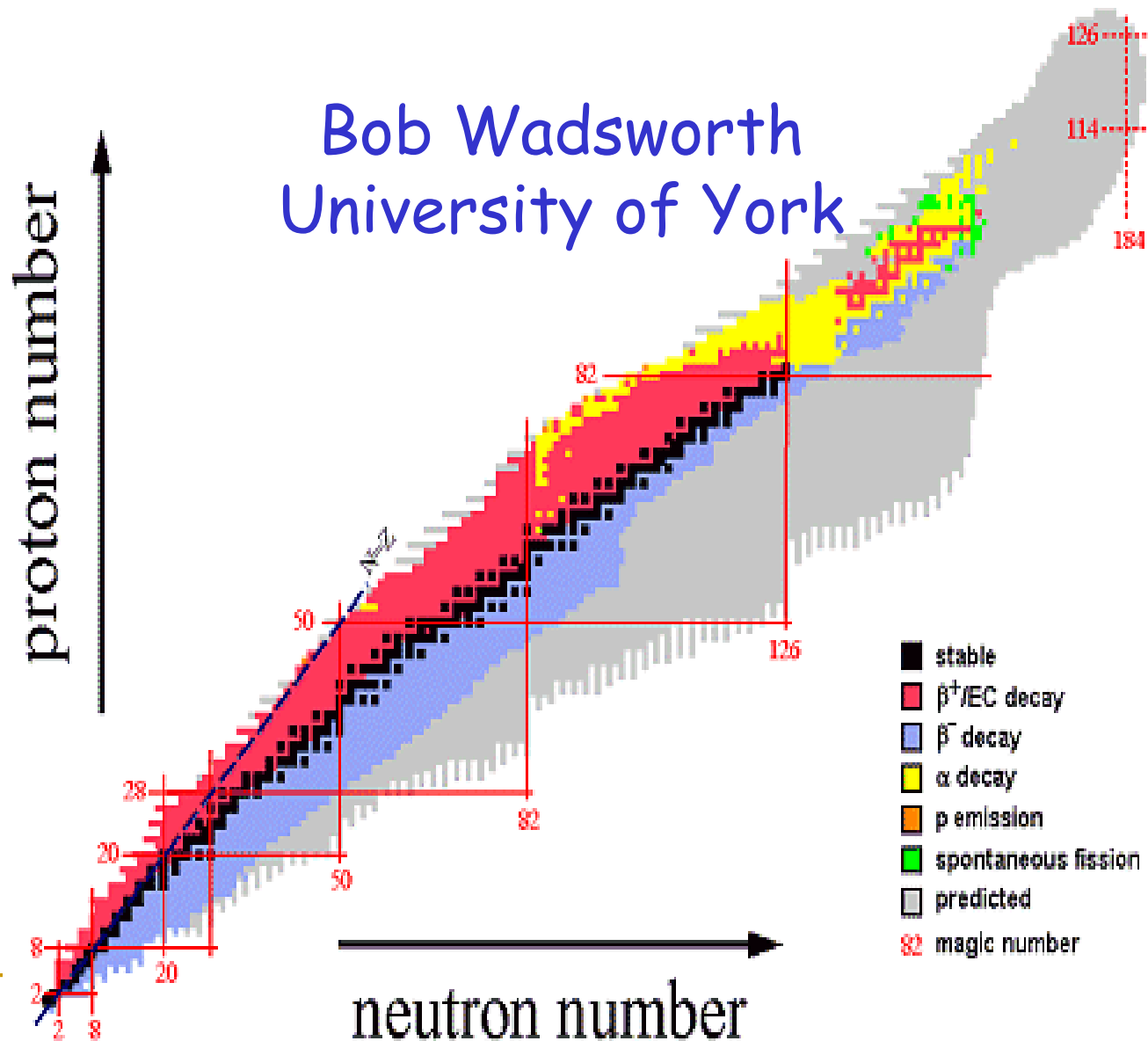


Spectroscopy near the N=Z line below ^{100}Sn

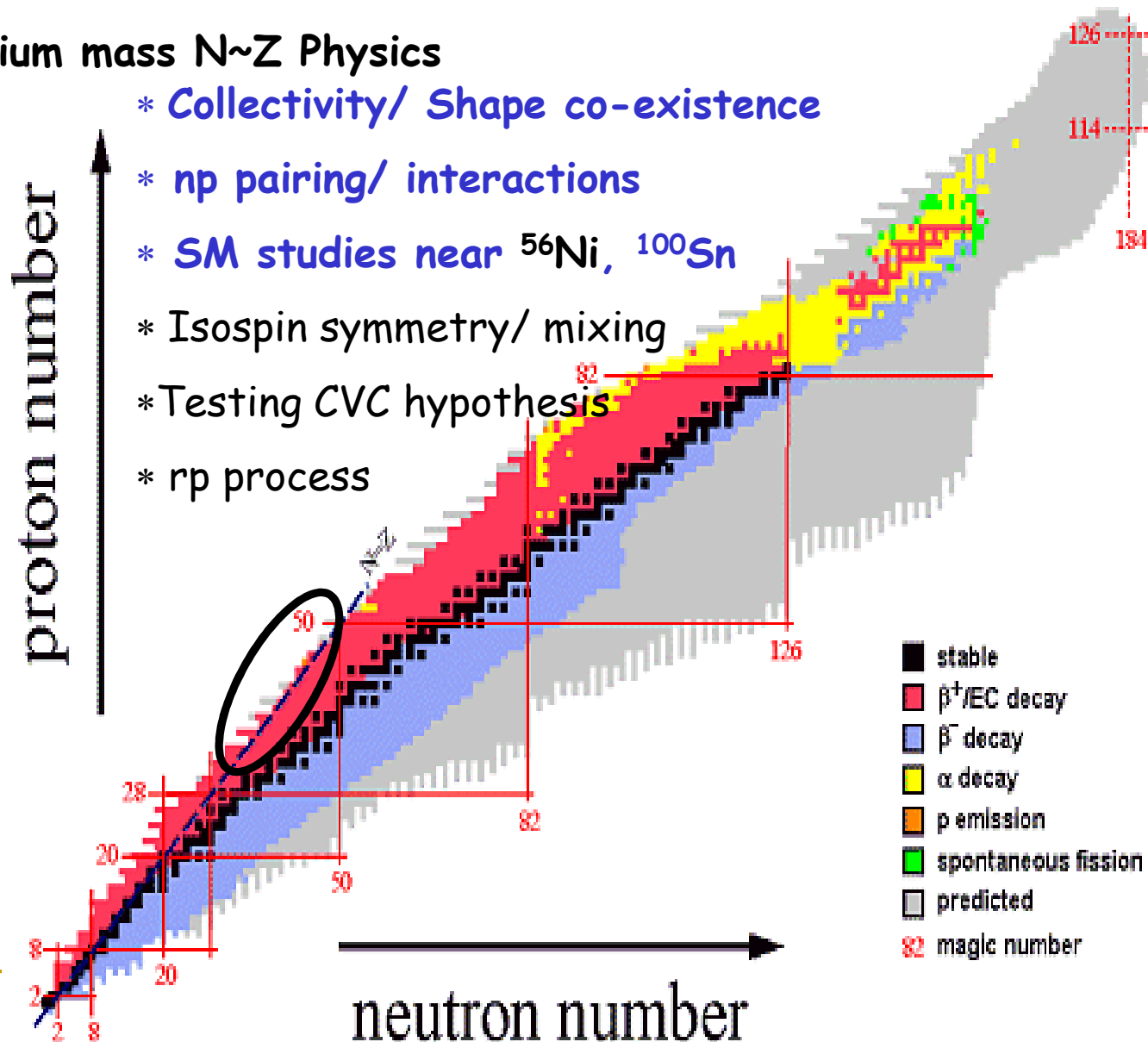
Bob Wadsworth
University of York



Spectroscopy near the N=Z line

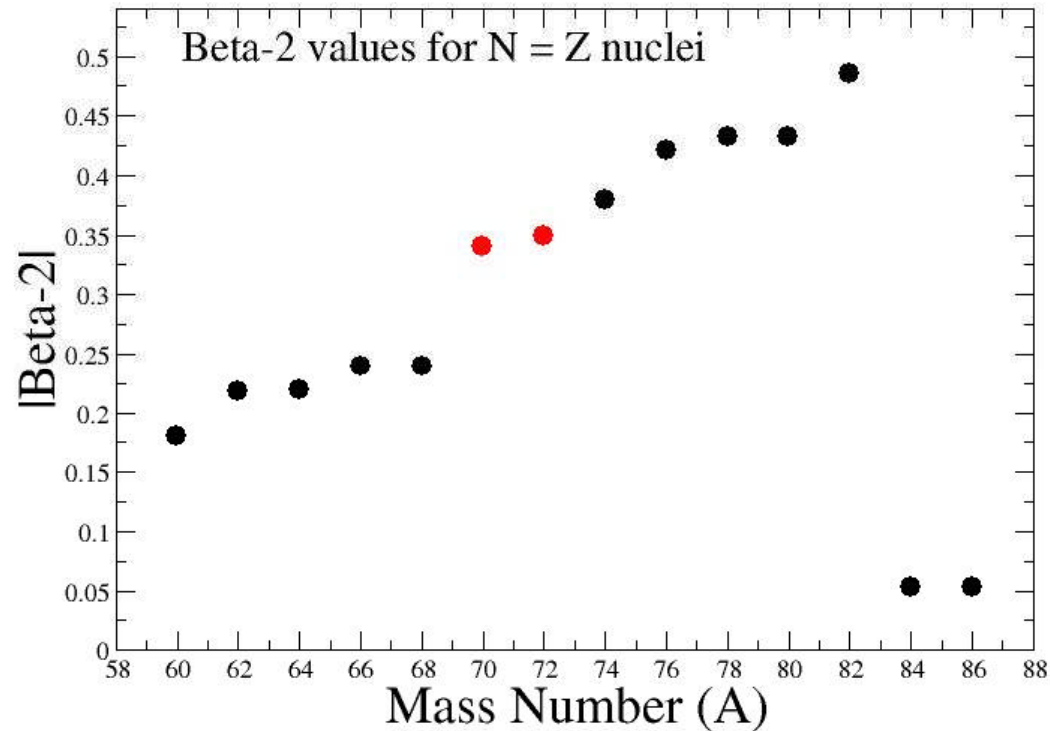
Medium mass N~Z Physics

- * Collectivity/ Shape co-existence
- * np pairing/ interactions
- * SM studies near ^{56}Ni , ^{100}Sn
- * Isospin symmetry/ mixing
- * Testing CVC hypothesis
- * rp process



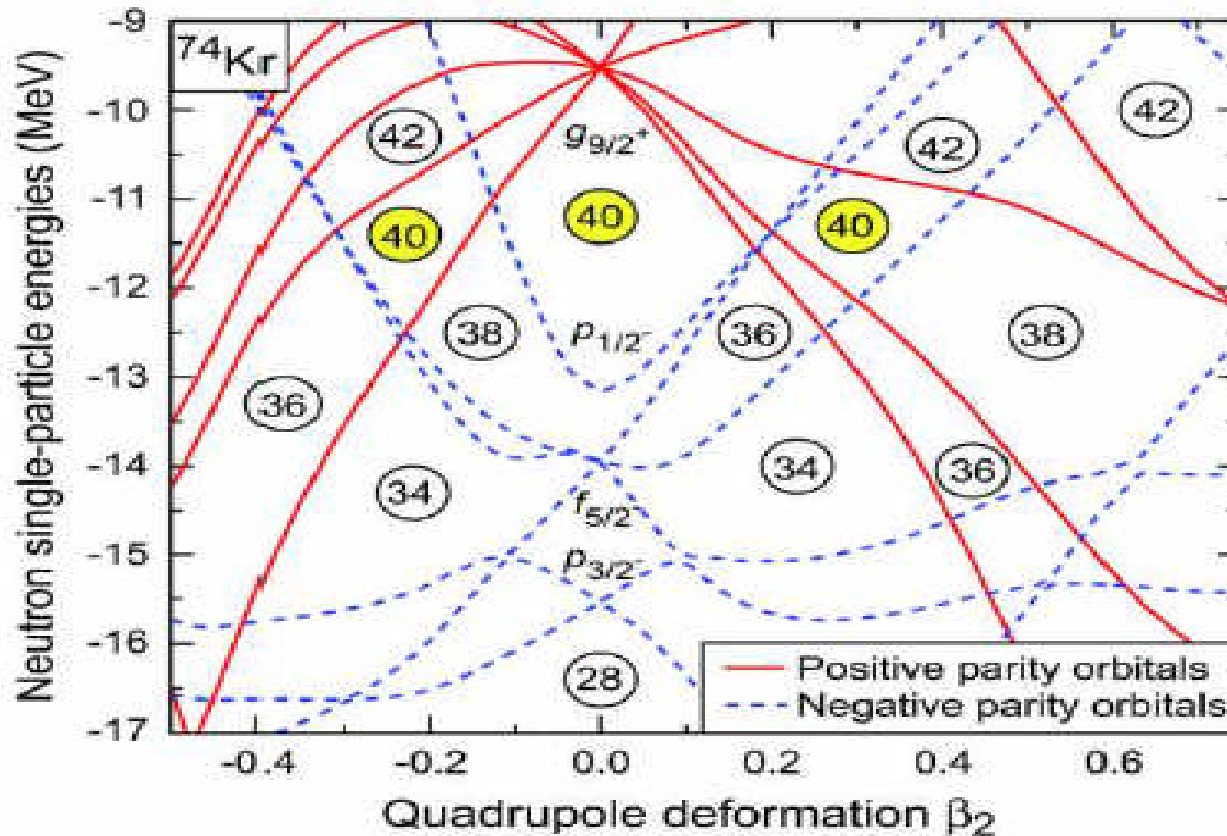
Collectivity along the $N=Z$ line

Finite-range droplet macroscopic model + folded-Yukawa single-particle microscopic model



Moller, Nix, Myers, Swiatecki
Atomic and Nuclear data tables, 59, 185 (1995)

Collectivity along the N=Z line



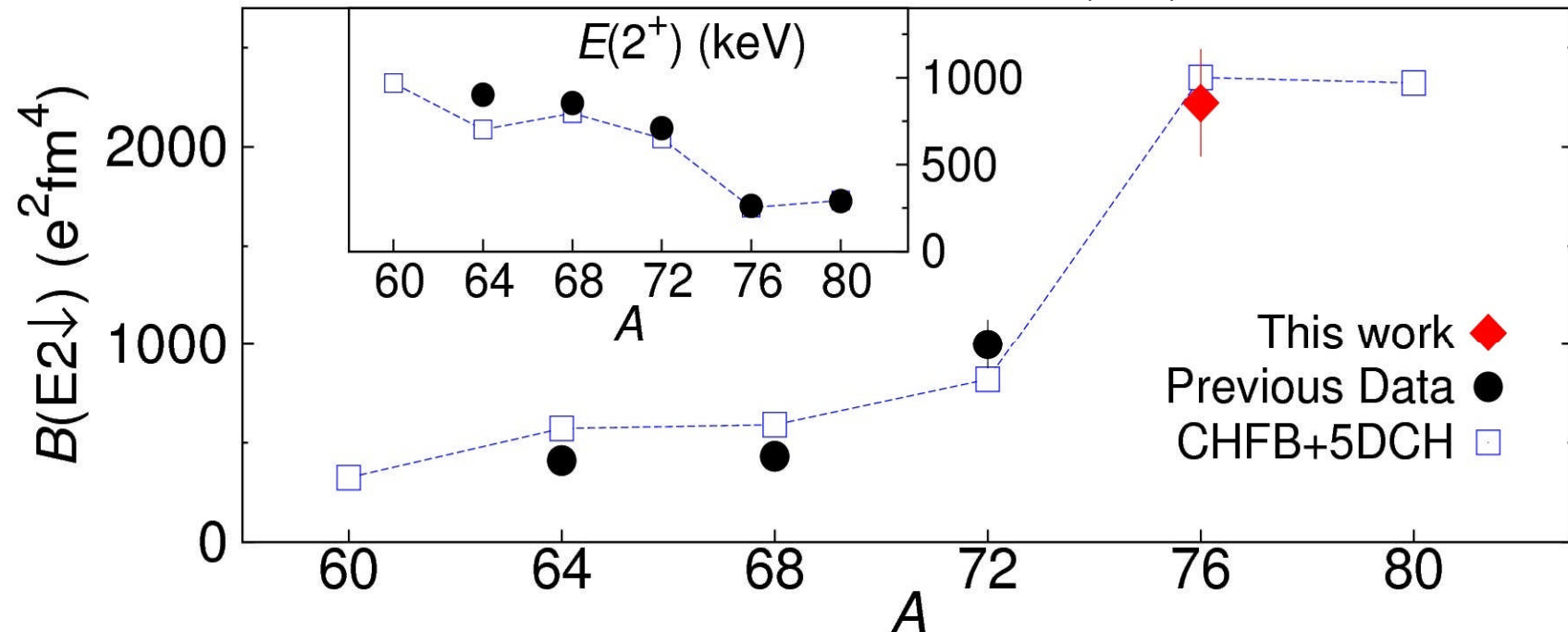
- Nuclear deformation and collectivity in the mass 70-80 is region is largely driven by proton, neutron occupancy of the $g_{9/2}$ orbit.
- Large shell-gaps at prolate, spherical and oblate shapes results in the potential for shape co-existence for many nuclei in the mid-mass region

Collectivity along the N=Z line

A Lemasson et al., PRC 85, 041303(R) (2012)

Constrained HF Bogoliubov theory with mapping to the 5D collective Hamiltonian

J P Delaroche et al., PRC 81, 014303 (2010)



Transition strengths difficult to measure:

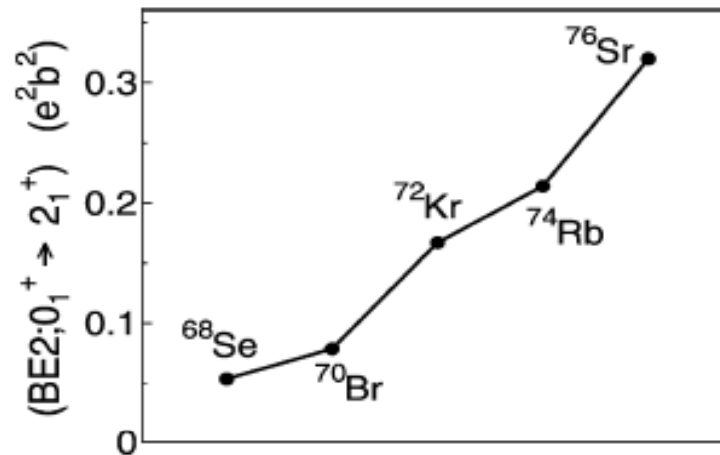
⁶⁸Se, ⁷²Kr Relativistic Coulex,

⁶⁴Ge Plunger lifetime expt following nucleon removal

⁷⁶Sr Doppler shift lineshape measurements + charge exchange reaction (#)

Collectivity along the N=Z line

Shell model using a truncated $f_{5/2}, p_{1/2}, g_{9/2}, d_{5/2}$ model space



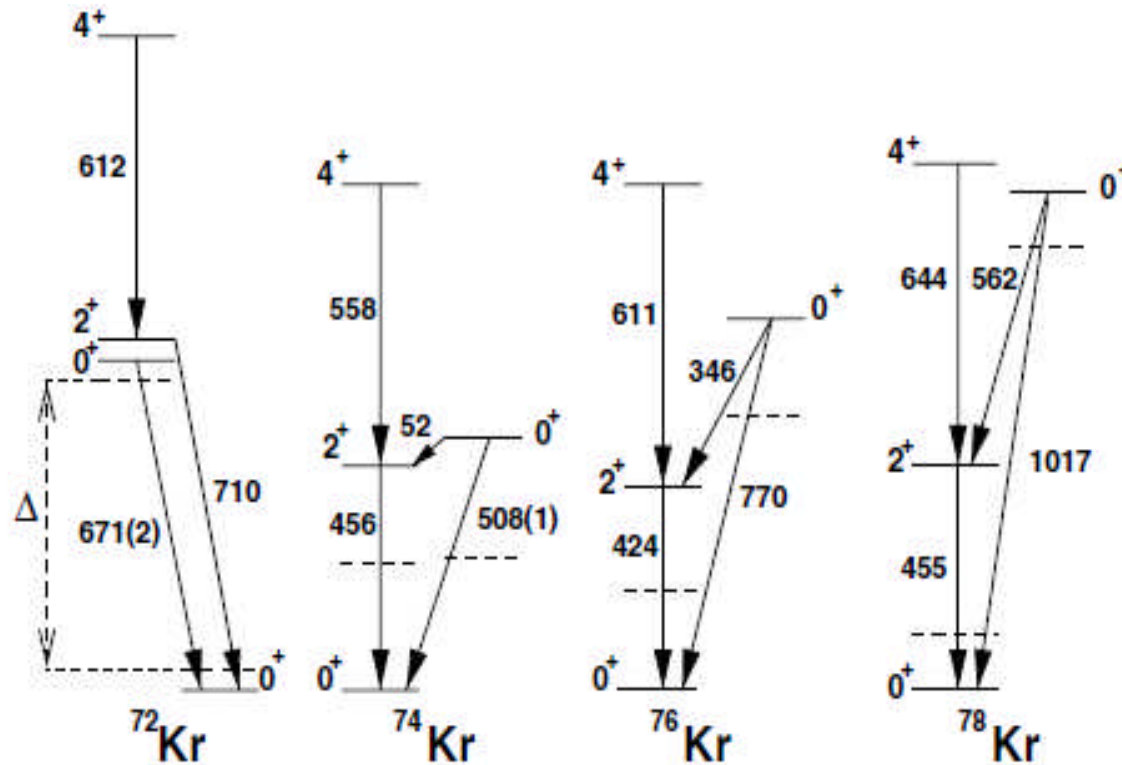
M Hasagawa, K Kaneko, T Mizusaki, Y Sun, PLB 656, 51 (2007)

Sharp increase in $B(E2)$ values beyond ^{70}Br is attributed to a sudden jump of protons and neutrons into the upper gd shell.

**Separator + recoil- β -tagging (or PPAC/ Ion Chamber) + Differential Plunger
Need a good efficient γ ray array**

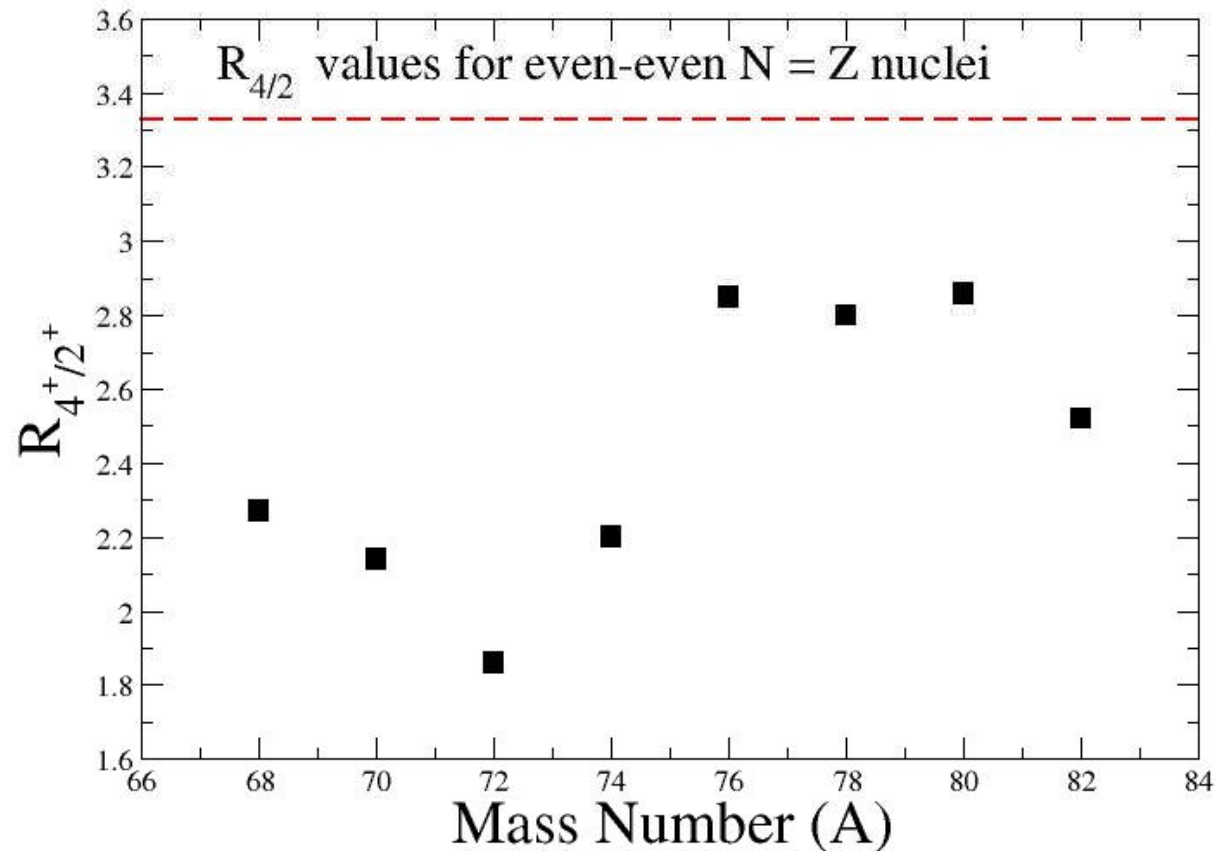
Shape co-existence along the N=Z line

Even-A Kr isotopes are known to exhibit shape co-existence, with excited 0^+ states already identified: **E Bouchez et al., PRL 90, 083502 (2003)**



Shape co-existence along the N=Z line

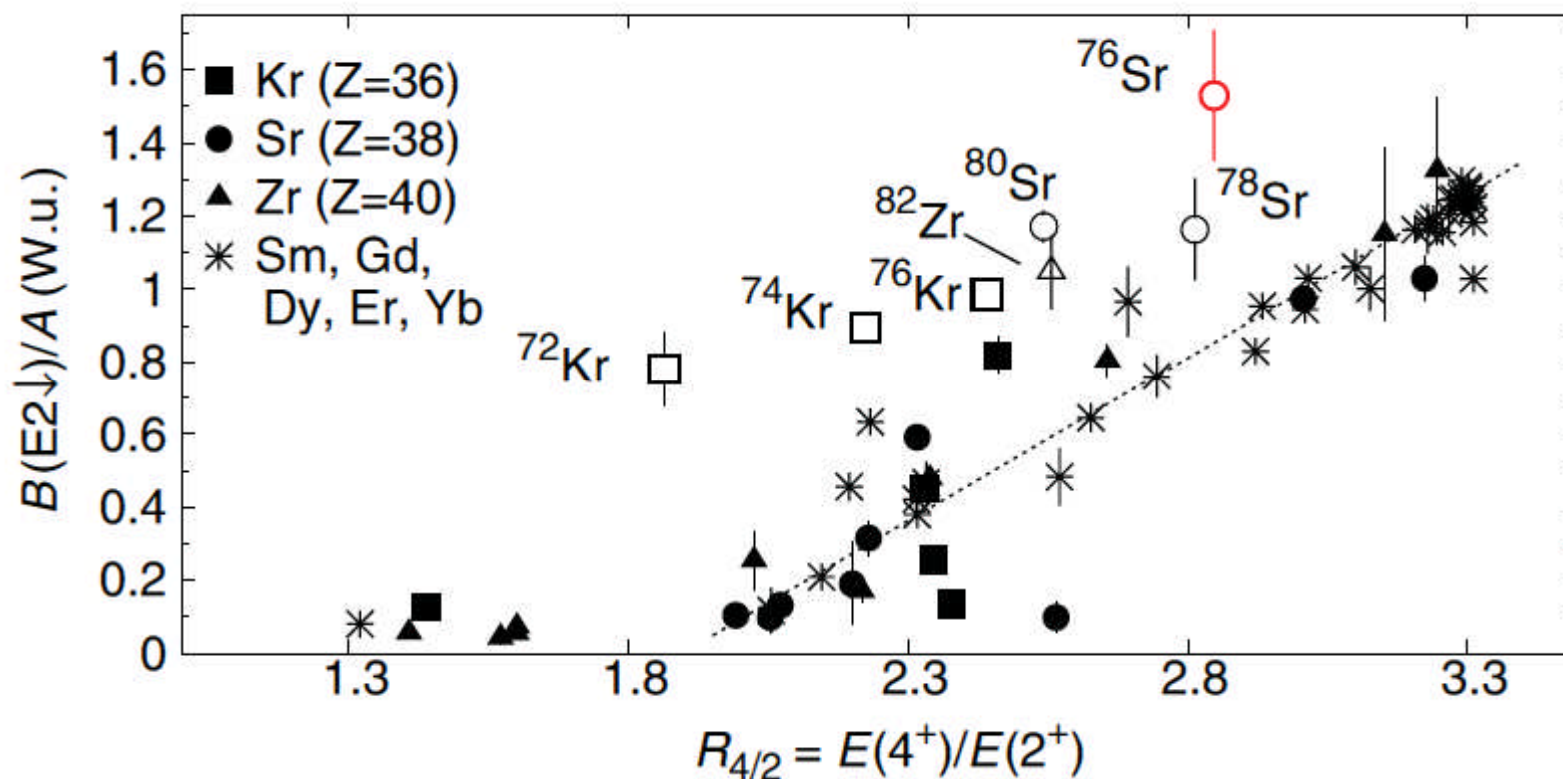
$$R_{4/2} = E_{4+}/E_{2+}$$



Why are the $R_{4/2}$ values well below the rotational limit for A~80 nuclei?

Shape co-existence along the N=Z line

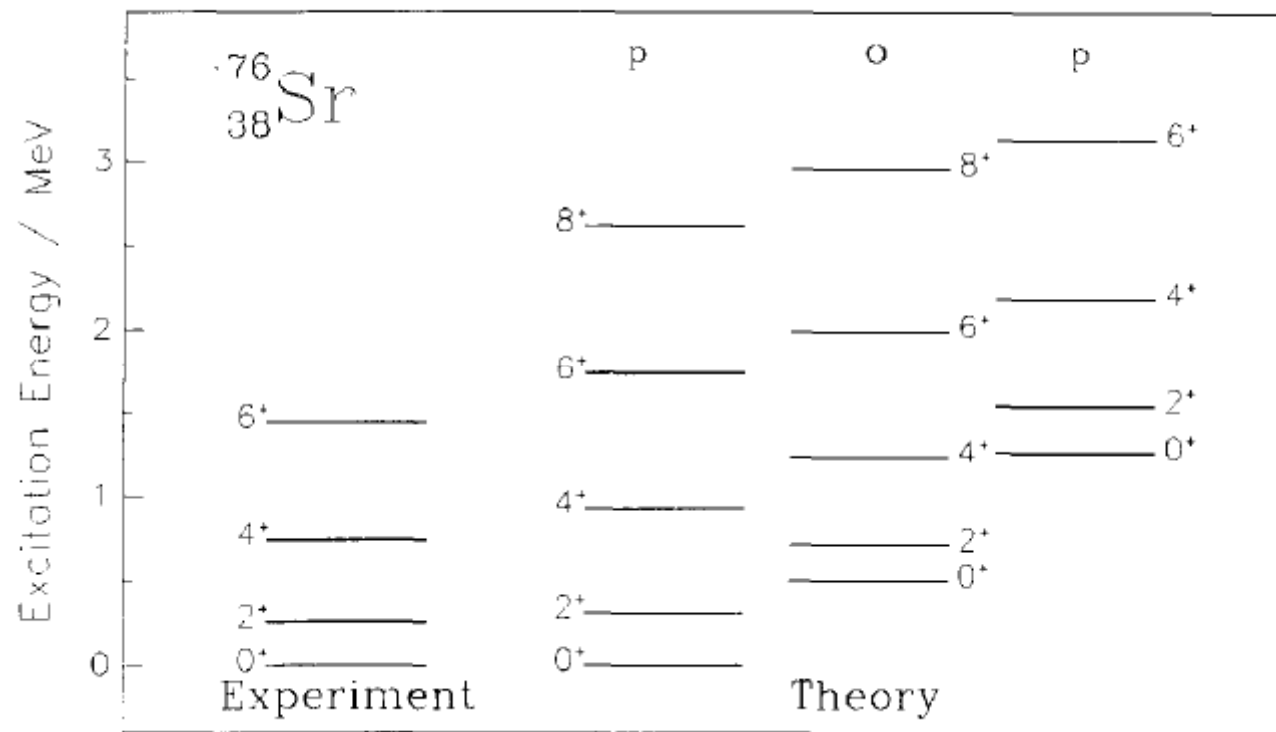
A Lemasson et al., PRC 85, 041303(R) (2012)



- * Clear correlation between $B(E2)/A$ and $R_{4/2}$ values for vibrational to rotational nuclei:
- * Deviations (OPEN SYMBOLS) occur in nuclei where shape co-existence is known or expected

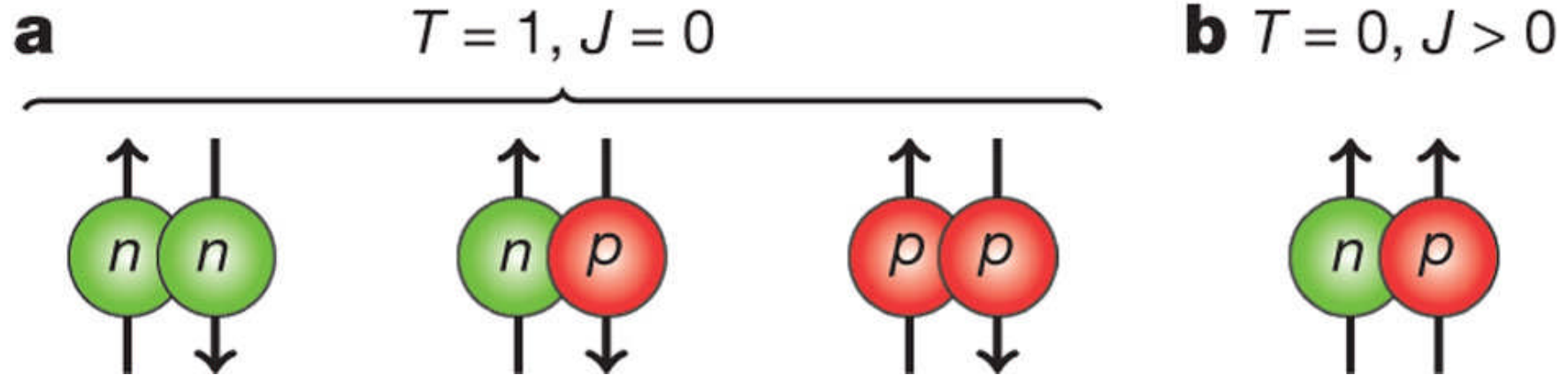
Shape co-existence along the N=Z line

0^+_2 in ^{76}Sr at ~ 0.5 MeV: A. Petrovici *et al.*, *Nucl. Phys. A605*, 290 (1996).



Nuclei such as ^{76}Sr , ^{78}Y , ^{80}Zr , ^{82}Nb , ^{84}Mo etc only have data on yrast states
 \Rightarrow clearly need information on non-yrast states to test the hypothesis that shape co-existence may be responsible for low $R_{4/2}$ values.

Neutron-proton pairing in $N=Z$ nuclei



- Studies of Binding energies in e-e and o-o nuclei indicate that $T=1$ np pairing is dominant, with no evidence for a $T=0$ (deuteron-like) pair condensate.

[P. Vogel, Nucl. Phys. A662 \(2000\) 148,](#)

[A.O. Macchiavelli et al PRC 61 \(2000\) 014303R](#)

- Comparison of data with mean-field calculations for $A=68-80$ nuclei suggests the presence of a strong isovector ($T=1$) np pair field at low spin, but no evidence for $T=0$ pairing.

[A Afanasjev, S Frauendorf, Phys. Rev. C 71, 064318 \(2005\)](#)

- Odd-odd nuclei such as $^{66}\text{As}, ^{70}\text{Br}, ^{74}\text{Rb}$ have $T = 1, 0^+$ grd states, but no low-lying [$J=1, T=0$] state, implying $T=0$ pairing mode is weak in this mid-mass region.

- Does $T=0$ pairing/ interaction play a role at low or high spin in heavier $N=Z$ nuclei?

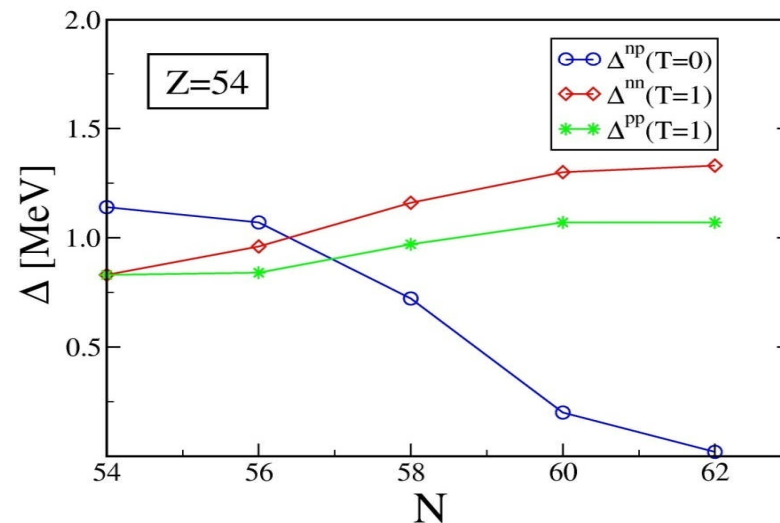
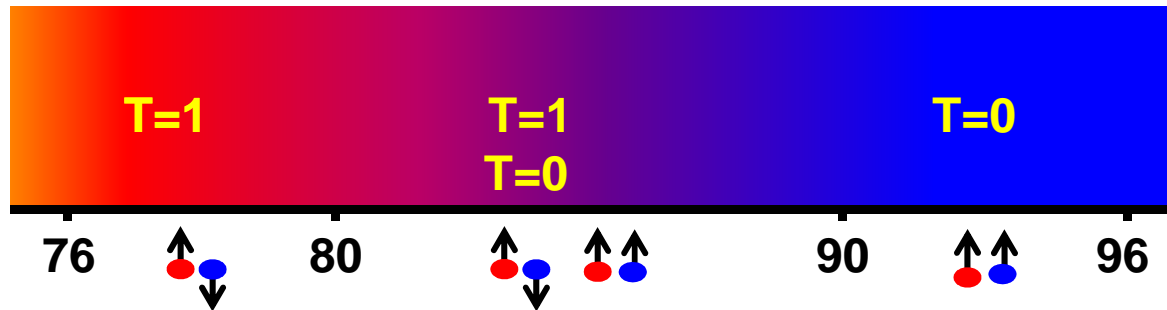
[A.L. Goodman PRC 58 R3051 \(1998\) and PRC 60, 014311 \(1999\)](#)

[W Satula, R Wyss PLB 393, 1 \(1997\) and PRL 86, 4488 \(2001\)](#)

[J Engel et al., PLB 389, 211 \(1996\) + Etc.](#)

Neutron-proton pairing in $N=Z$ nuclei

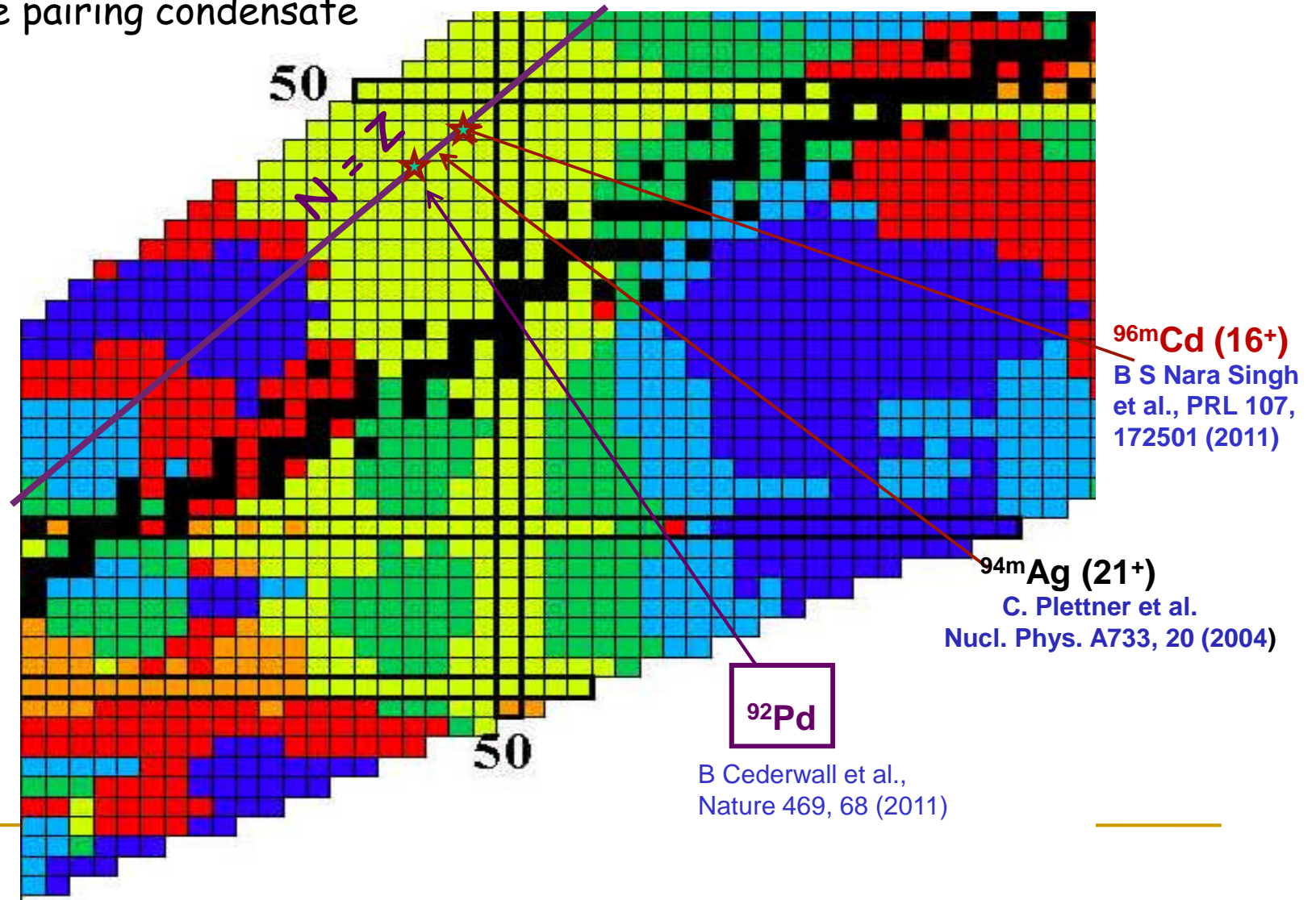
A. L. Goodman, PRC 60, 014311 (1999) -
studies of ground states of $e-e$ $N=Z$, $A=76-96$ nuclei



Calculation by W. Satula, R. Wyss, Phys. Rev. Lett. Vol. 86, 4488 (2001)

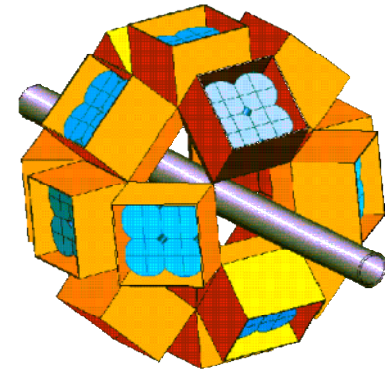
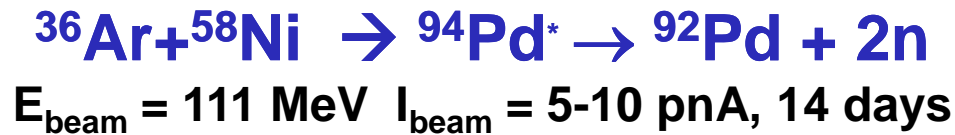
Evidence for isoscalar np interaction

As yet there is no data that definitively supports the presence of np T=0 BCS type pairing condensate



Evidence for isoscalar np interaction

Excited states in ^{92}Pd populated via fusion-evaporation at the Coulomb barrier (GANIL).



Detector systems:

EXOGRAM

$\epsilon \sim 0.11$

$\epsilon(\text{any charged particle}) \sim 0.66 \rightarrow$ veto efficiency for

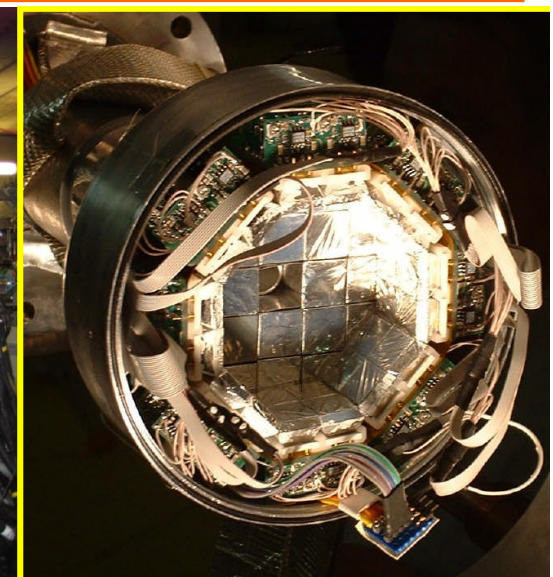
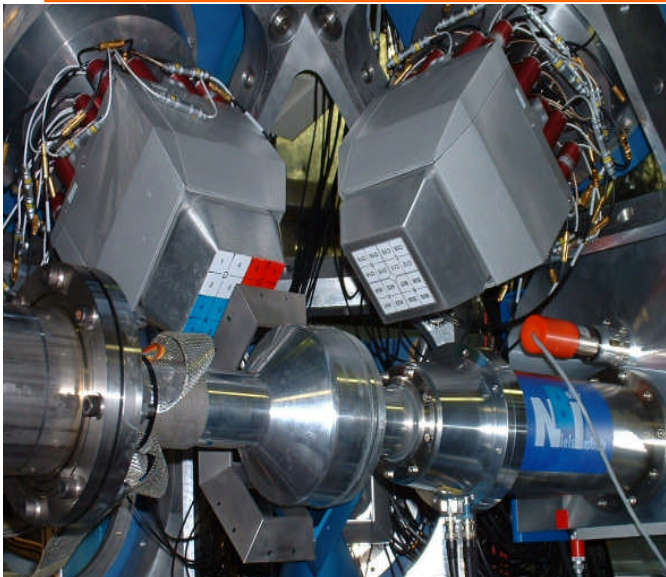
NEUTRON WALL

$\epsilon(n) \sim 0.25$ $\epsilon(\text{"clean" } 2n) \sim 0.03$

DIAMANT CsI(Tl) array

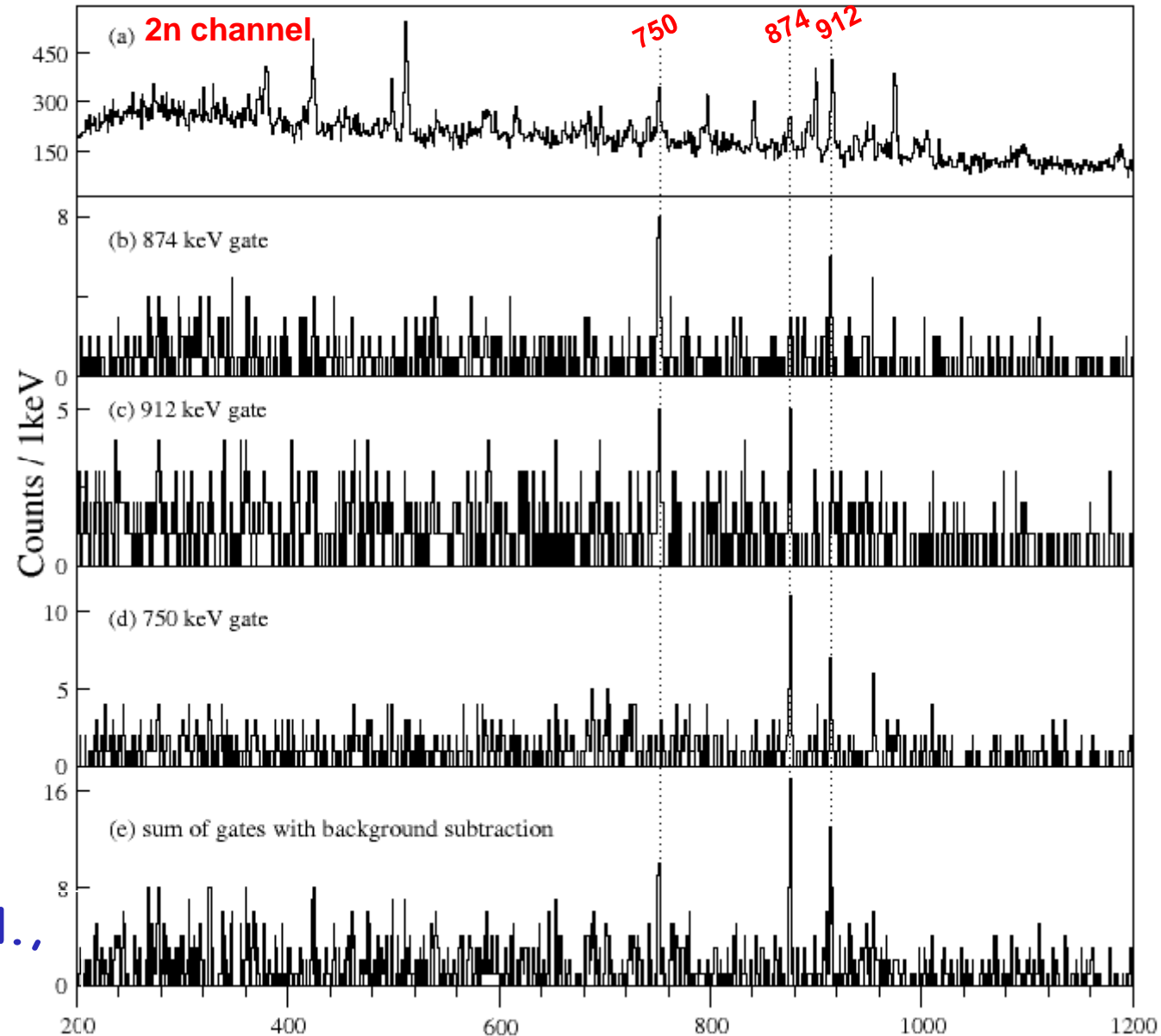
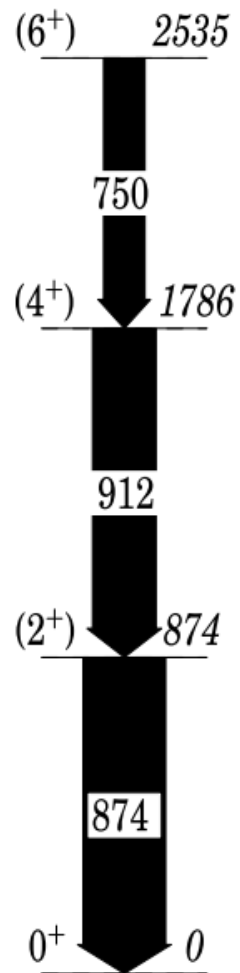
$\epsilon(p) \sim 0.50$ $\epsilon(\alpha) \sim 0.40$

particle mult. $>1 = 88\%$



^{92}Pd level scheme

H Al-Azri (York PhD student)

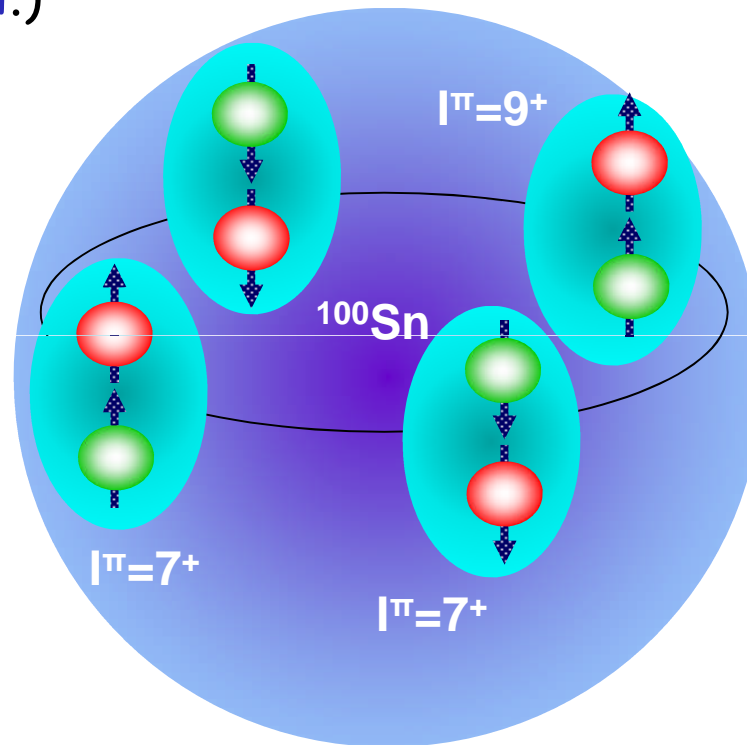


B Cederwall et al.,
Nature 469, 68
(2011)

Evidence for isoscalar np interaction

Shell Model Calculations in $p_{1/2}, g_{9/2}$ space predict strong np interactions

→ Spin-aligned T=0 np coupling scheme for N=Z nuclei below ^{100}Sn
(J. Blomqvist et al.)



$^{92}\text{Pd}_{\text{gs}}$

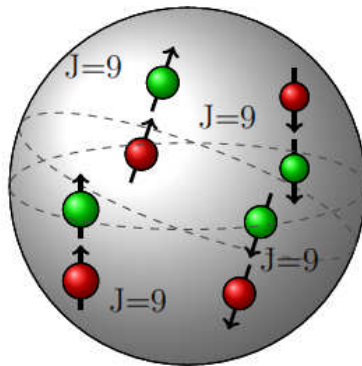
• 4-deuteron hole-like pairs coupled to $J=9$, each with a different angular momentum projection $M = +9, -9, +7, -7$ to satisfy the Pauli Principle.

Aligned np coupling: $\Psi_{\text{G.S.}} = [(\{v g_{9/2}^{-1} \times \pi g_{9/2}^{-1}\}_{9+})^2]_{0+} \times [(\{v g_{9/2}^{-1} \times \pi g_{9/2}^{-1}\}_{7+})^2]_{0+}$

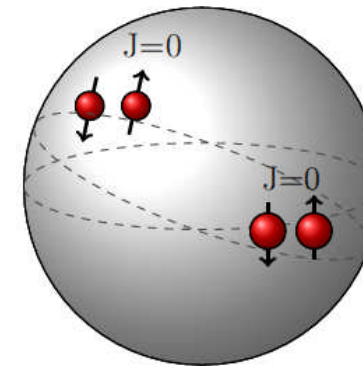
• Similar results confirmed for ^{96}Cd - S Zerguine and P Van Isacker, PRC 83, 064311 (2011)

Evidence for isoscalar np interaction

Effect of isoscalar np interaction at N=Z



Shell model calculations by J. Blomqvist et al.

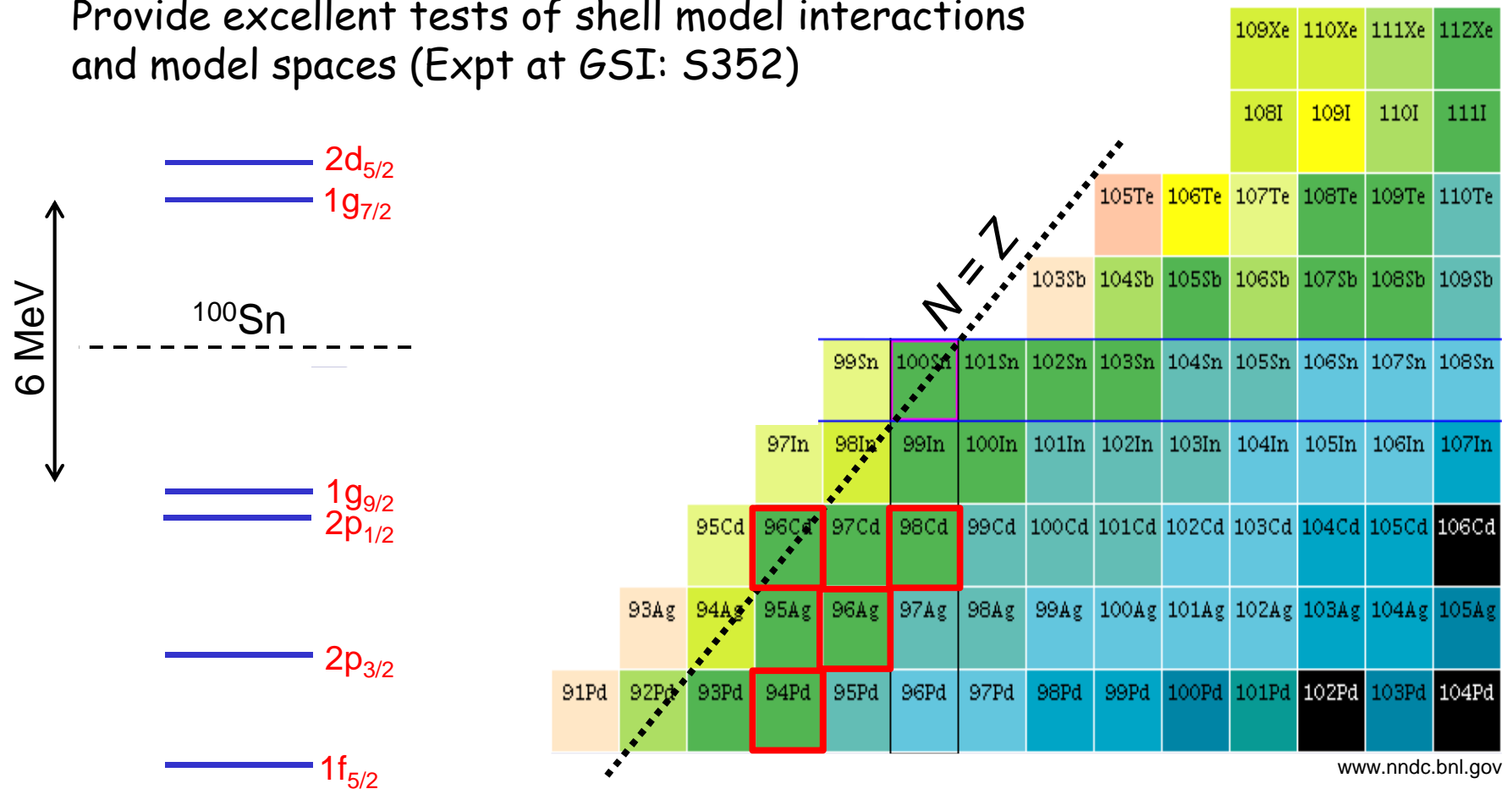


	10^+ 4072		10^+ 4065	10^+ 4052		10^+ 4052	10^+ 4065		10^+ 3862	10^+ 3796		10^+ 4131	10^+ 3784
	8^+ 3127	10^+ 3257						10^+ 3257					
(6^+) 2536	6^+ 2466	8^+ 2600	8^+ 2749	8^+ 2633	8^+ 2635	8^+ 2588	8^+ 2792	8^+ 2750	8^+ 2704	8^+ 2636	8^+ 2530		
(4^+) 1786	6^+ 2110	4^+ 2079	6^+ 2212	4^+ 2223	6^+ 2223	4^+ 2128	6^+ 2374	6^+ 2330	6^+ 2380	6^+ 2224	6^+ 2099		
	4^+ 1708	4^+ 1518					4^+ 1709	4^+ 1682	4^+ 1720	8.2			
	2^+ 878	2^+ 797	2^+ 1171	2^+ 1417	2^+ 1405	2^+ 1199				2^+ 1460	2^+ 1415		
	2^+ 874						2^+ 864	2^+ 861	2^+ 814	7.5			
0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0	0^+ 0
^{92}Pd exp	^{92}Pd SM	^{92}Pd T=0	^{92}Pd T=1	^{92}Pd no np	^{94}Pd no np	^{94}Pd T=1	^{94}Pd T=0	^{94}Pd SM	^{94}Pd exp	^{96}Pd SM	^{96}Pd exp		

Calculations done in several model spaces, i.e., $0g9/2$, $0g9/2-1p1/2$ and $0g9/2-1p1/2-0f5/2-1p3/2$ which all give similar results. Int. parameters determined to reproduce exp energies in $^{94,95}\text{Pd}$, $^{93,94}\text{Rh}$

Isomers in $N \cong Z$ nuclei below ^{100}Sn

Provide excellent tests of shell model interactions and model spaces (Expt at GSI: S352)



Isomers in $N \cong Z$ nuclei below ^{100}Sn

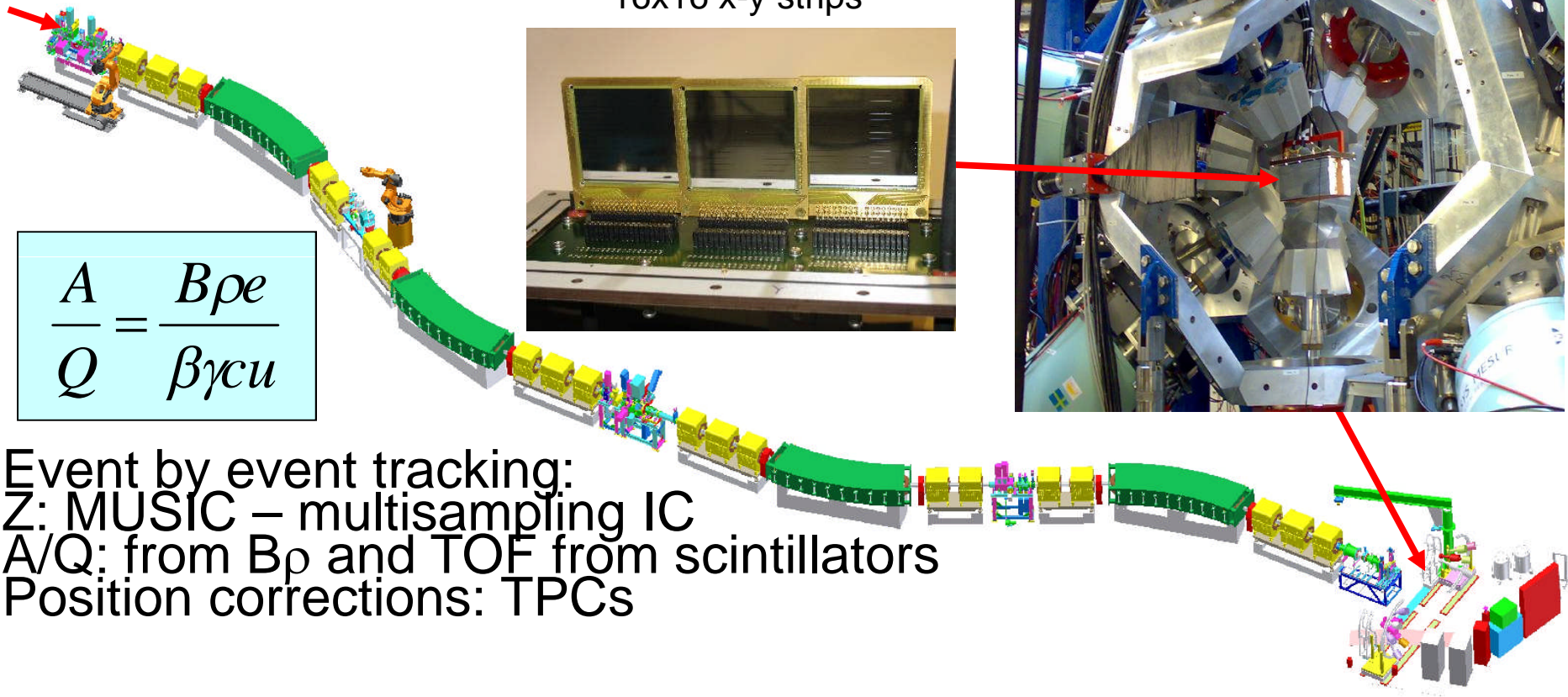
^{124}Xe : 850 MeV/u, $\sim 10^9$ pps

on ^9Be target

9 DSSSD, 1mm thick, $5 \times 5 \text{ cm}^2$
16x16 x-y strips

15 CLUSTERs x 7 crystals

$\epsilon_\gamma = 11\%$ at 1.3MeV

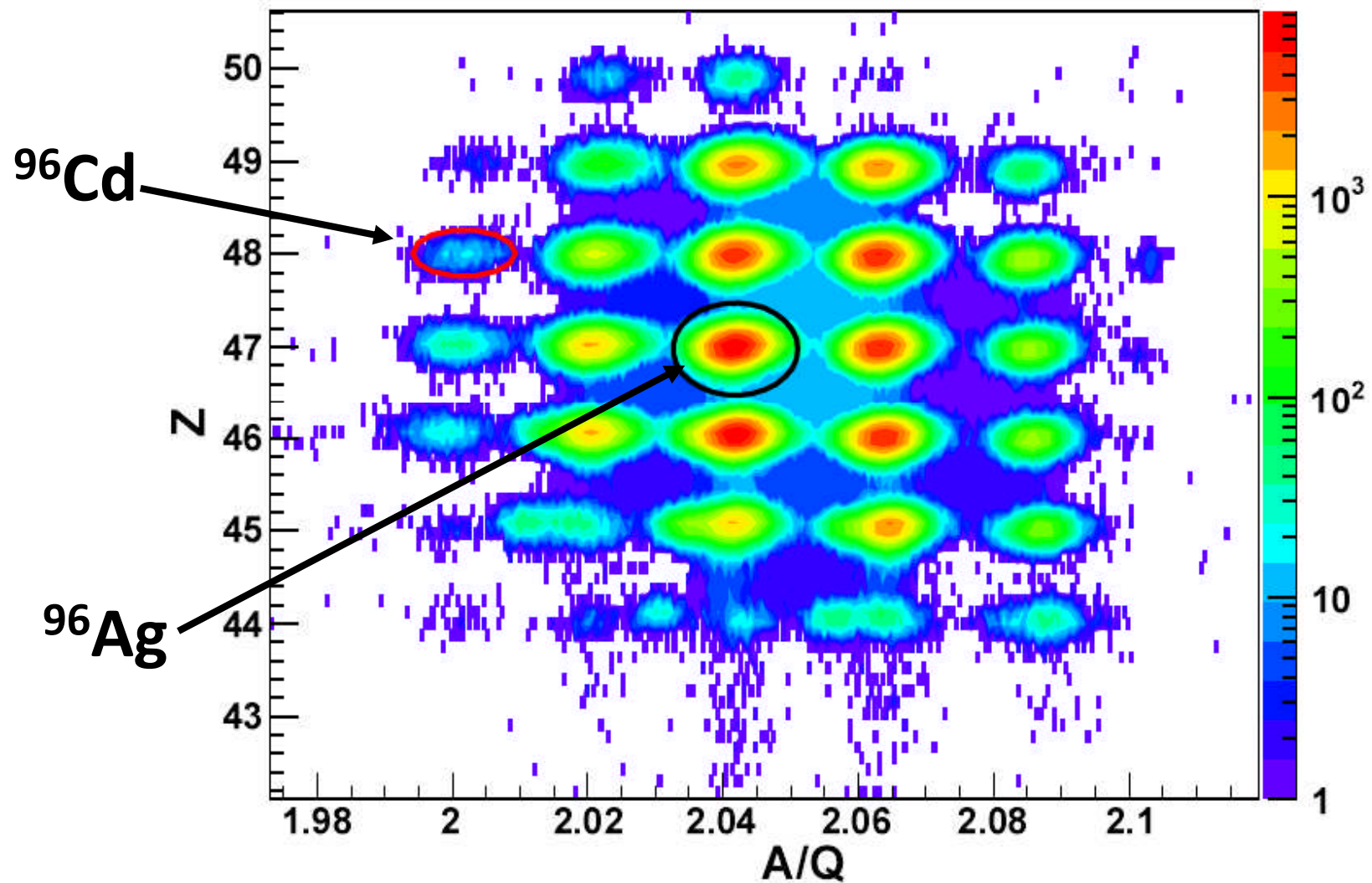


$$\frac{A}{Q} = \frac{B\rho e}{\beta\gamma cu}$$

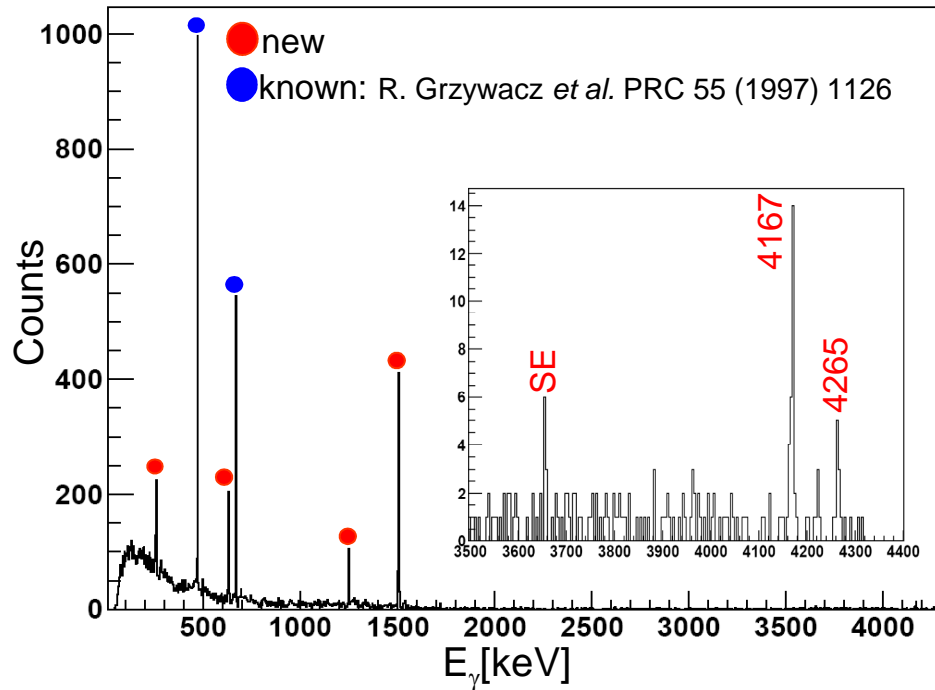
Event by event tracking:
 Z: MUSIC – multisampling IC
 A/Q: from $B\rho$ and TOF from scintillators
 Position corrections: TPCs

H. Geissel *et al.* NIM B70 (1992) 286
 S. Pietri *et al.* Nucl. Inst. and Meth. B261(2007) 1079
 R. Kumar *et al.* Nucl. Inst. and Meth. A598 (2009) 754

Isomers in $N \cong Z$ nuclei below ^{100}Sn



Isomers in ^{96}Ag



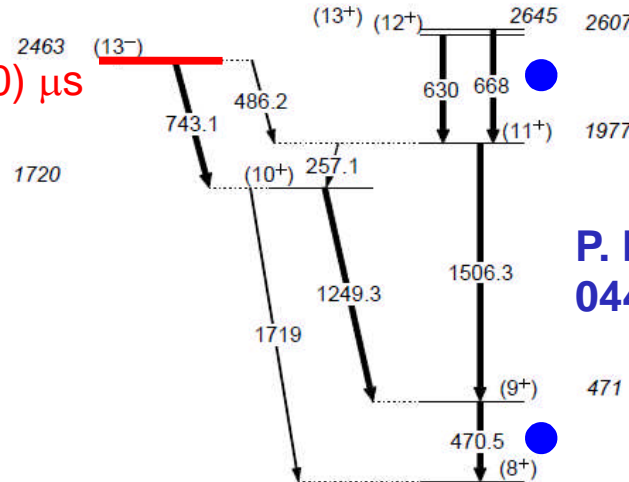
(19^+) (17^+) 96 $6813+x$ $0.16(3) \mu\text{s}$
 $6909+x$

4167

4265

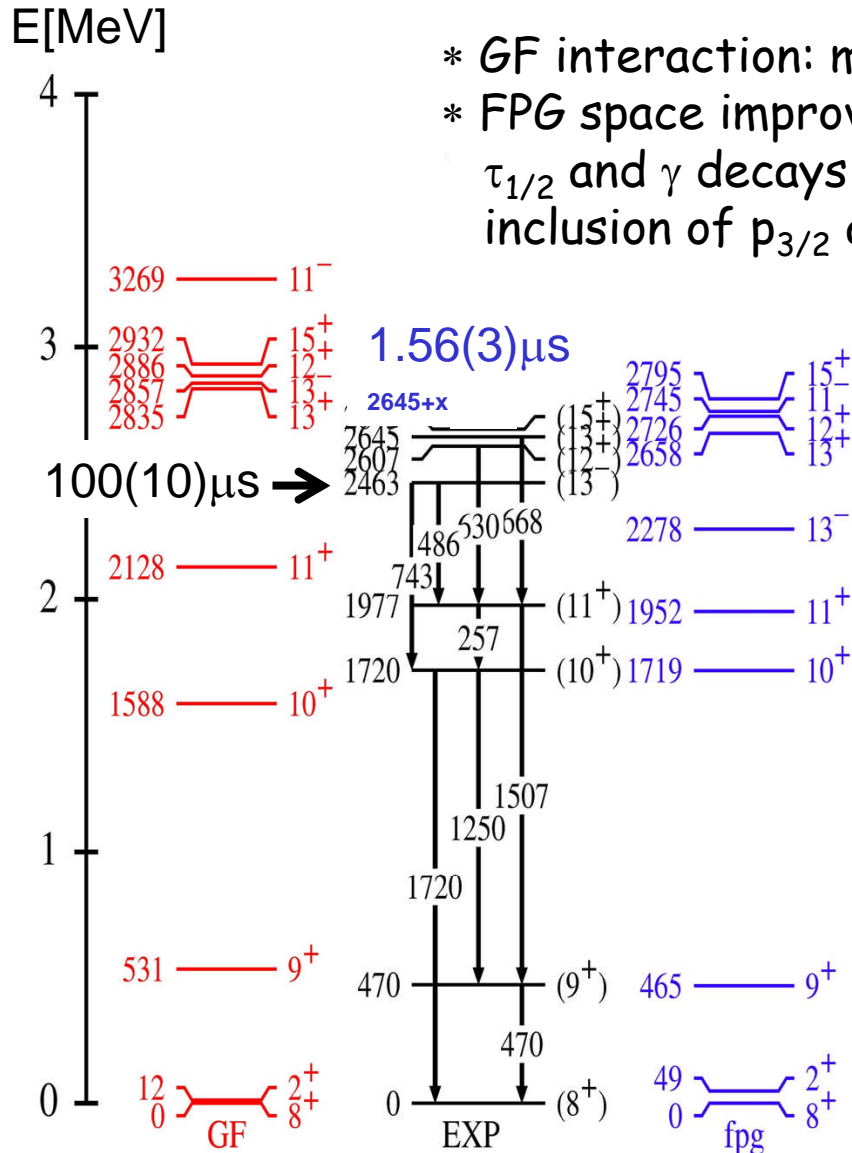
(13^-) 2463 $100(10) \mu\text{s}$

(13^+) (12^+) 2645 2607 (15^+) $2645+x$ $1.56(3) \mu\text{s}$



P. Boutachkov *et al.*, PRC 84, 044311 (2011)

Isomers in ^{96}Ag



- * GF interaction: model space: $\pi\nu(g_{9/2}, p_{1/2})$
- * FPG space improves position of 10^+ , 11^+ and 13^- levels, also $\tau_{1/2}$ and γ decays (E3 and M2) of new 100 μs isomer requires inclusion of $p_{3/2}$ and $f_{5/2}$ orbitals to calc B(E3), B(M2)'s

fpg: GF + SNA
 $\pi\nu$ 1p-1h excitation
 from $f_{5/2}$ and $p_{3/2}$

TBME from OXBASH package
 (SNA+GF) and SPE tuned to ^{100}Sn

J_i^π	J_f^π	σL	B($\sigma\lambda$) W.u.		
			Expt	GF	FPG
13^-	11^+	M2	$9.6(14) \times 10^{-5}$		3.6×10^{-5}
		E3	0.62(9)		0.53
	10^+	E3	0.145(17)		0.057
15^+	13^+	E2 ^a	2.45(6)	2.99	2.93

Electric trans calc with p/n eff. Charges of (1.5/0.5)e
 a = assuming $E_\gamma = 50$ keV

Core excited isomer in ^{96}Ag

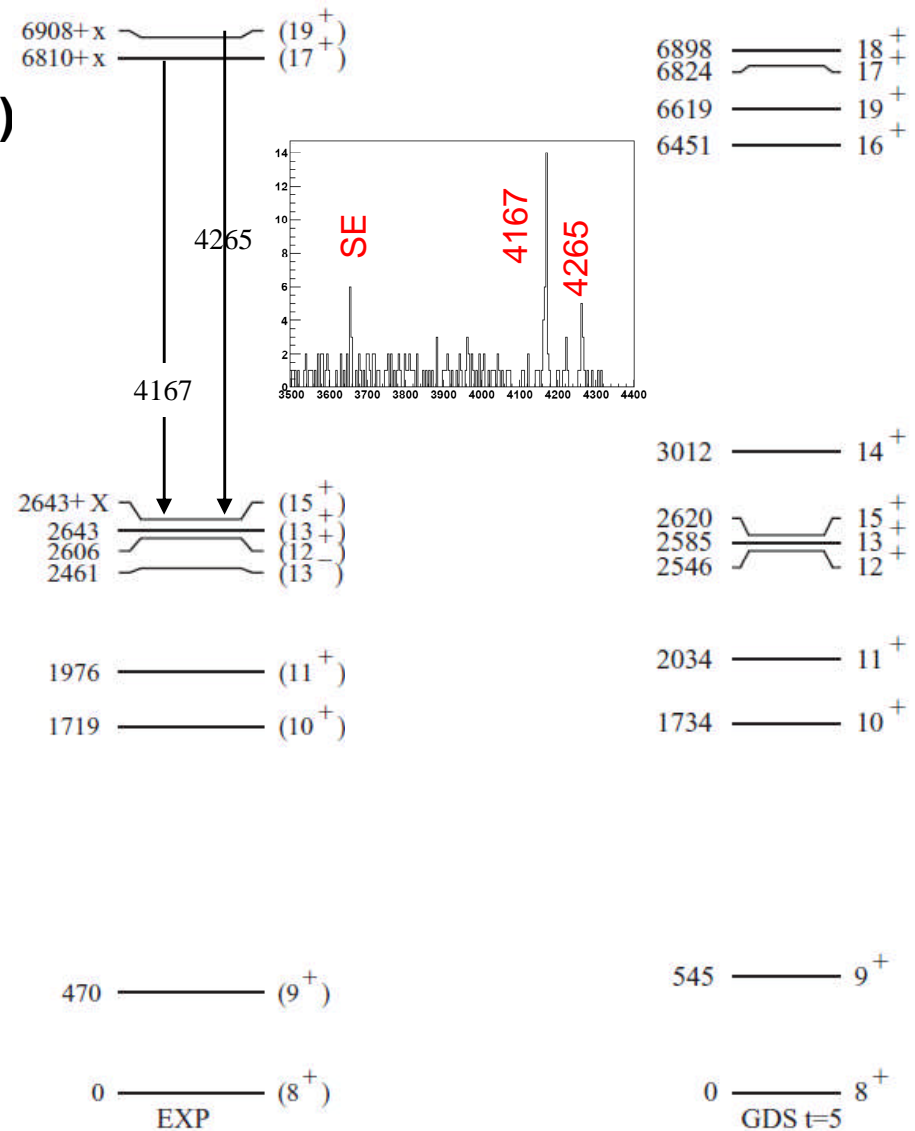
P Boutachkov et al., PRC 84, 044311 (2011)

GDS

LSSM calculations with
5p-5h, $t=5$, excitations
Antoine+Nathan codes

$J_i^{\pi_i}$	$J_f^{\pi_f}$	σL	$B(\sigma\lambda)$ W.u.	
			Expt	GDS
19 ⁺	17 ⁺	E2	4.7(10)	3.6
	15 ⁺	E4	0.9(6)	0.7

p/n eff. Charges of (1.5/0.5)e



Incorrect order of 17⁺, 19⁺ may result from incorrect details of the interaction. Fine tuning of residual p-h interaction is needed.

Similar effect noted for 10⁺, 12⁺ states in ^{98}Cd

— A. Blazhev et al., Phys Conf. Series 205, 012035 (2010)

Evidence for $T=0$ np Int. in ^{96}Cd

- Long standing SM predictions of the presence of 16^+ and $25/2^+$ spin-gap isomers in $^{96,97}\text{Cd}$, for example:

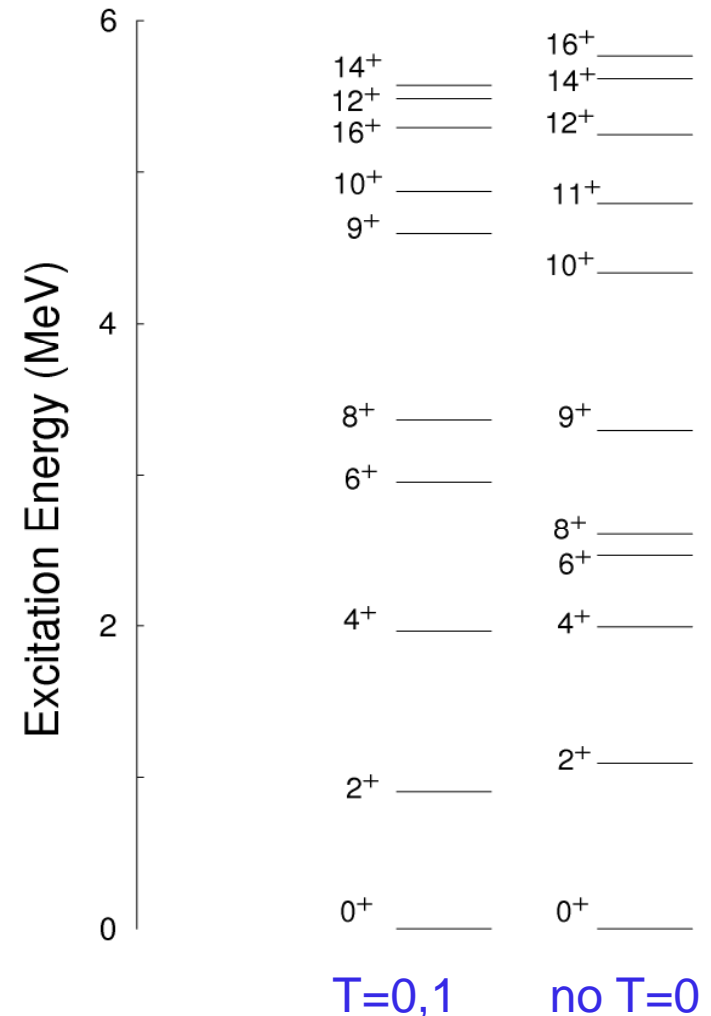
K Ogawa, Phys Rev C 28, 958 (1983)

- Spin gap isomer results from extra BE due to the large attractive n-p interaction for maximally aligned hole-hole configs.

Existence of isomer provides evidence for the importance of the $T = 0$ np interaction at high-spin

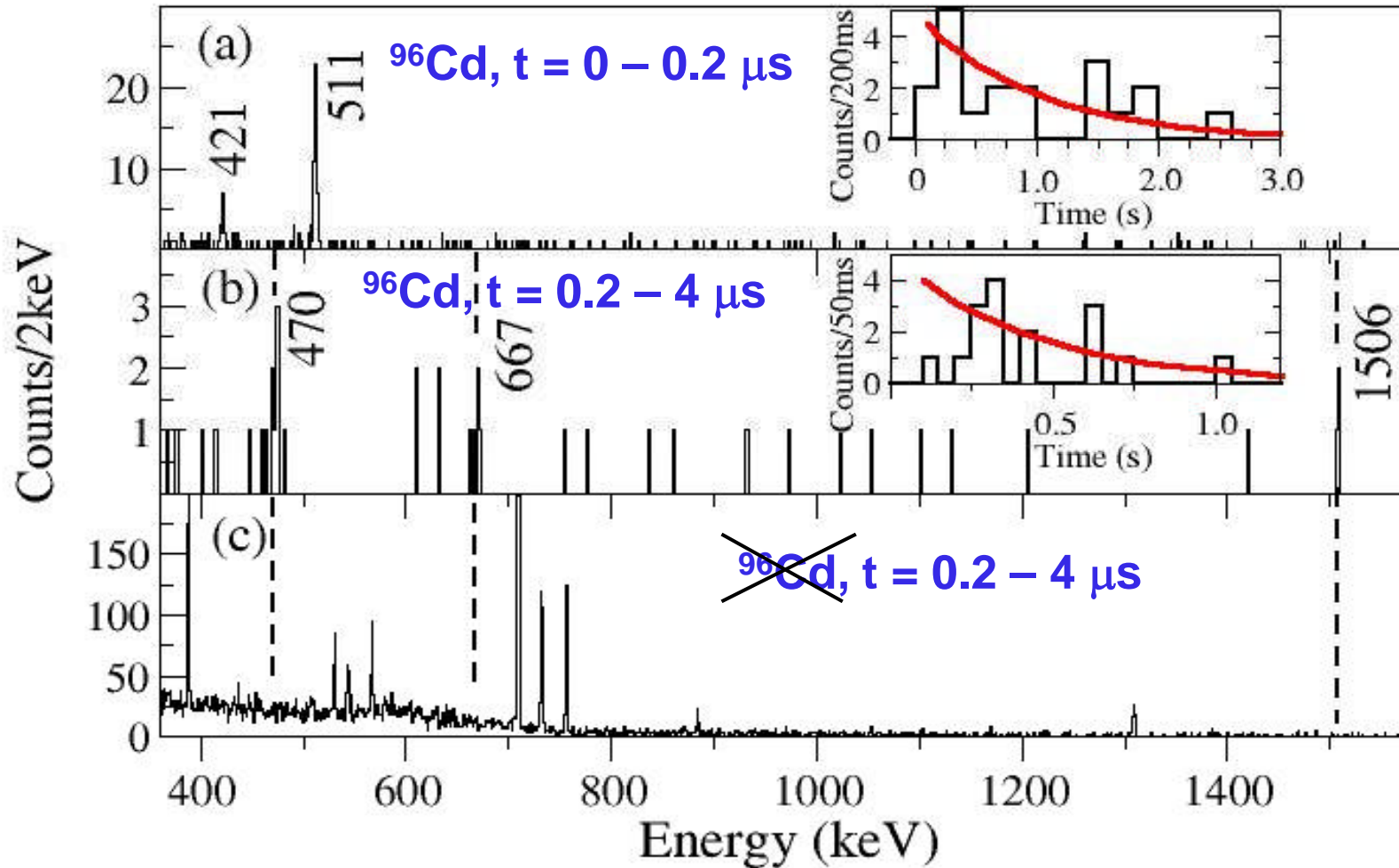
Cf ^{92}Pd results – B Cederwall et al., Nature 469, 68 (2011)

Gross and Frenkel, Nucl. Phys. A 267, 85 (1976)



SM Calculations by H Grawe
GF int., $g_{9/2}$, $p_{1/2}$ model space

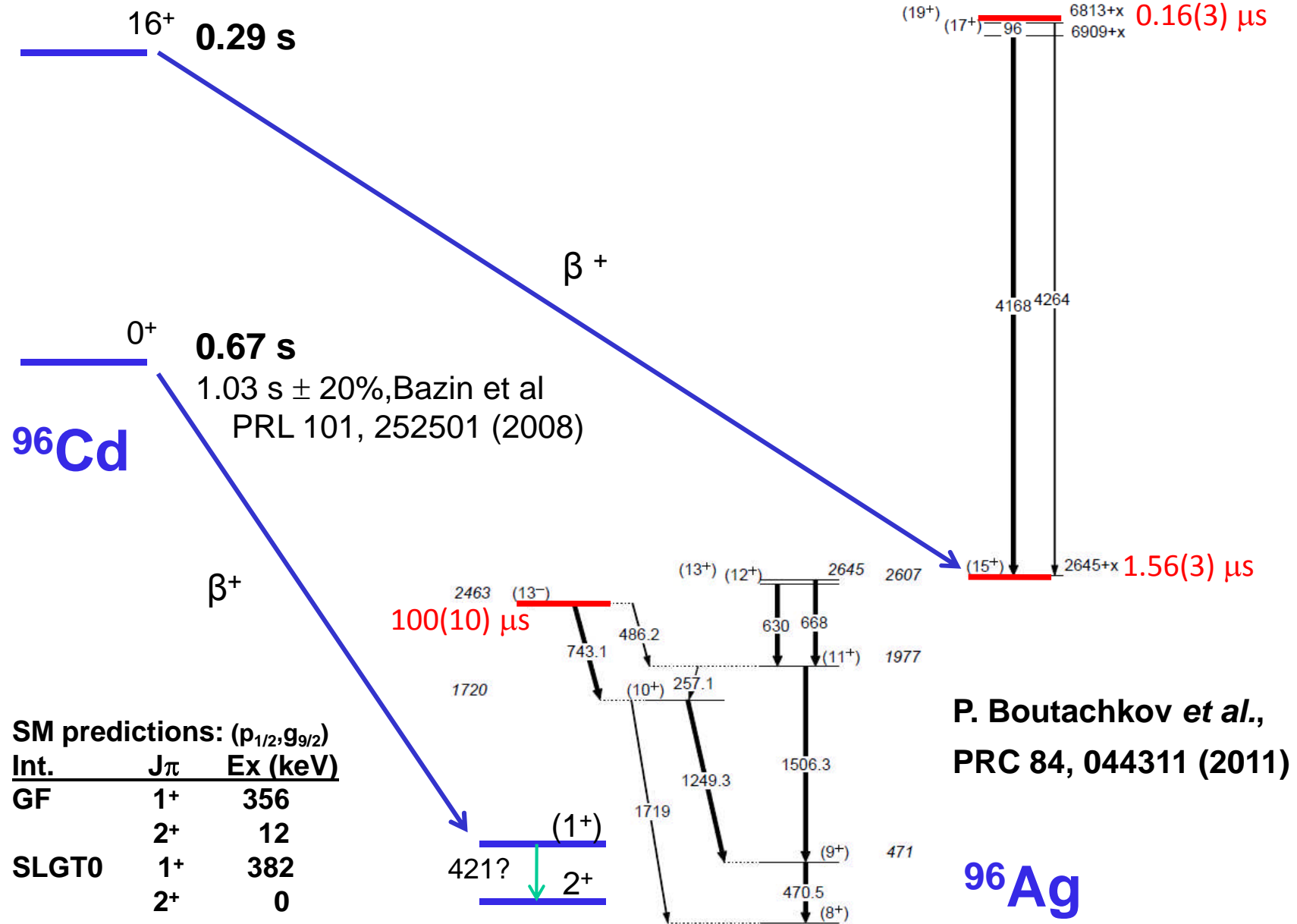
Spin-gap isomer ^{96}Cd



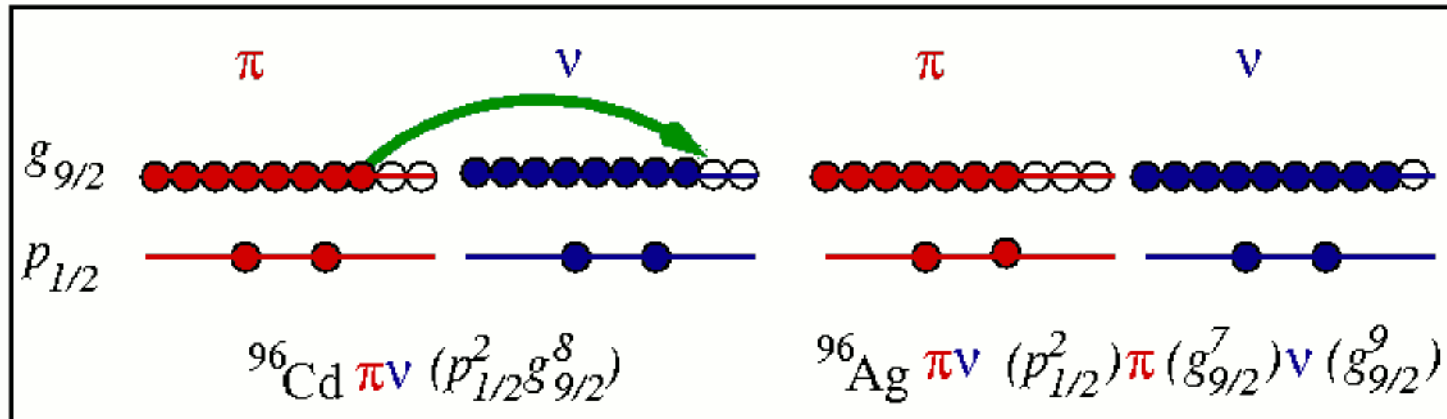
$T_{1/2}(421) = 0.67 \pm 0.15 \text{ sec}, \quad T_{1/2}(470, 1506, 667) = 0.29^{+0.11}_{-0.10} \text{ sec}$

B.S Nara Singh et al., PRL 107, 172502 (2011)

Spin-gap isomer ^{96}Cd



Spin-gap isomer ^{96}Cd



^{96}Ag , 15^+ state: 100% of GT strength in $p_{1/2}$, $g_{9/2}$ model space

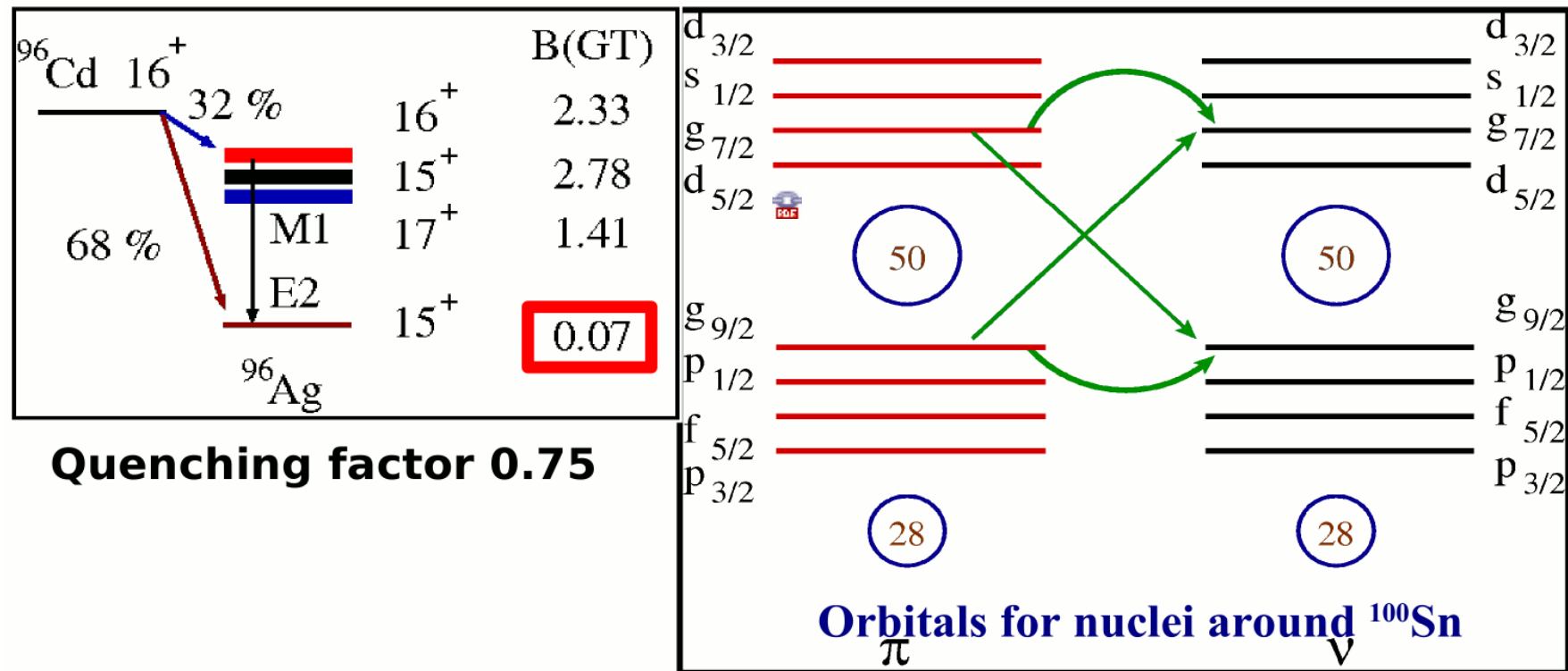
$B_{\text{GF}} = 0.14$ with quenching factor of 0.6 (Herndl and Brown NPA627, 35 (1997))

$$B_{\text{exp}} = [3860(18) * I_{\beta}] / (f T_{1/2}) = 0.19^{+0.08}_{-0.07}$$

with $T_{1/2} = 0.29$ secs

Spin-gap isomer ^{96}Cd – LSSM Calcs with Core excitations

GT strength is fragmented, due to the mixed nature of the states



High statistics are needed to obtain the B(GT) distribution

Summary

- **Mapping of collectivity along the $N=Z$ line is underway, but still lots to do:**
 - ⇒ *lifetime measurements/ mapping $B(E2)$ values*
 - ⇒ *role of shape co-existence in the mid-mass $A \sim 66-84$ region.*
- **Evidence that isoscalar np coupling is important at both low and high spin for $N=Z$ nuclei close to ^{100}Sn . But no direct evidence yet of $T=0$ np (BCS type) pair condensate.**
 - ⇒ *Need to measure lifetimes of low-lying states in $A \sim 90$ $N=Z$ nuclei and*
 - ⇒ *extend/ identify yrast bands in nuclei such as $^{92}\text{Pd}/^{96}\text{Cd}$ as well as investigate $T=0,1$ states in ^{90}Rh , ^{94}Ag , ^{98}In etc.*
- **Several isomers/ γ rays observed in $N \sim Z$ mass 90 nuclei in recent years, including core-excited states –**
 - ⇒ *these data provide stringent tests of model spaces and shell model interactions, but*
 - ⇒ *more data required to help tune the interactions used in SM calculations*

Significant interest to try and extend studies to $N < Z$ nuclei to investigate isomers/ isospin symmetry/ effects of weak binding in the mass 60-100 region.